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## Introduction

- The axial-vector current examines various sides of baryon properties (e.g. spin content, decay width, etc)
- The axial-vector structures of the baryon decuplet are not well known.
- Mass difference of light baryons, electromagnetic properties etc were described well in the chiral quark-soliton model.
- We would like to explain the axial-vector form factors of the baryon decuplet within the chiral quark-soliton model.



Sea quark contribution

#### Valence quark contribution

### Framework



 Baryons can be considered as Nc valence quarks, which are bound by the mesonic mean-fields at large Nc (E. Witten, NPB160, 57 (1979)).

Effective chiral action:  $S_{\text{eff}} = -N_c \text{Tr} \ln \left[ i \gamma^{\mu} \partial_{\mu} + i \hat{m} + i M U^{\gamma_5} \right]$ 

- The Hedgehog ansatz is applied to the pseudo-Nambu-Goldston boson field.
- The mean-field can be found by solving the equations of motion self-consistently.
- Witten's trivial embedding is used to preserve the hedgehog symmetry in the flavor SU(3) (E. Witten, NPB223, 422 (1983)).
- We take into account the zero-mode quantization of the soliton (Rotational and translational zero modes).

$$\frac{\delta E_{cl}}{\delta U} = 0 \longrightarrow E_{cl} P(r)$$

 $U_{SU(2)}(r) = e^{i\mathbf{n}\cdot\boldsymbol{\tau}P(r)/f_{\pi}}$ 

$$U(r) = \begin{pmatrix} U_{SU(2)}(r) & 0\\ 0 & 1 \end{pmatrix}.$$

 $U(\boldsymbol{x},t) = A(t)U_c(\boldsymbol{x} - \boldsymbol{Z}(t))A^{\dagger}(t)$ 

### Framework

$$\begin{split} & \left\{ \begin{array}{c} & \left\{ B_{J'}^{(10)}(p') \left| A_{a}^{\mu}(0) \right| B_{J}^{(10)}(p) \right\} \bullet \\ & \left\{ g_{1}^{(a)B}(q^{2})\eta_{\alpha\beta} + h_{1}^{(a)B}(q^{2})\frac{q_{\alpha}q_{\beta}}{4M_{B}^{2}} \right\} \\ & + \frac{q_{\mu}}{2M_{B}} \left\{ g_{3}^{(a)B}(q^{2})\eta_{\alpha\beta} + h_{3}^{(a)B}(q^{2})\frac{q_{\alpha}q_{\beta}}{4M_{B}^{2}} \right\} \right] \gamma^{5}u^{\beta}(p,J) \\ & \left\{ \int dA \int d^{3}z \, e^{i\vec{q}\cdot\vec{z}} \langle B_{J'}^{\prime(10)} \left| A \rangle \mathcal{F}_{\mu}^{a}(\vec{z},A) \langle A \left| B_{J}^{(10)} \right\rangle \right\} \\ & \left\{ g_{1}^{(a)B}(Q^{2}) = \frac{M_{B}}{E} \int d^{3}r \langle B(p',\frac{3}{2}) \right| \left[ j_{0}(Q|\mathbf{r}|) \{ \hat{\mathcal{F}}_{A}^{a}(\mathbf{r}) \}_{10} - j_{2}(Q|\mathbf{r}|) \left\{ \sqrt{2\pi}Y_{2} \otimes \hat{\mathcal{F}}_{A}^{a}(\mathbf{r}) \right\}_{10} \right] |B(p,\frac{3}{2}) \rangle \\ & g_{3}^{(a)B}(Q^{2}) = -\frac{4M_{B}^{2}}{EQ^{2}} \int d^{3}r \langle B(p',\frac{3}{2}) | \left[ (E-M_{B})j_{0}(Q|\mathbf{r}|) \{ \hat{\mathcal{F}}_{A}^{a}(\mathbf{r}) \}_{10} \right] \\ & + (2E+M_{B})j_{2}(Q|\mathbf{r}|) \left\{ \sqrt{2\pi}Y_{2} \otimes \hat{\mathcal{F}}_{A}^{a}(\mathbf{r}) \right\}_{10} \right] |B(p,\frac{3}{2}) \rangle \end{split}$$

The form factor  $h_{1,3}^{(a)B}(q^2)$  are in fact the same as  $g_{1,3}^{(a)B}(q^2)$  apart from the kinematical factors

$$g_{1(3)}^{(a)B}(Q^2) = (g_{1(3)}^{(a)B}(Q^2))^{(\text{sym})} + (g_{1(3)}^{(a)B}(Q^2))^{(\text{op})} + (g_{1(3)}^{(a)B}(Q^2))^{(\text{wf})}$$

The flavor SU(3) symmetry breaking contributions



The triplet axial-vector form factors, a=3





The singlet axial-vector form factors, a=0





#### The octet axial-vector form factors, a=8





#### Alexandrou et al, PRD87, 114513, 2013





$g_1^{(3)B}(0)$	$\Delta^{++}$	$\Delta^+$	$\Delta^0$	$\Delta^{-}$	$\Sigma^{*+}$	$\Sigma^{*0}$	$\Sigma^{*-}$	$\Xi^{*0}$	Ξ*-	$\Omega^{-}$
$m_{ m s}=0~{ m MeV}$	2.0064	0.6688	-0.6688	-2.0064	1.338	0	-1.338	0.669	-0.669	0
$m_{ m s}=180~{ m MeV}$	2.1333	0.7111	-0.7111	-2.1333	1.440	0	-1.440	0.729	-0.729	0
LQCD [4] $(m_{\pi} = 131.2(13) \text{ MeV})$	-	-	_	_	1.1740(380)	-	-	0.5891(198)	-	_
LQCD [4] $(m_{\pi} = 213 \text{ MeV})$	1.9777(1458)	0.5181(981)	-0.6499(973)	-1.7090(1422)	1.1929(521)	-0.1367(685)	-1.2633(516)	0.5869(216)	-0.6682(382)	_
LQCD [4] $(m_{\pi} = 256 \text{ MeV})$	1.6956(1897)	0.5670(1479)	-0.5929(1167)	-1.7322(1718)	1.1462(720)	0.0148(542)	-1.0646(661)	0.5785(278)	-0.5424(303)	_
LQCD [3] $(m_{\pi} = 297 \text{ MeV})$	-	0.604(38)	_	_	-	-	-	-	-	_
LQCD [4] $(m_{\pi} = 302 \text{ MeV})$	1.9574(1552)	0.6374(976)	-0.4798(1063)	-1.4374(1331)	1.2839(636)	0.0654(444)	-1.0423(619)	0.6204(256)	-0.5459(299)	_
LQCD [3] $(m_{\pi} = 353 \text{ MeV})$	-	0.640(26)	_	_	_	_	_	_	_	_
LQCD [4] $(m_{\pi} = 373 \text{ MeV})$	1.7602(1035)	0.5215(639)	-0.5676(635)	-1.5872(1270)	1.1478(558)	-0.0130(323)	-1.1139(485)	0.5741(243)	-0.5702(230)	_
LQCD [3] $(m_{\pi} = 411 \text{ MeV})$	-	0.571(18)	_	_	_	_	_	_	_	_
LQCD [4] $(m_{\pi} = 432 \text{ MeV})$	1.8520(875)	0.6129(478)	-0.5949(489)	-1.8108(868)	1.2228(473)	0.0124(244)	-1.1765(450)	0.6059(213)	-0.5885(223)	_
LQCD [3] $(m_{\pi} = 490 \text{ MeV})$	-	0.578(13)	_	_	_	-	_	_	-	_
LQCD [3] $(m_{\pi} = 563 \text{ MeV})$	_	0.5887(98)	_	_	_	_	_	_	_	_
$\chi \mathrm{PT}~[6]$	$2.25^{*}$	-	_	_	_	-	-	_	_	_
RCQM[GBE] [7, 8]	$2.24^{*}$	_	_	_	$1.499^\dagger$	_	_	$0.75^{\dagger}$	_	_
LCSR [9]	$2.70\pm0.6^*$	_	_	_	_	_	_	_	_	_
PCQM [10]	$1.863^{*}$	-	-	-	$1.242^\dagger$	-	_	$0.621^\dagger$	-	-

[3] C. Alexandrou et al, PRD 87, 114513 (2013).
[4] C. Alexandrou et al, PRD 94, 034502 (2016).
[6] F. Jiang and B. C. Tiburzi, PRD 78, 017504 (2008).
[7,8] Ki-Seok Choi et al, PRD 82, 014007 (2010), FBS54, 1055 (2013).
[9] A. Kucukarslan et al, PRD 90, 054002 (2014).
[10] X. Y. Liu et al, PRC 97, 055206 (2018).

### Additional numerical results

$m_{\rm s}=180~{\rm MeV}$	$\Delta^{++}$	$\Delta^+$	$\Delta^0$	$\Delta^{-}$	$\Sigma^{*+}$	$\Sigma^{*0}$	$\Sigma^{*-}$	$\Xi^{*0}$	$\Xi^{*-}$	$\Omega^{-}$
$g_3^{(3)B}(0)$	346.1	115.4	-115.4	-346.1	303.9	0	-303.9	193.7	-193.7	0
$g_3^{(0)B}(0)$	7.822	7.822	7.822	7.822	1.622	1.622	1.622	-8.204	-8.204	-21.936
$g_3^{(8)B}(0)$	50.8	50.8	50.8	50.8	-60.0	-60.0	-60.0	-251.9	-251.9	-542.8
$\langle r_A^2 \rangle_B \; [{\rm fm}^2]$	0.447	0.447	0.447	0.447	0.438	_	0.438	0.431	0.431	_
$M_A \; [\text{GeV}]$	1.023	1.023	1.023	1.023	1.033	-	1.033	1.041	1.041	-

$$g_1^{(3)B}(Q^2) = \frac{g_1^{(3)B}(0)}{\left(1 + \frac{Q^2}{M_A^2}\right)^2}, \quad \langle r_A^2 \rangle_B = \frac{-6}{g_1^{(3)B}(0)} \left. \frac{\partial g_1^{(3)B}(Q^2)}{\partial Q^2} \right|_{Q^2 = 0}$$

## Summary & Outlook

#### • Summary

- We performed the axial-vector form factor calculations within the chiral quarksoliton model.
- We discussed the flavor SU(3) symmetry breaking contributions to the axialvector form factors.
- We compared the axial-vector form factors and constants with Lattice data and other model calculations.
- We have calculated the axial radii and masses for baryon decuplet as well.
- Outlook
- The axial-vector form factors with the quark contents.
- Transition form factors between the light baryons.
- The Axial-vector form factors for the heavy baryons.
- Tensor form factors.

## Thank you for listening!!!





C. Alexandrou et al, PRD 94, 034502 (2016).



C. Alexandrou et al, PRD 94, 034502 (2016).



C. Alexandrou et al, PRD 94, 034502 (2016).



**Triplet axial-vector form factors** 













## Axial-vector constants for the baryon decuplet

$g_1^{(3)B}(0)$	$\Delta^{++}$	$\Delta^+$	$\Delta^0$	$\Delta^{-}$	$\Sigma^{*+}$	${\Sigma^{*}}^{0}$	$\Sigma^{*-}$	$\Xi^{*0}$	Ξ*-	$\Omega^{-}$
$m_{ m s}=0~{ m MeV}$	1.9671	0.6557	-0.6557	-1.9671	1.3114	0	-1.3114	0.6557	-0.6557	0
$m_{ m s}=180~{ m MeV}$	1.9876	0.6625	-0.6625	-1.9876	1.3279	0	-1.3279	0.6654	-0.6654	0
LQCD [4] $(m_{\pi} = 131.2(13) \text{ MeV})$	-	-	_	_	1.1740(380)	-	-	0.5891(198)	-	_
LQCD [4] $(m_\pi = 213 \text{ MeV})$	1.9777(1458)	0.5181(981)	-0.6499(973)	-1.7090(1422)	1.1929(521)	-0.1367(685)	-1.2633(516)	0.5869(216)	-0.6682(382)	_
LQCD [4] $(m_\pi = 256 \text{ MeV})$	1.6956(1897)	0.5670(1479)	-0.5929(1167)	-1.7322(1718)	1.1462(720)	0.0148(542)	-1.0646(661)	0.5785(278)	-0.5424(303)	_
$\mathrm{LQCD}~[3]~(m_{\pi}=297~\mathrm{MeV})$	_	0.604(38)	_	_	_	_	_	_	_	_
LQCD [4] $(m_{\pi} = 302 \text{ MeV})$	1.9574(1552)	0.6374(976)	-0.4798(1063)	-1.4374(1331)	1.2839(636)	0.0654(444)	-1.0423(619)	0.6204(256)	-0.5459(299)	_
LQCD [3] $(m_{\pi} = 353 \text{ MeV})$	-	0.640(26)	_	_	_	-	-	_	_	_
LQCD [4] $(m_{\pi} = 373 \text{ MeV})$	1.7602(1035)	0.5215(639)	-0.5676(635)	-1.5872(1270)	1.1478(558)	-0.0130(323)	-1.1139(485)	0.5741(243)	-0.5702(230)	_
LQCD [3] $(m_{\pi} = 411 \text{ MeV})$	_	0.571(18)	_	_	_	_	_	_	_	_
LQCD [4] $(m_{\pi} = 432 \text{ MeV})$	1.8520(875)	0.6129(478)	-0.5949(489)	-1.8108(868)	1.2228(473)	0.0124(244)	-1.1765(450)	0.6059(213)	-0.5885(223)	_
LQCD [3] $(m_{\pi} = 490 \text{ MeV})$	_	0.578(13)	_	_	_	_	_	_	_	_
LQCD [3] $(m_{\pi} = 563 \text{ MeV})$	_	0.5887(98)	_	_	-	-	-	_	-	_
$\chi \mathrm{PT}$ [6]	$2.25^{*}$	_	_	_	_	_	_	_	_	_
RCQM[GBE] [7, 8]	$2.24^{*}$	_	_	_	$1.499^\dagger$	_	_	$0.75^\dagger$	_	_
LCSR [9]	$2.70\pm0.6^*$	_	_	_	_	_	_	_	_	_
PCQM [10]	$1.863^{*}$	-	-	-	$1.242^\dagger$	-	-	$0.621^\dagger$	_	_

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[10] X. Y. Liu et al, PRC 97, 055206 (2018).

### Additional numerical results

$m_{\rm s}=180\;{\rm MeV}$	$\Delta^{++}$	$\Delta^+$	$\Delta^0$	$\Delta^{-}$	$\Sigma^{*+}$	$\Sigma^{*0}$	$\Sigma^{*-}$	$\Xi^{*0}$	Ξ*-	$\Omega^{-}$
$g_3^{(3)B}(0)$	267.7	89.2	-89.2	-267.7	228.0	0	-228.0	141.1	-141.1	0
$g_3^{(0)B}(0)$	24.76	24.76	24.76	24.76	1.795	1.795	1.795	-34.01	-34.01	-83.59
$g_3^{(8)B}(0)$	100.4	100.4	100.4	100.4	-10.29	-10.29	-10.29	-185.8	-185.8	-432.6
$\langle r_A^2 \rangle_B \; [{ m fm}^2]$	0.507	0.507	0.507	0.507	0.498	_	0.498	0.490	0.490	_
$M_A \; [{ m GeV}]$	0.960	0.960	0.960	0.960	0.969	-	0.969	0.976	0.976	-

$$g_1^{(3)B}(Q^2) = \frac{g_1^{(3)B}(0)}{\left(1 + \frac{Q^2}{M_A^2}\right)^2}, \quad \langle r_A^2 \rangle = \frac{12}{M_A^2}$$

### Preparation for the form factor calculation

#### Using parameter

 $m_{\pi} = 139.57 \text{MeV}, \ f_{\pi} = 93 \text{MeV}, \ M_c = 420 \text{MeV}, \ m_s = 180 \text{MeV}$  $D = 8 \text{fm}, \ \Lambda_1 = 0.377 \text{GeV}, \ \Lambda_2 = 1.428 \text{GeV}, \ m_0 = 6.13 \text{MeV}$ 

#### Calculated values

$M_{\rm cl}({\rm GeV})$	$I_1(\mathrm{fm})$	$I_2(\mathrm{fm})$	$K_1(\mathrm{fm})$	$K_2(\mathrm{fm})$	$\Sigma_{\pi N}(\text{MeV})$	$\langle \bar{q}q \rangle ({ m MeV}^3)$
1.2957	1.084	0.519	0.410	0.264	43.89	$-(239.51)^3$