



Strong Force Fields and Stability of the Nucleon and a singly heavy baryon

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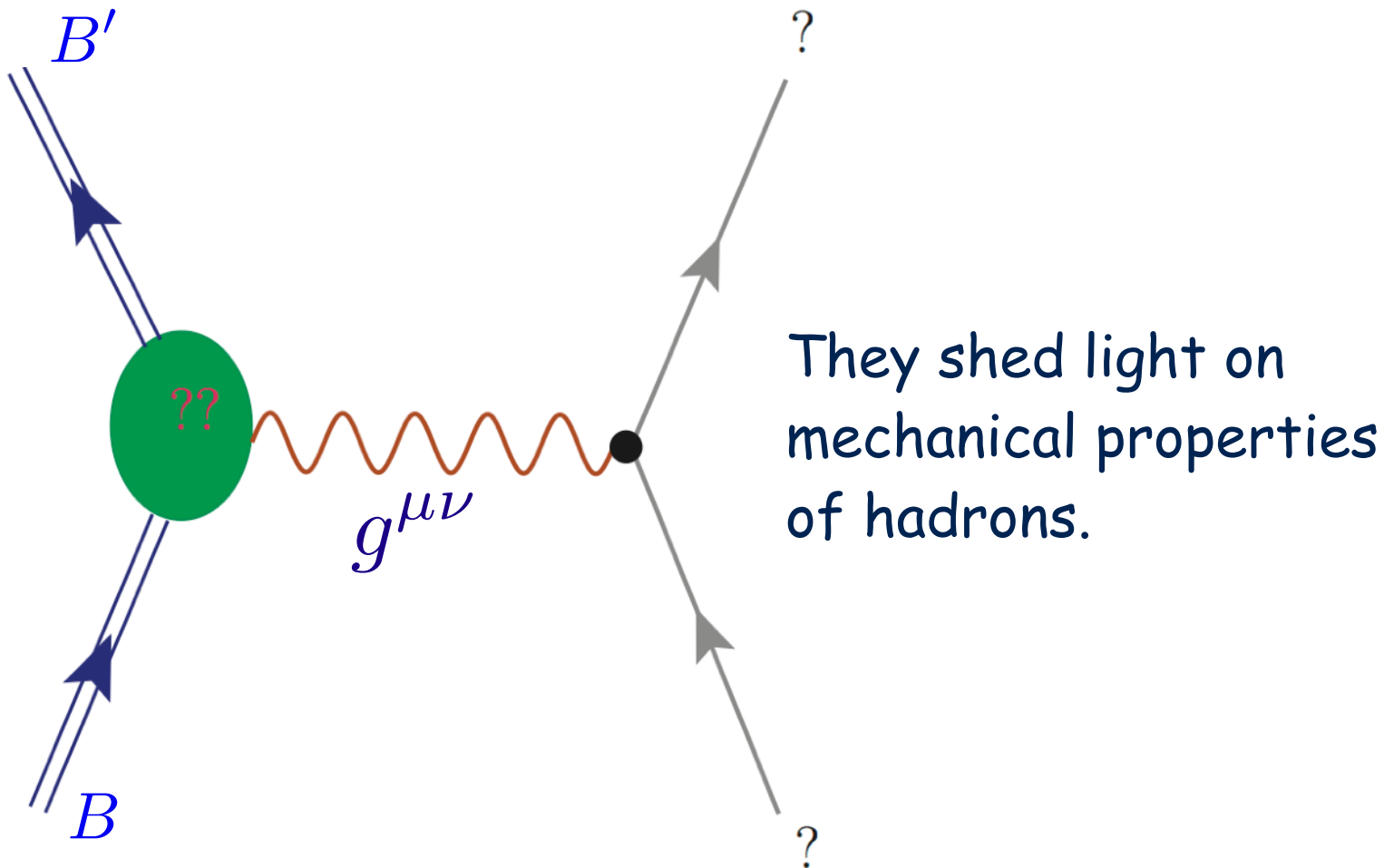
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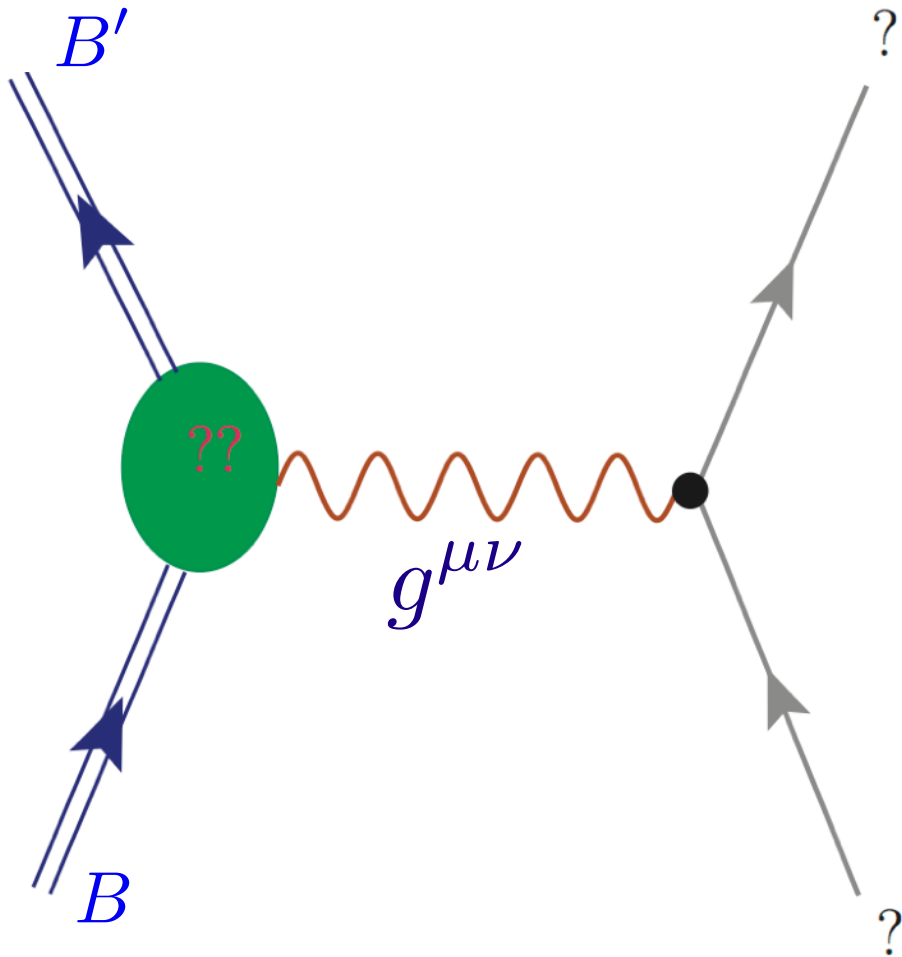
Gravitational form factors

- Why Gravitational form factors?



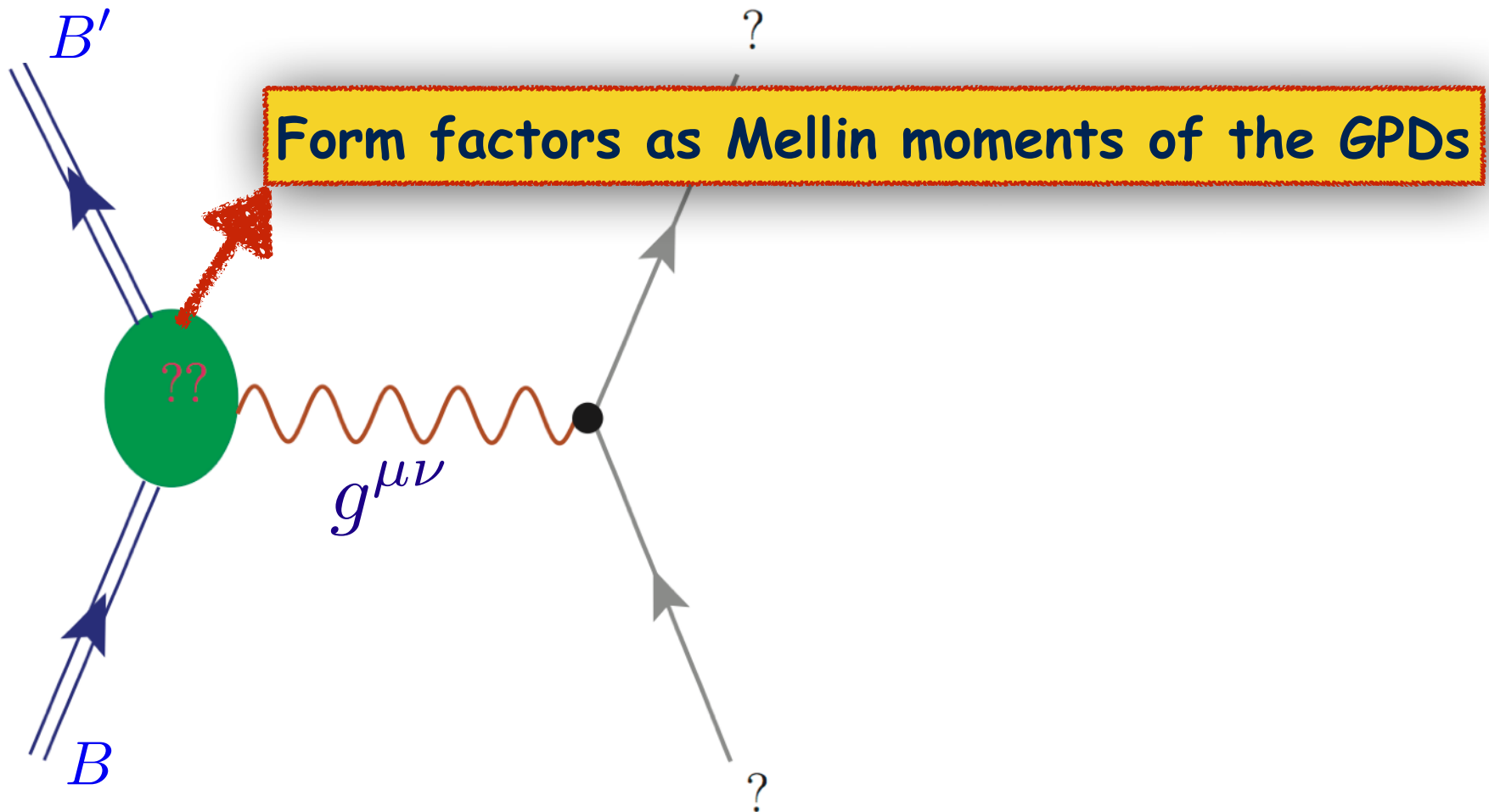
Gravitational form factors

Probes are unknown for the **Gravitational form factors!**



Gravitational form factors

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Gravitational form factors

- Unpolarized Generalized Parton Distributions

$$\int \frac{d\lambda}{2\pi} e^{i\lambda x} \langle p', \sigma' | \bar{\psi}_q \left(-\frac{\lambda n}{2} \right) \not{n} \psi_q \left(\frac{\lambda n}{2} \right) | p, \sigma \rangle$$

$$= \boxed{H^q(x, \xi, t)} \bar{u}(p', \sigma') \not{n} u(p, \sigma) + \boxed{E^q(x, \xi, t)} \bar{u}(p', \sigma') \frac{i\sigma^{\mu\nu} n_\mu \Delta_\nu}{2M_B} u(p, \sigma)$$

$$n^2 = 1$$

The first moment: Electromagnetic form factors

$$\Delta = (p' - p)$$

- The second moment: Gravitational form factors

$$n \cdot \Delta = -2\xi$$

$$\int_{-1}^1 dx x \sum_q H^q(x, \xi, t) = M_2^Q(t) + \frac{4}{5} d_1^Q(t) \xi^2,$$

$$\int_{-1}^1 dx x \sum_q E^q(x, \xi, t) = 2J^Q(t) - M_2^Q(t) - \frac{4}{5} d_1^Q(t) \xi^2$$

Gravitational form factors

$$\langle p', \sigma' | \hat{T}_{\mu\nu}^a(0) | p, \sigma \rangle = \bar{u}(p', \sigma') \left[M_2^a(t) \frac{P_\mu P_\nu}{M_B} + J^a(t) \frac{i(P_\mu \sigma_{\nu\rho} + P_\nu \sigma_{\mu\rho}) \Delta^\rho}{2M_B} \right. \\ \left. + d_1^a(t) \frac{\Delta_\mu \Delta_\nu - g_{\mu\nu} \Delta^2}{5M_B} + \bar{c}^a(t) M_B g_{\mu\nu} \right] u(p, \sigma)$$

- Components of the Energy-momentum tensor

$$\int d^3r \sum_{a=Q,G} T_{00}^a(\mathbf{r}, \sigma', \sigma) = M_B M_2(0) = M_B, \quad \text{: mass of a baryon}$$

$$\int dr^3 \sum_{a=Q,G} J_{\sigma'\sigma}^{a,i}(\mathbf{r}) = 2\hat{S}_{\sigma'\sigma}^i J(0) = \hat{S}_{\sigma'\sigma}^i \quad \text{: spin of a baryon}$$

Strong force field and Stability

- Conservation of the EMT

$$\partial^i T_{ij} = \frac{r_j}{r} \left[\frac{2}{3} \frac{\partial s(r)}{\partial r} + \frac{2s(r)}{r} + \frac{\partial p(r)}{\partial r} \right] = 0$$

$p(r)$: Pressure density

$s(r)$: Shear-force density

- Global Stability condition (von Laue condition)

$$\int_0^\infty dr r^2 p(r) = 0$$

Strong force field and Stability

- Strong force fields

$$dF_{(r,\theta,\phi)}^i = T^{ij} dS_{(r,\theta,\phi)} e^j_{(r,\theta,\phi)}$$

$$\frac{dF_r}{dS_r} = \frac{2}{3}s(r) + p(r), \quad \frac{dF_\theta}{dS_\theta} = \frac{dF_\phi}{dS_\phi} = -\frac{1}{3}s(r) + p(r)$$

- Local stability condition

$$\frac{2}{3}s(r) + p(r) > 0$$

- 2D von Laue stability condition

$$\int_0^\infty dr \, r \left(-\frac{1}{3}s(r) + p(r) \right) = 0$$

Baryon in pion mean fields

- * A **baryon** can be viewed as a state of N_c quarks bound by mesonic **mean fields** (E. Witten, NPB, 1979 & 1983).

Its mass is proportional to N_c , while its width is of order $O(1)$.

- Mesons are weakly interacting (Quantum fluctuations are suppressed by $1/N_c$: $O(1/N_c)$).

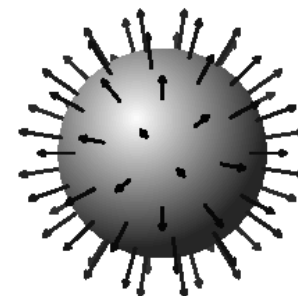
Meson mean-field approach (Chiral Quark-Soliton Model)

- * Baryons as a state of N_c quarks bound by mesonic mean fields.

$$S_{\text{eff}} = -N_c \text{Tr} \ln (i\not{D} + iMU\gamma^5 + i\hat{m})$$

- * **Key point: Hedgehog Ansatz**

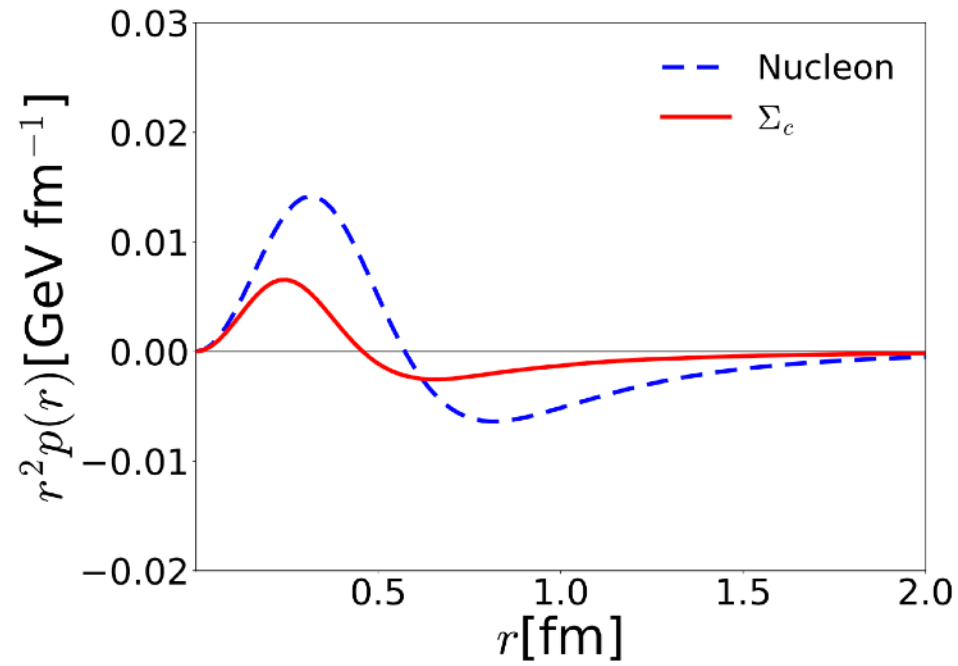
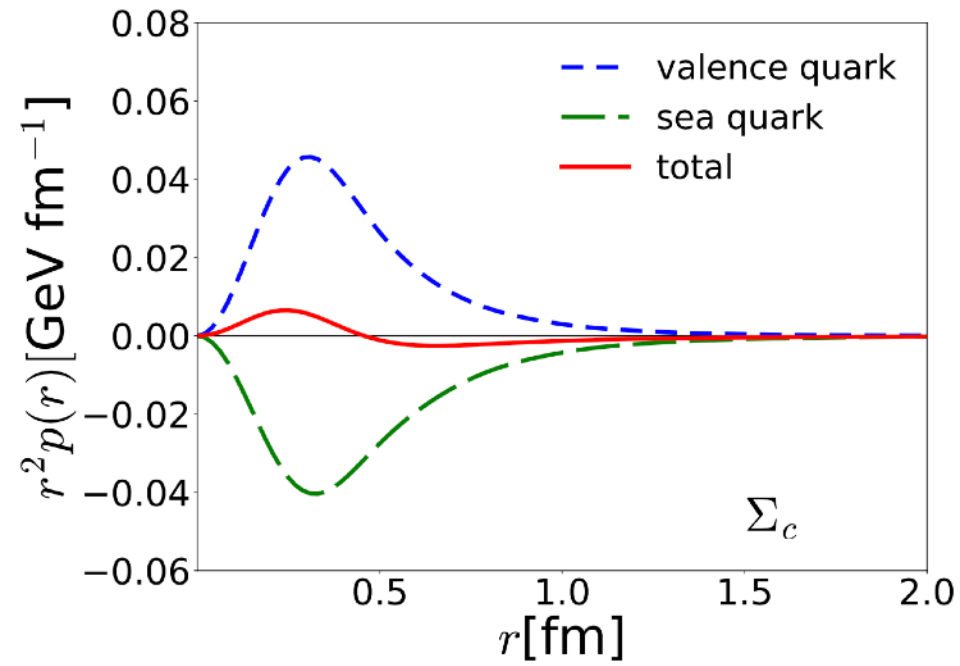
$$\pi^a(\mathbf{r}) = \begin{cases} n^a F(r), & n^a = x^a/r, \quad a = 1, 2, 3 \\ 0, & a = 4, 5, 6, 7, 8. \end{cases}$$



hedgehog

- It breaks spontaneously $SU(3)_{\text{flavor}} \otimes O(3)_{\text{space}} \rightarrow SU(2)_{\text{isospin+space}}$

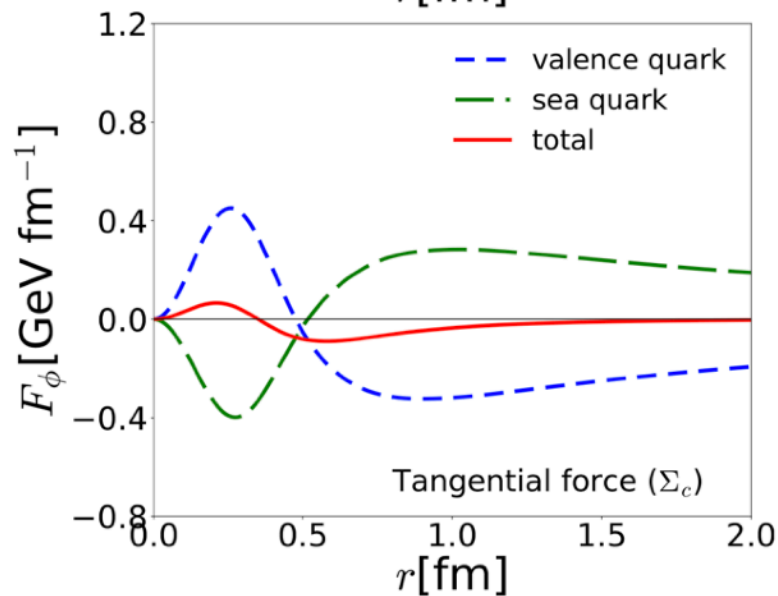
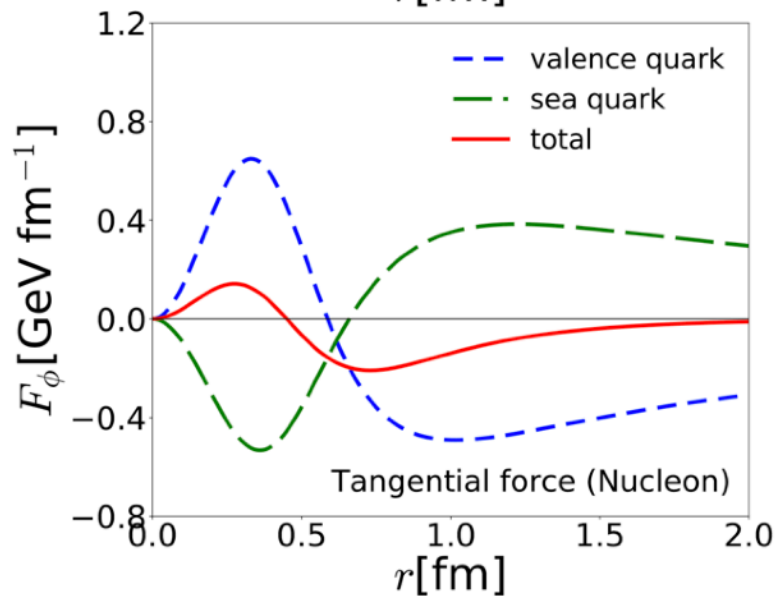
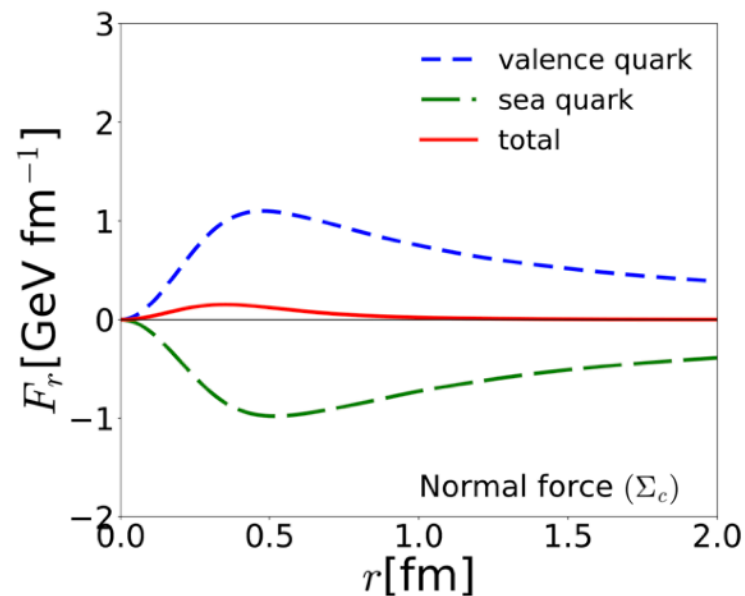
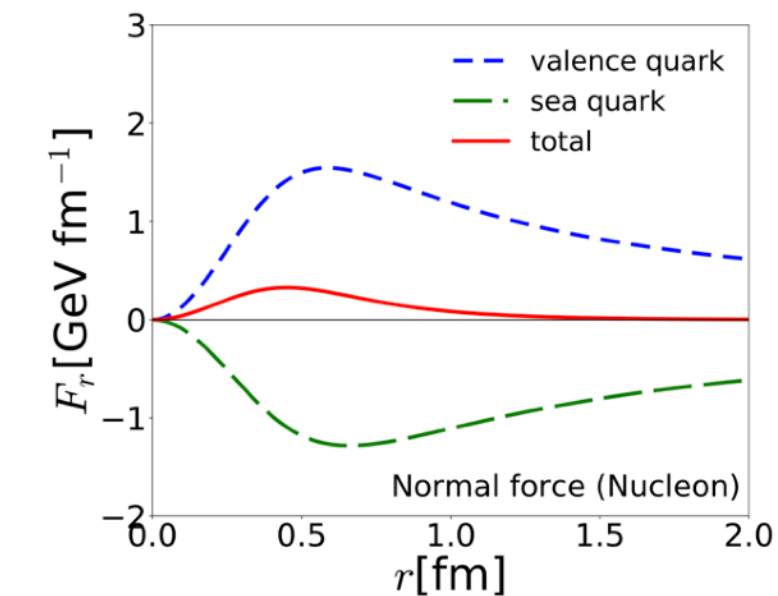
Pressure density



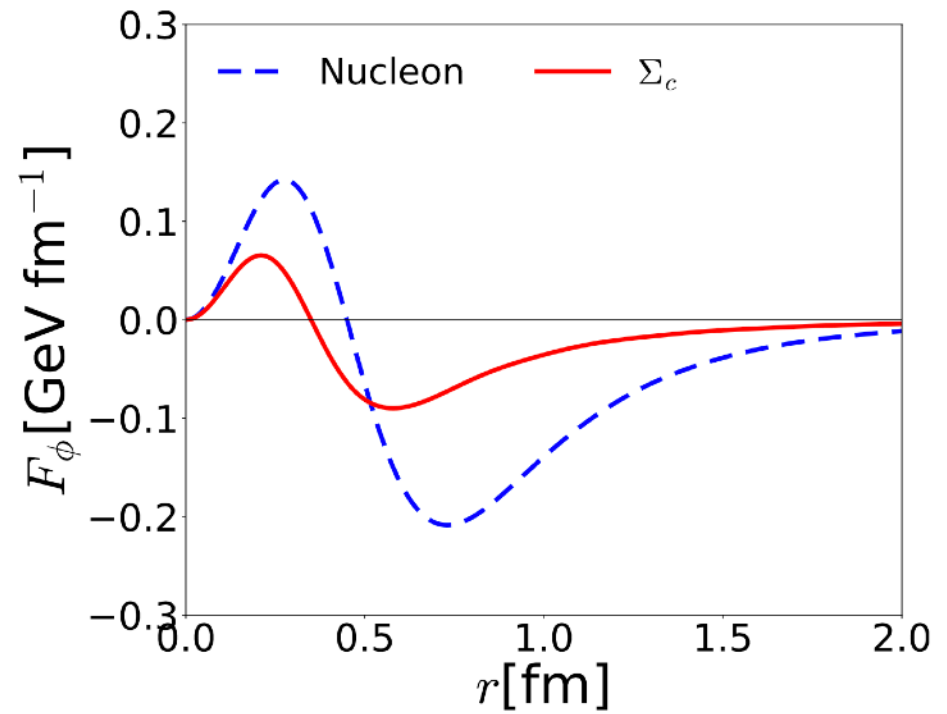
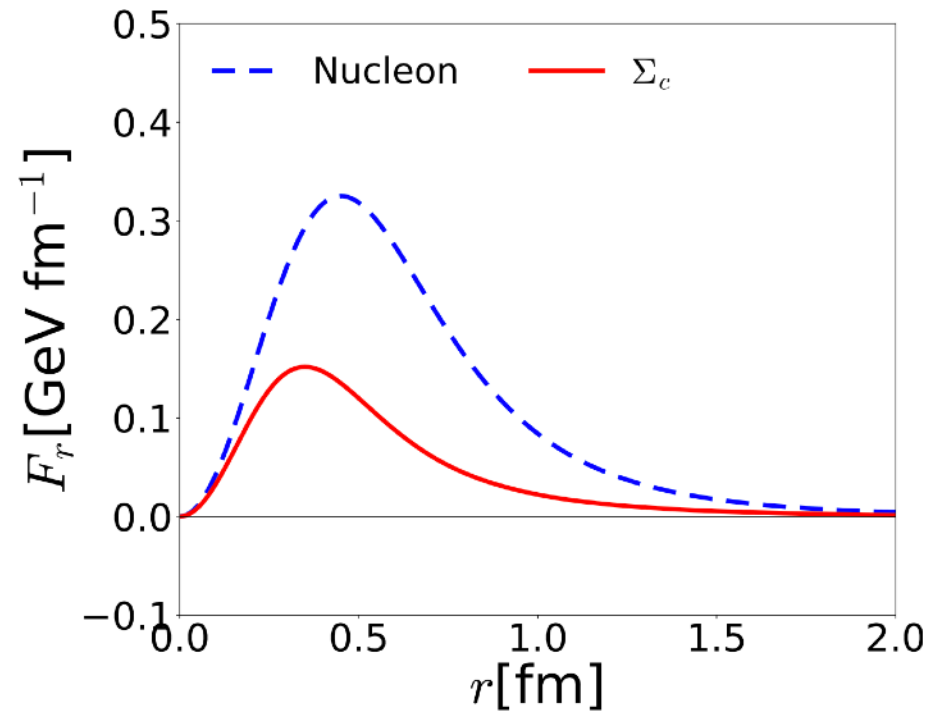
$$\int_0^{\infty} dr r^2 p(r) = 0$$

This condition is a nontrivial one.

Strong force field

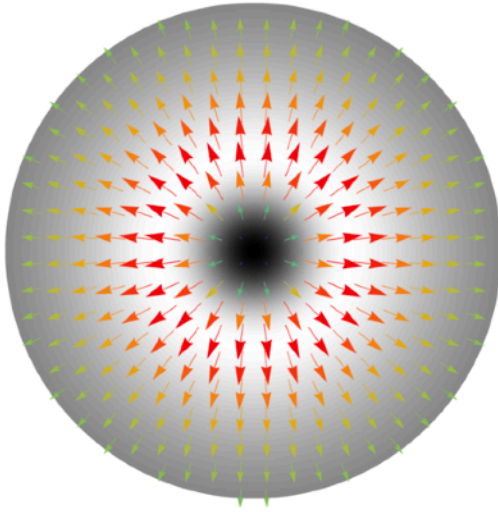


Strong force fields

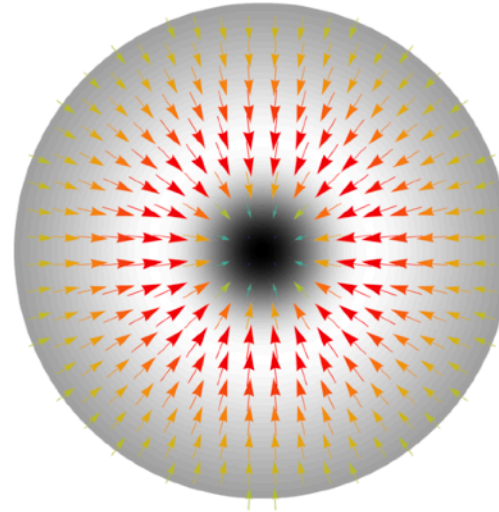


Strong force fields

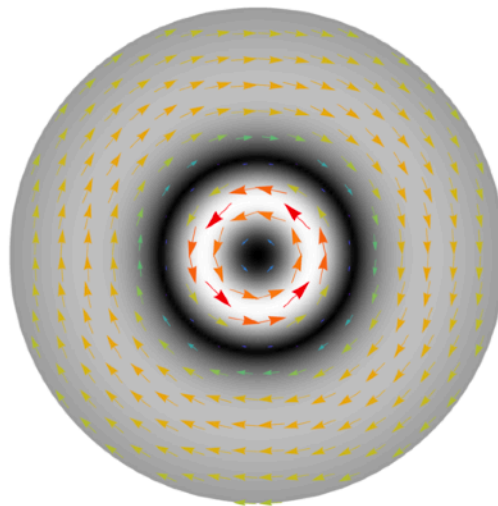
Normal force in the nucleon (valence quark)



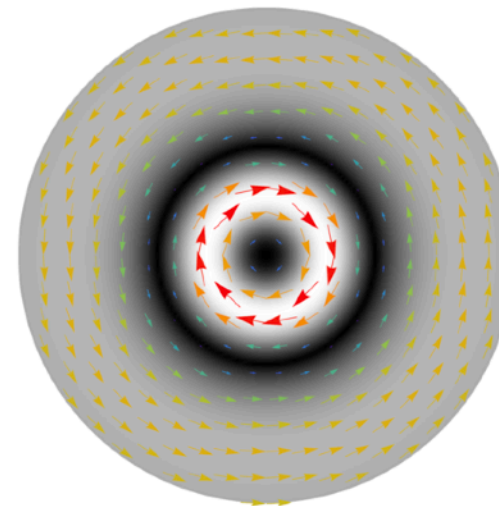
Normal force in the nucleon (sea quark)



Tangential force in the nucleon (valence quark)

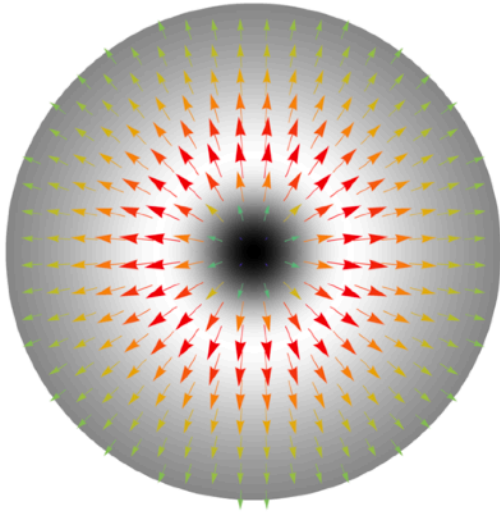


Tangential force in the nucleon (sea quark)

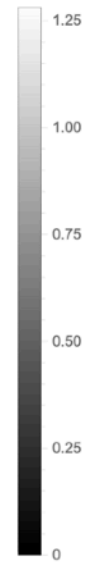
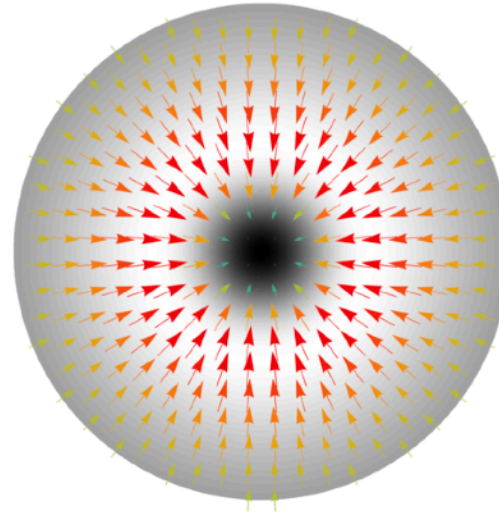


Strong force fields

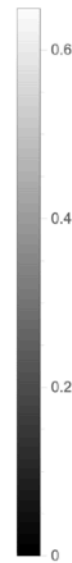
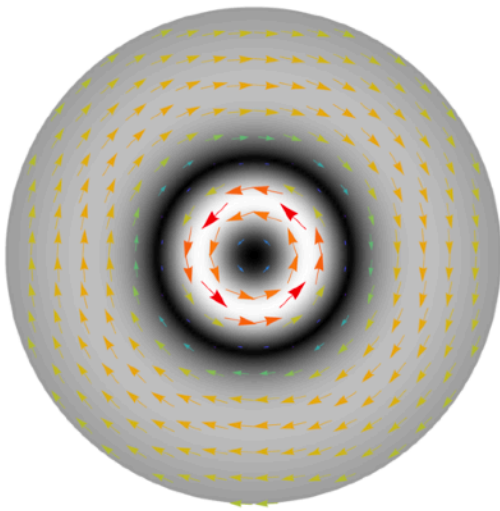
Normal force in Σ_c (valence quark)



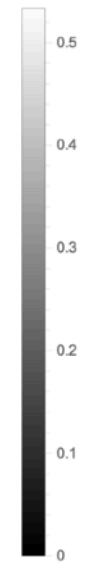
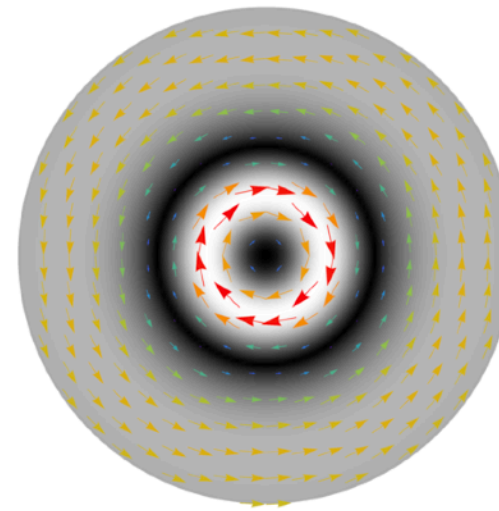
Normal force in Σ_c (sea quark)



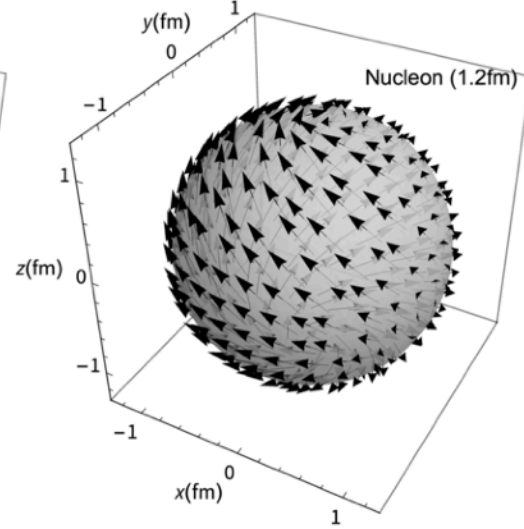
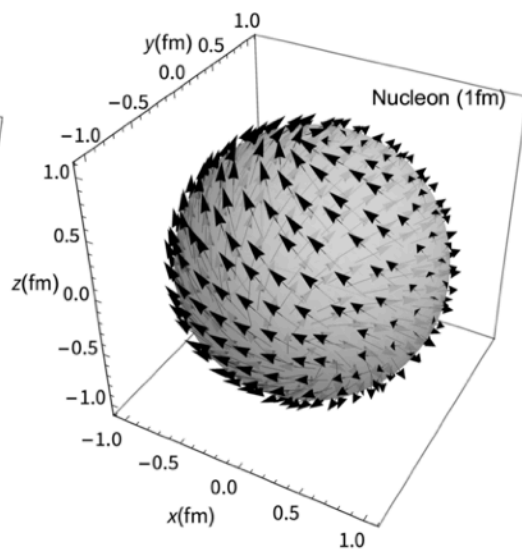
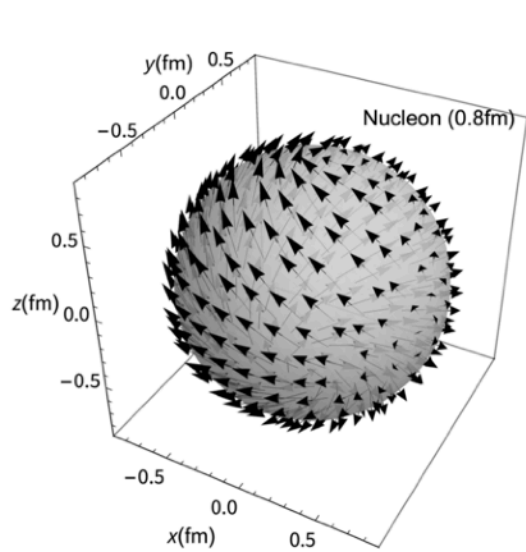
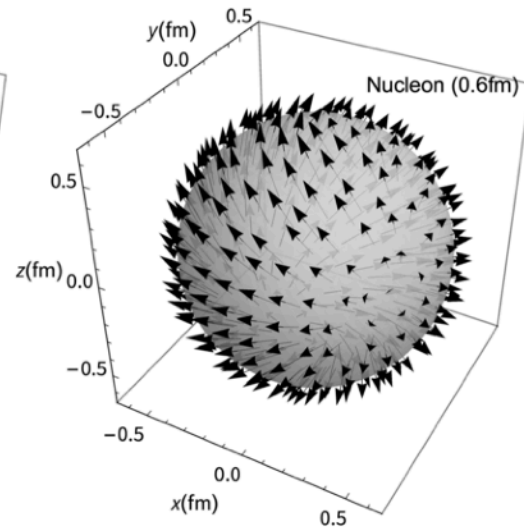
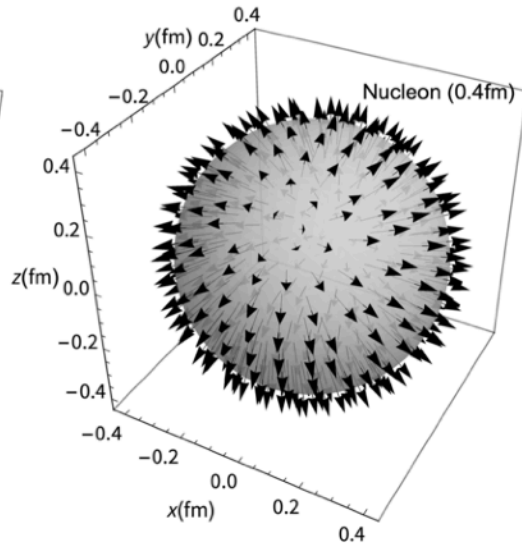
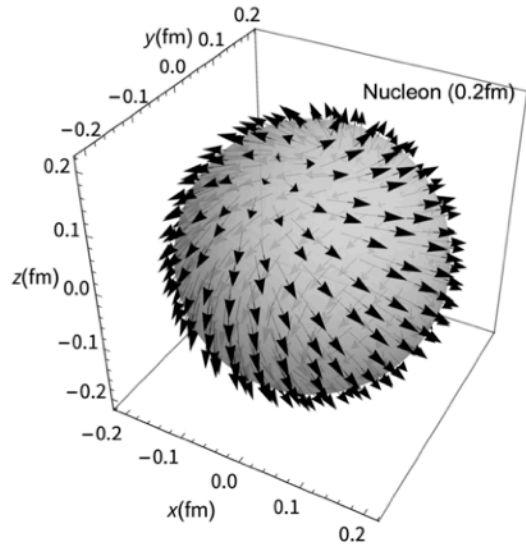
Tangential force in Σ_c (valence quark)



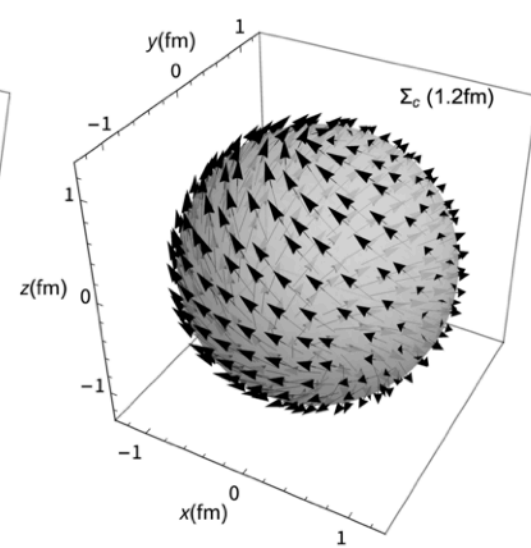
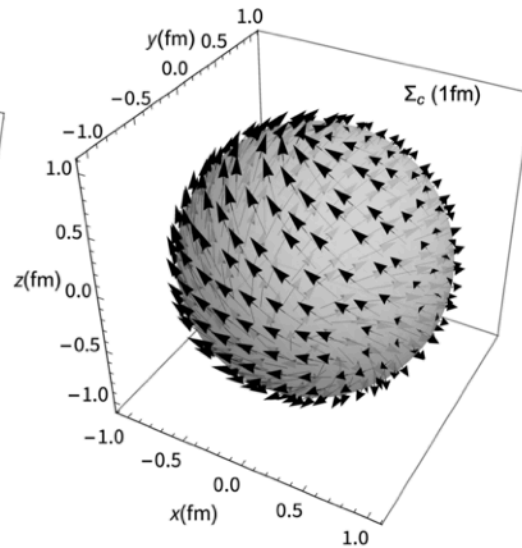
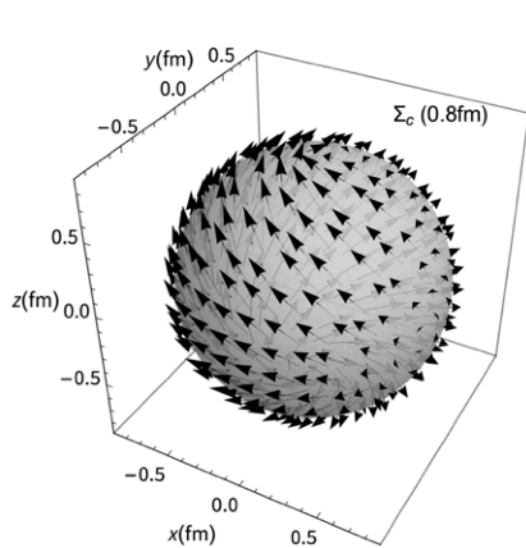
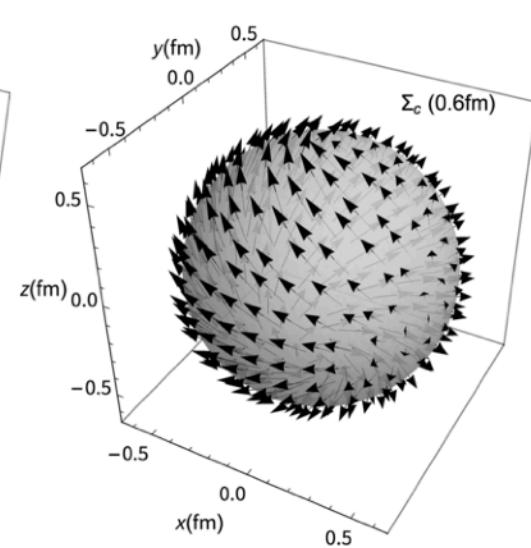
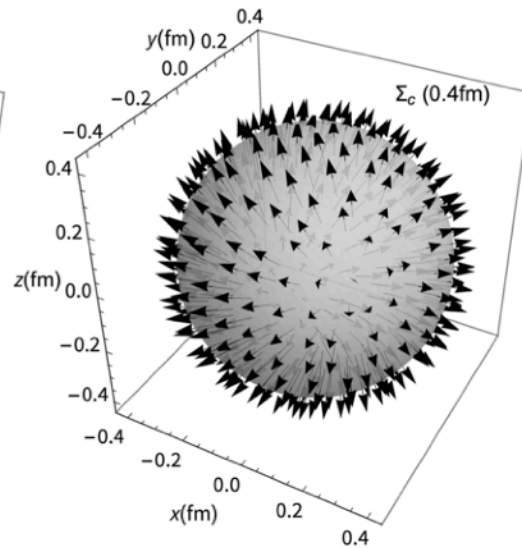
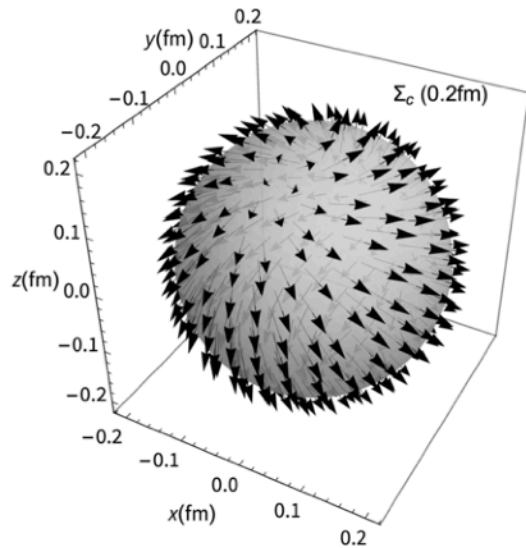
Tangential force in Σ_c (sea quark)



Strong force fields



Strong force fields



Summary & Outlook

Summary

- In the present talk, we aimed at showing how the nucleon and singly heavy baryons acquire their stabilities as particles.
- As shown in the results, the vacuum fluctuations are essential to secure the stabilities of baryons.
- This may give a clue on the mechanism of the quark confinement.
- The form factors were also computed (see the backup slide).

One can construct a theory of **baryon seismology**, based on the present work.

**Though this be madness,
yet there is method in it.**

Hamlet Act 2, Scene 2

by Shakespeare

Thank you very much for the attention!

Gravitational form factors

