

NuSYM 2018

Gravitational waves and Tidal deformability of Neutron Stars

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in collaboration with

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(Daegu Univ.), and Chang-Hwan Lee (Pusan Nat'l Univ.)

First detection of GW from a BNS

GW170817

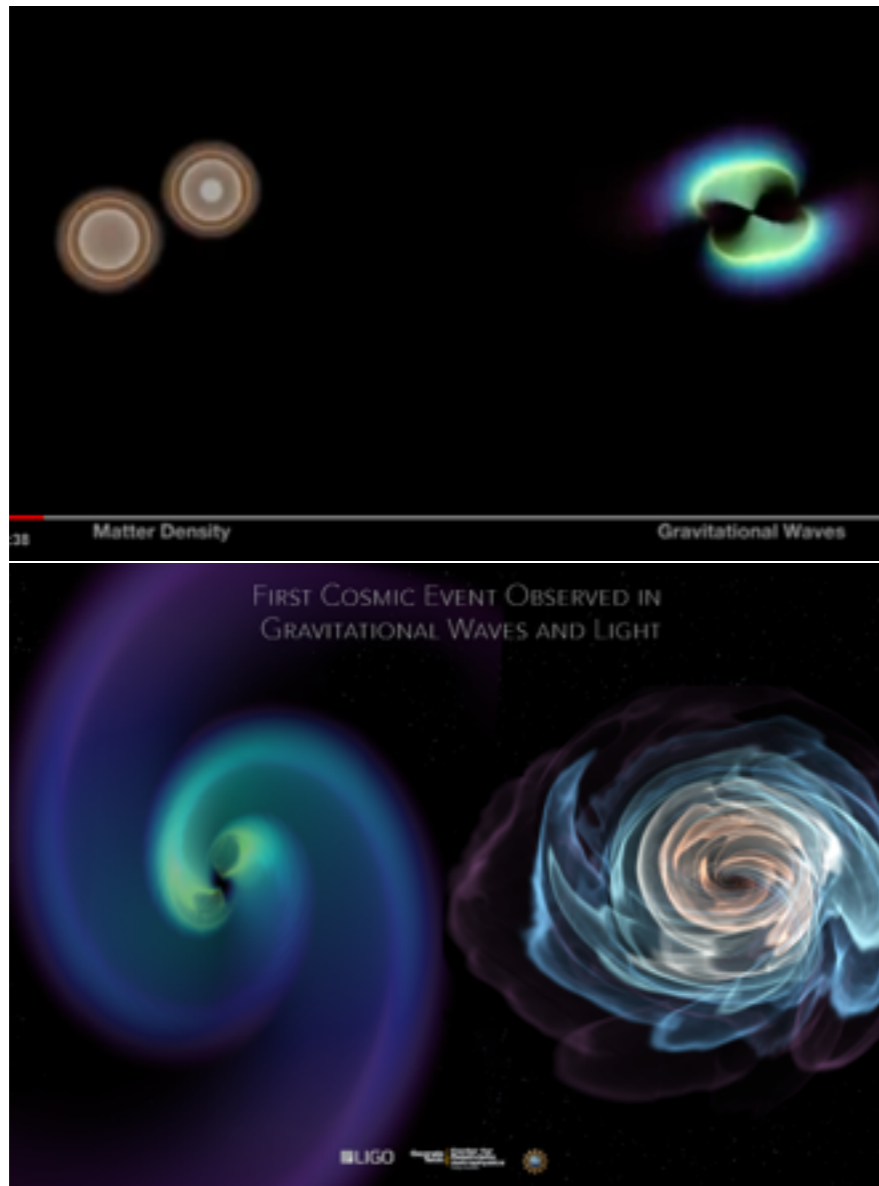
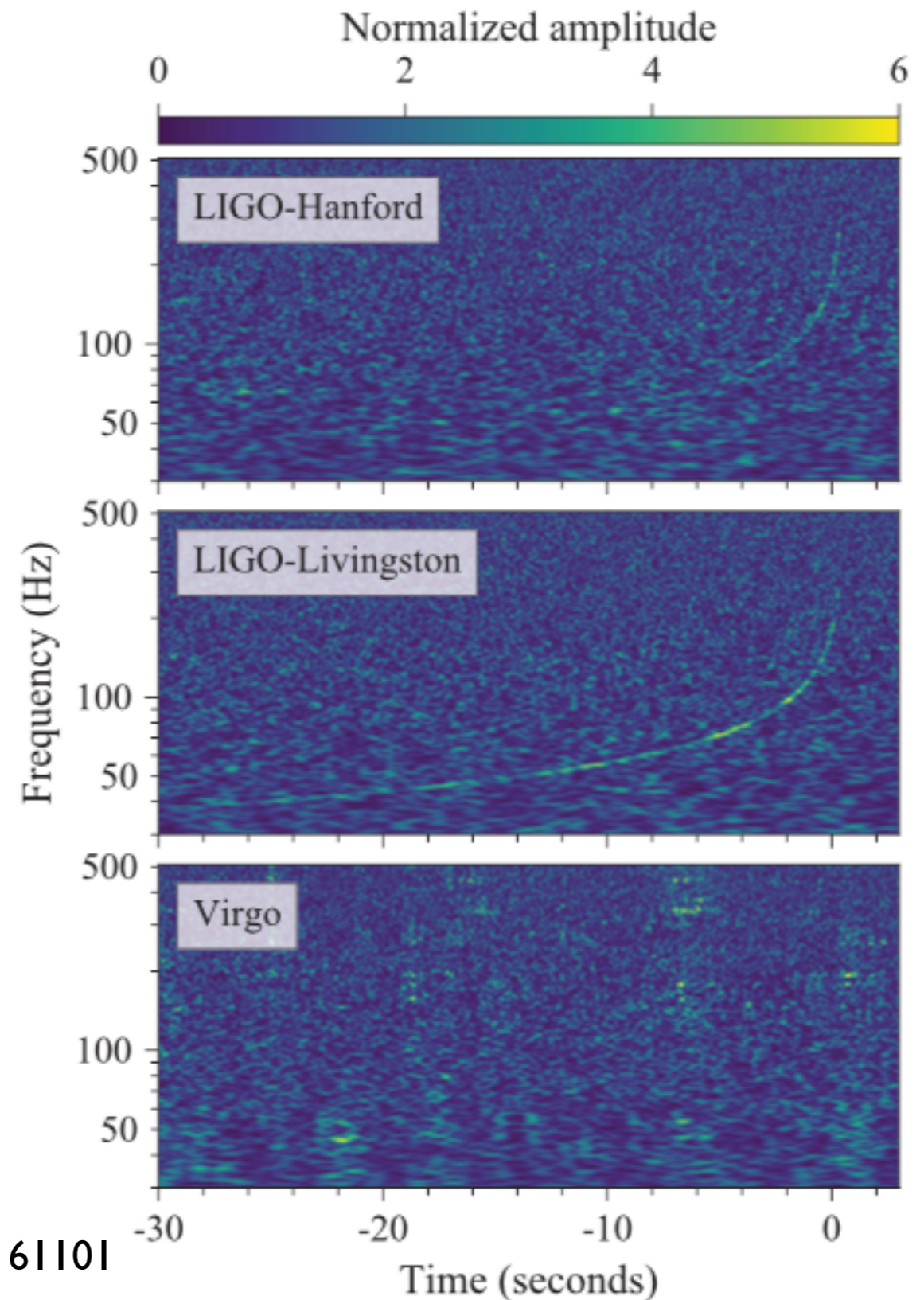
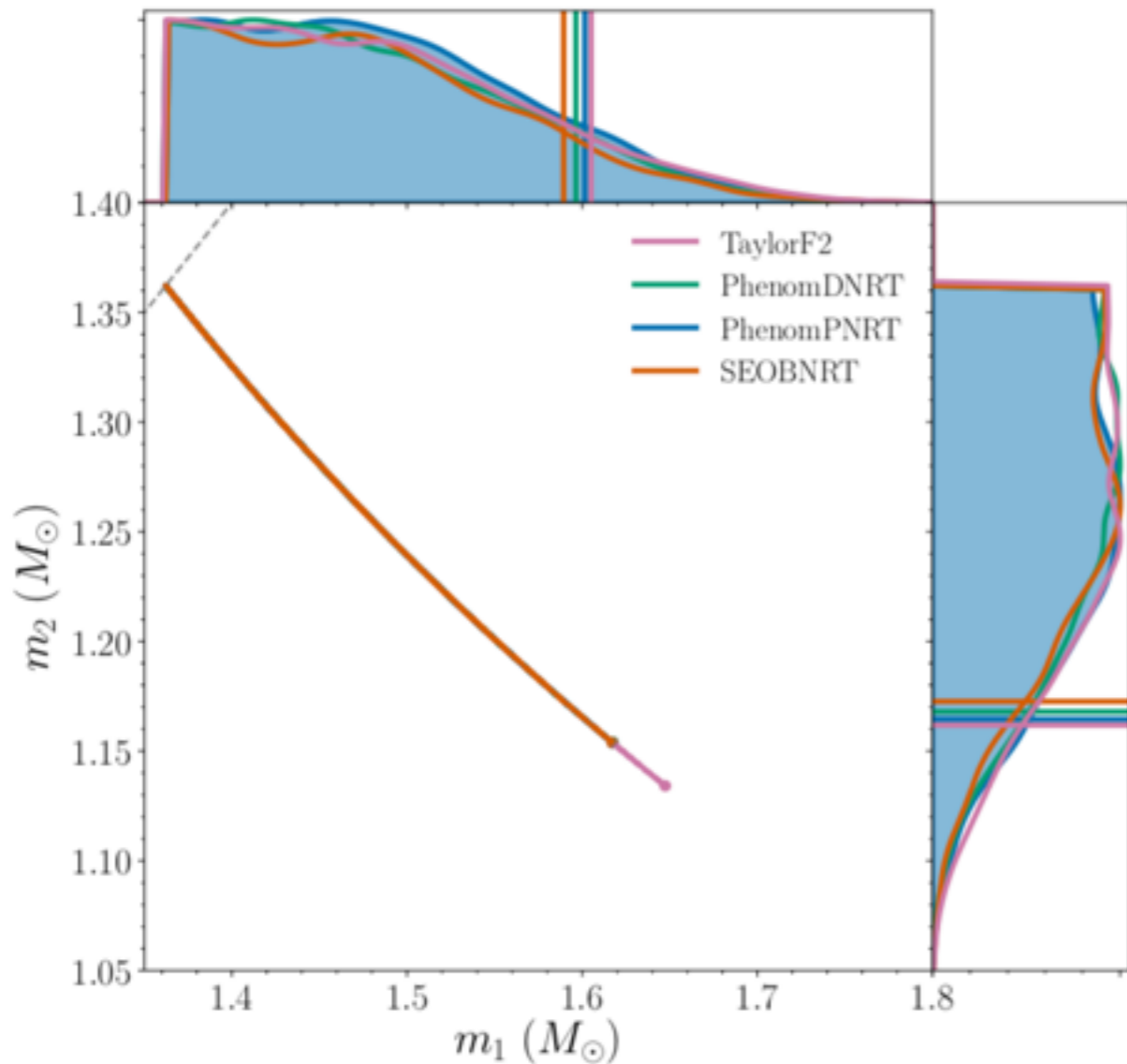


Image credit: Karan Jani/Georgia Tech.

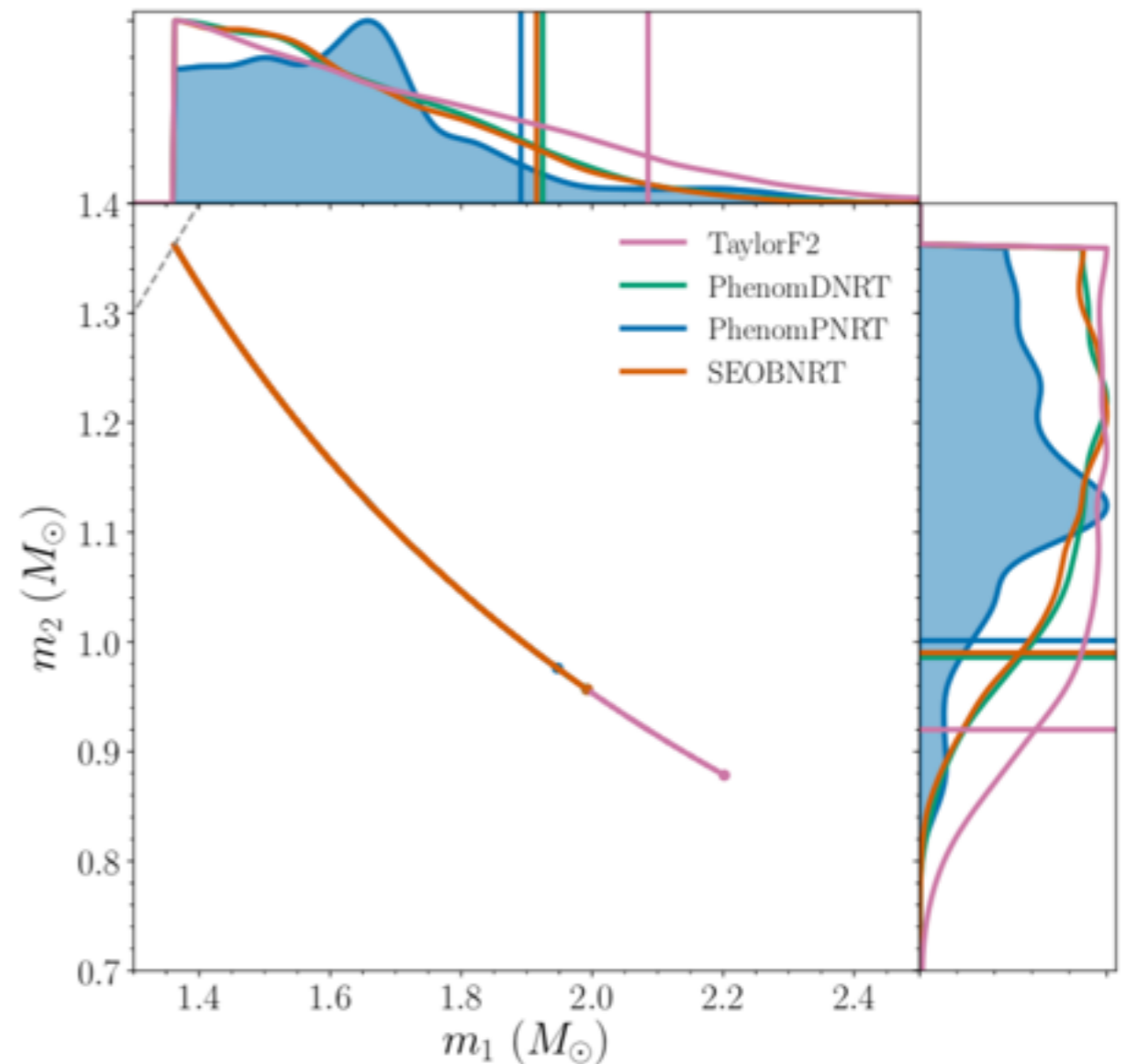


Parameter Estimation (I)

low spin prior:
 $m_1 = 1.36 \sim 1.60 M_\odot$
 $m_2 = 1.16 \sim 1.36 M_\odot$



high spin prior:
 $m_1 = 1.36 \sim 1.89 M_\odot$
 $m_2 = 1.00 \sim 1.36 M_\odot$

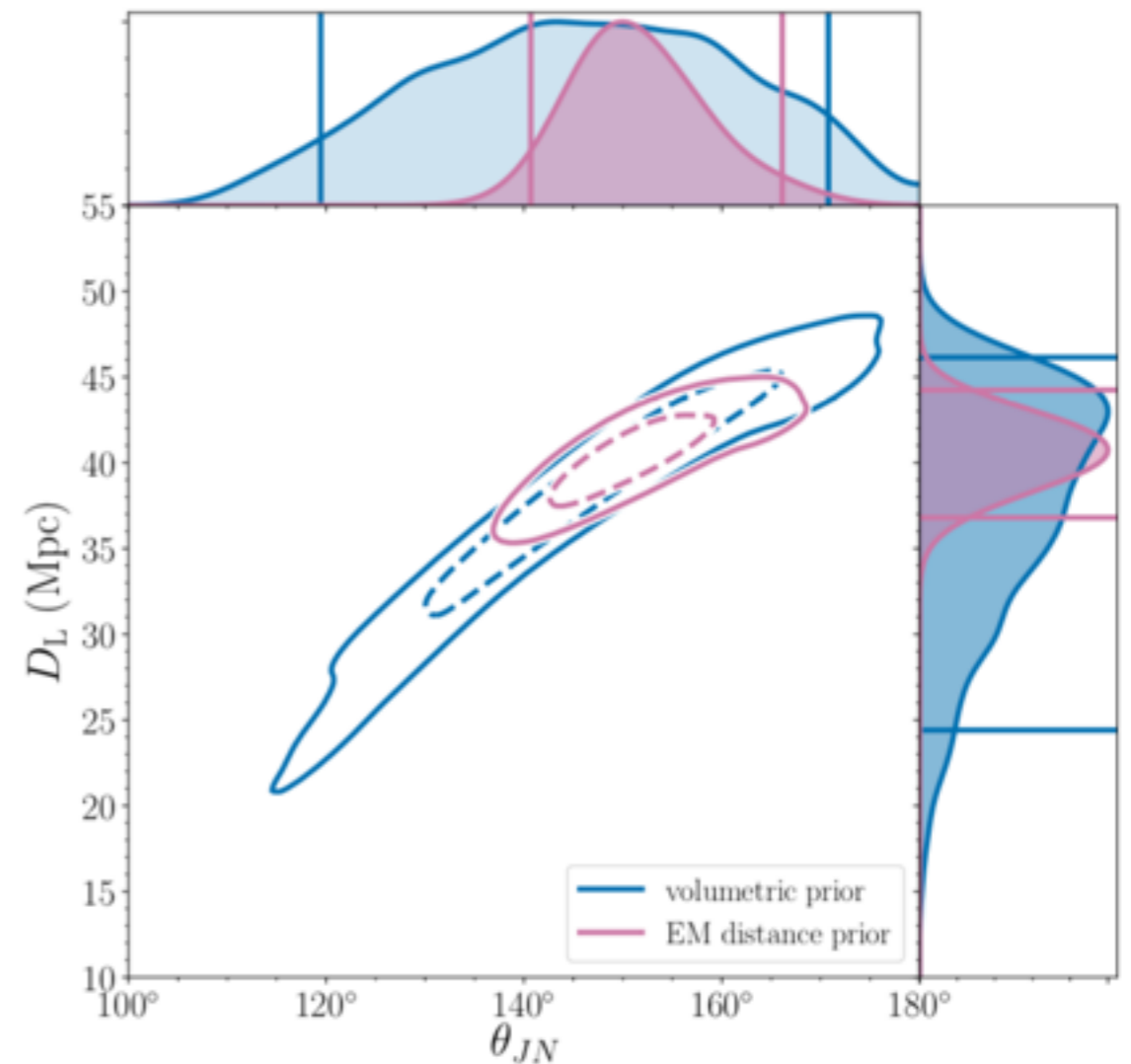
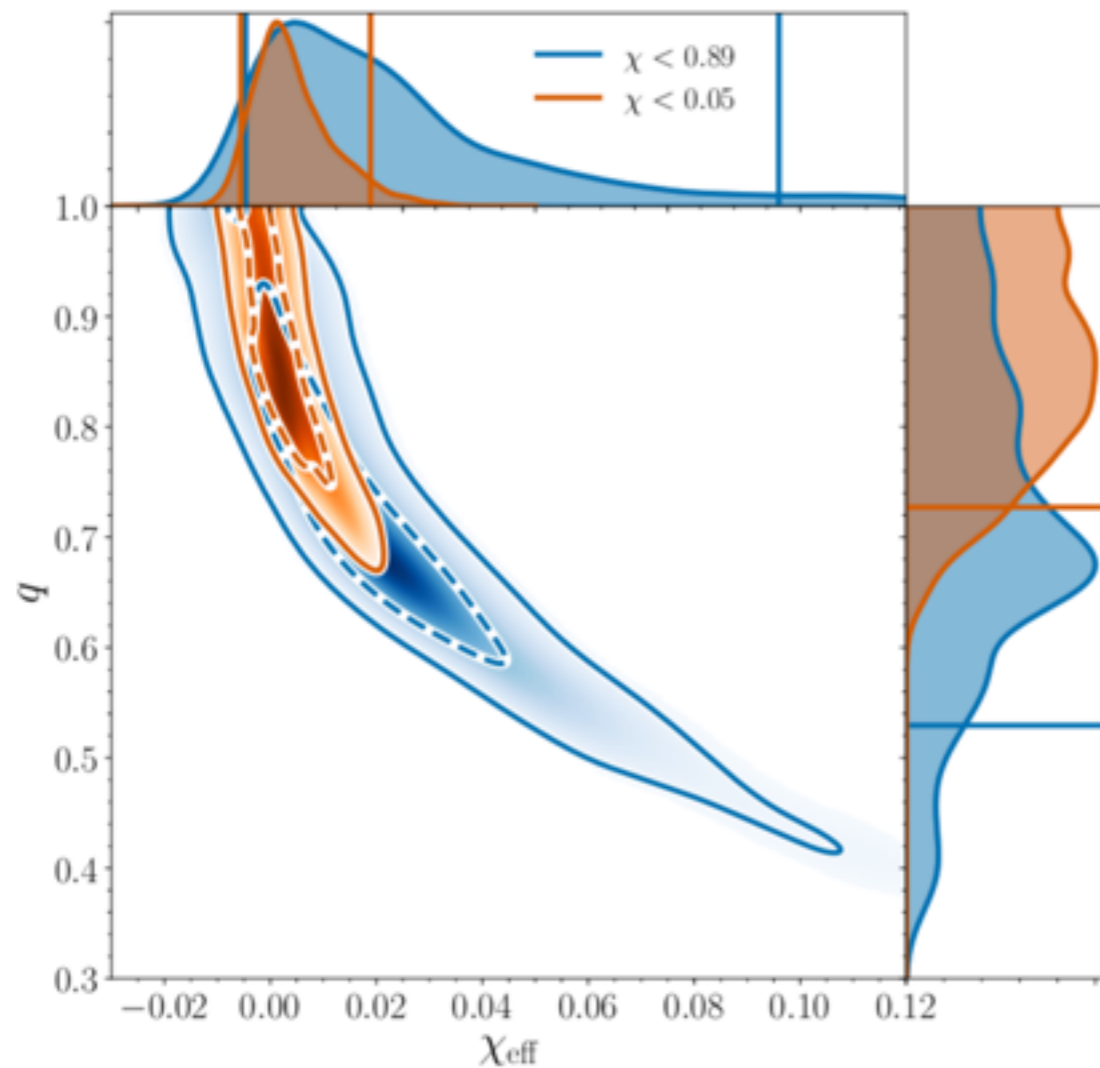


Abbott et al. (LSC and Virgo), arxiv:1805.11579

Parameter Estimation (2)

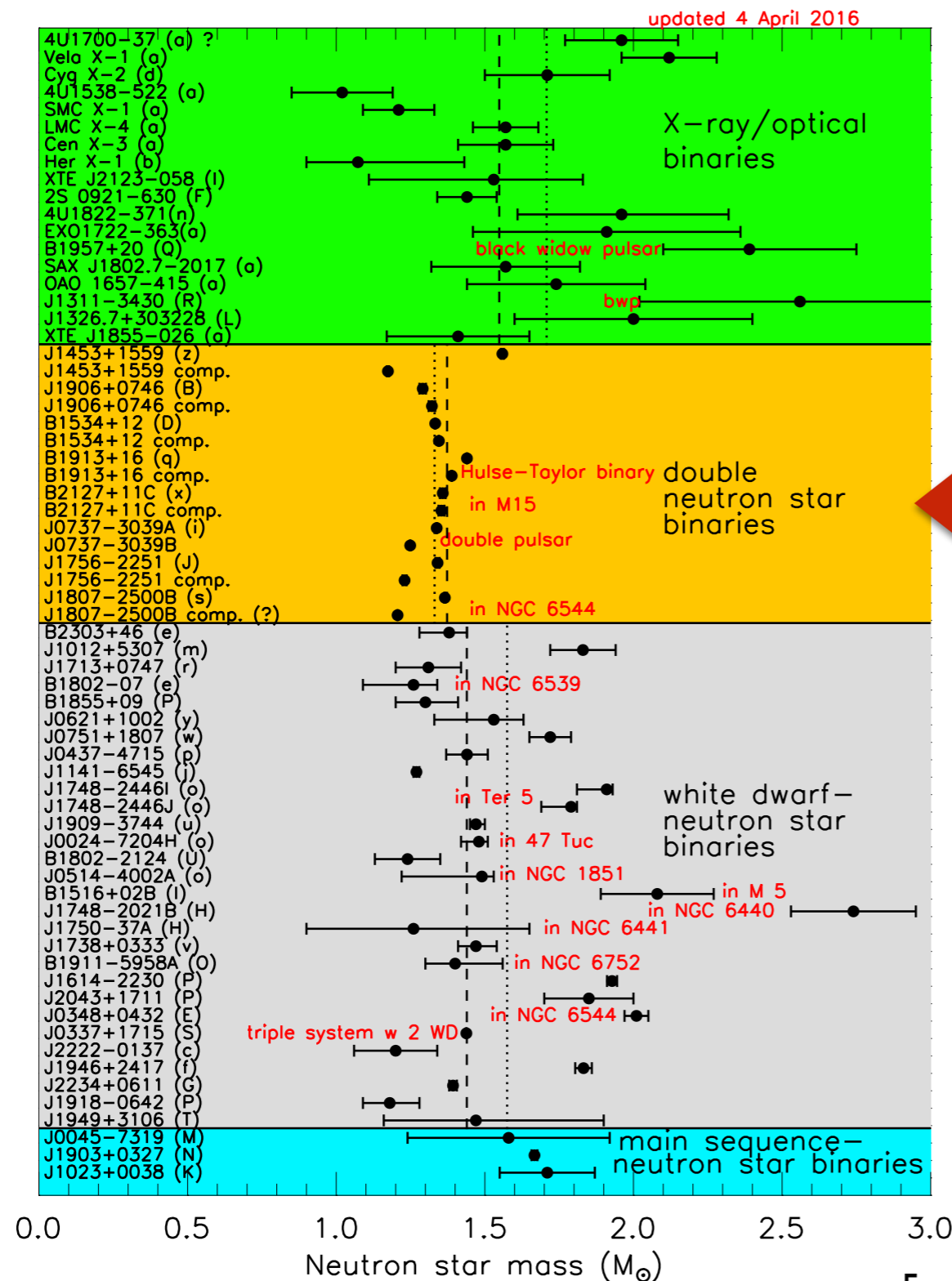
mass ratio $q =$
low spin: $0.73 \sim 1.0$
high spin: $0.53 \sim 1.0$

$D_L \sim 40$ Mpc
(consistent with D of NGC4993)



Abbott et al. (LSC and Virgo), arxiv:1805.11579

Neutron Star of Known Mass

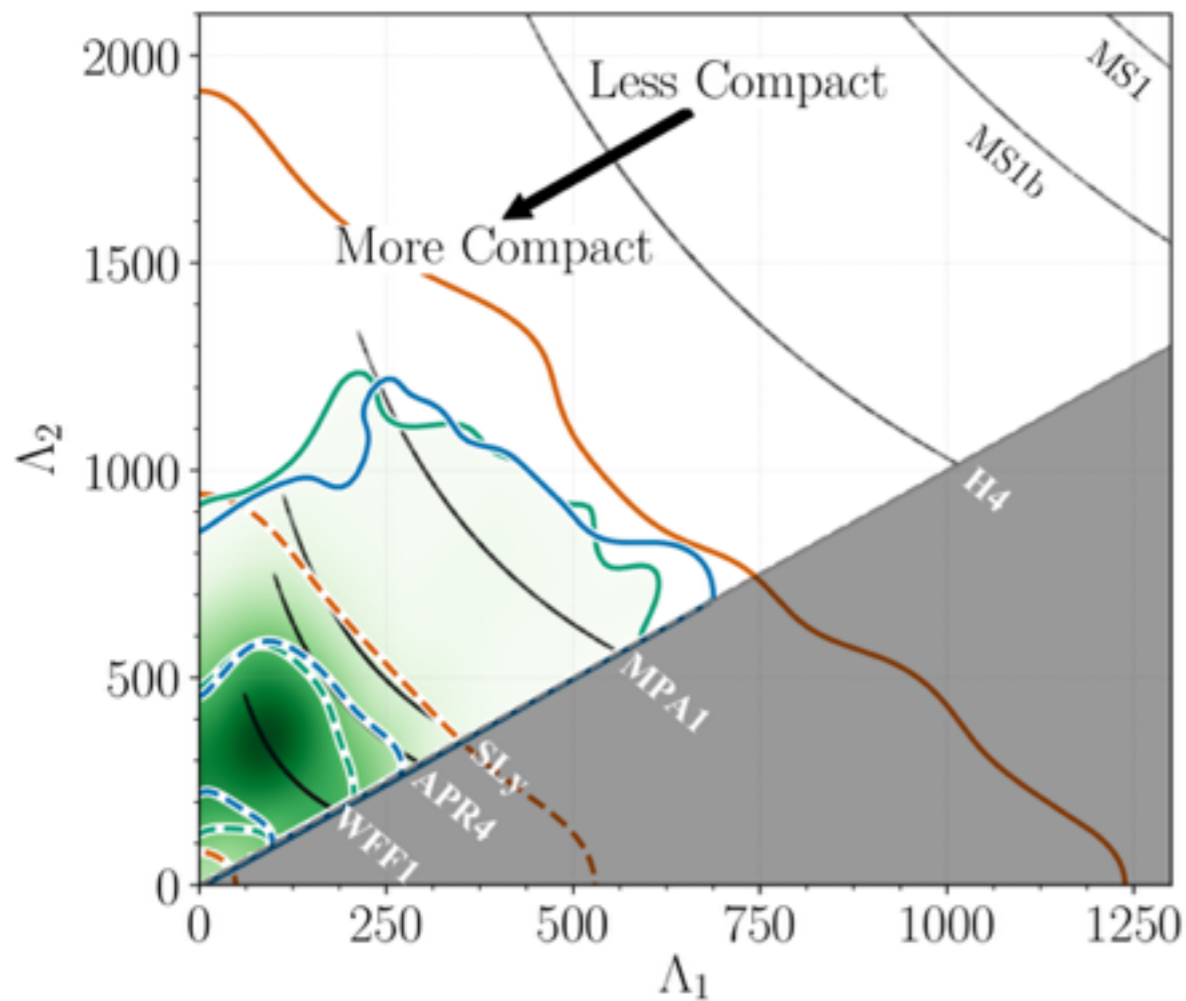


GW170817:
BNS
 $M_1: 1.36 \sim 1.60 M_{\odot}$
 $(1.36 \sim 2.26)$
 $M_2: 1.17 \sim 1.36 M_{\odot}$
 $(0.86 \sim 1.36)$

J. Lattimer, Annu.Rev.Nucl.Part.Sci.62,485(2012)
 and <https://stellarcollapse.org> by C. Ott

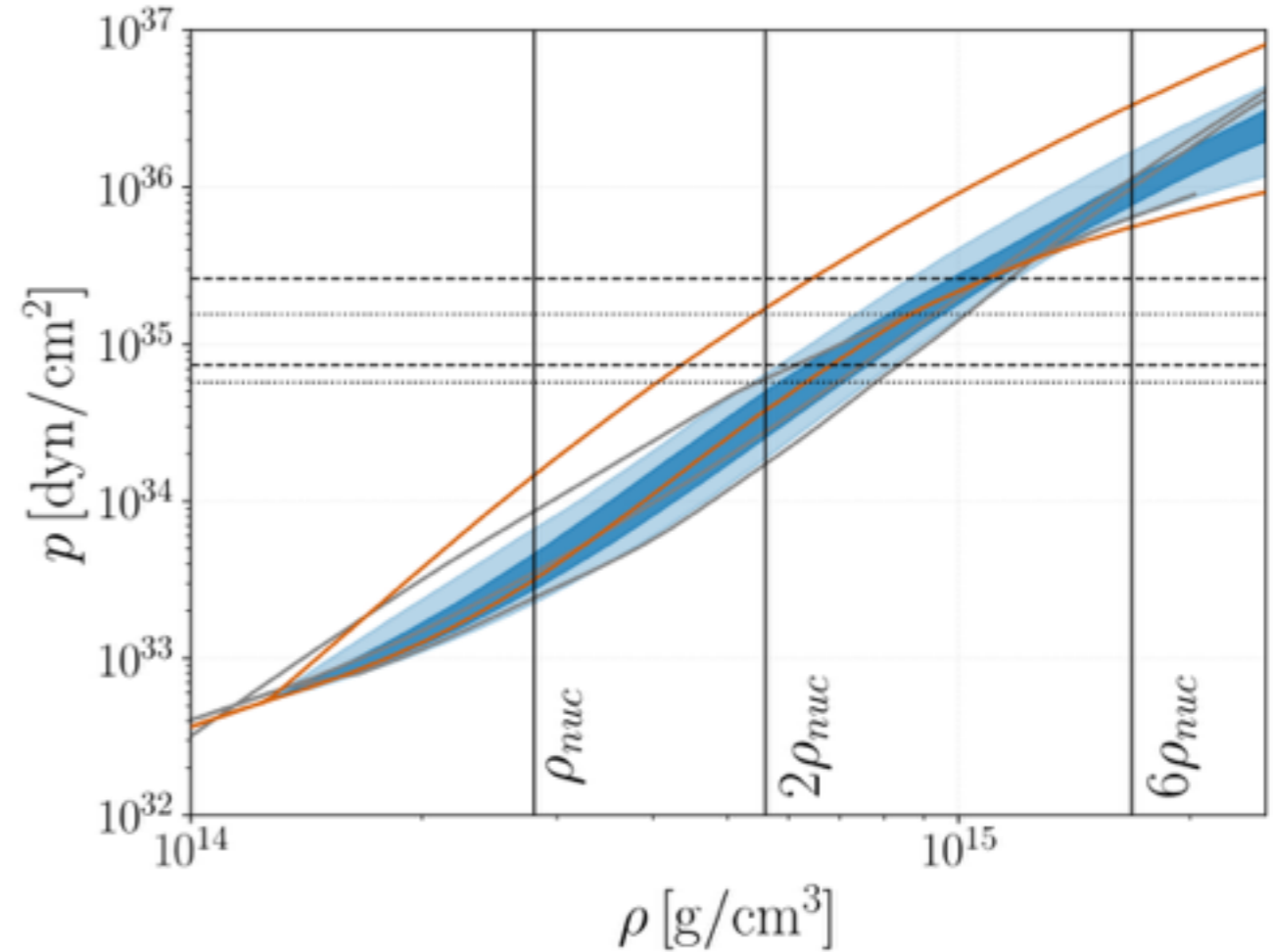
A new constraint by GW Obs. (I)

$$\Lambda(1.4M_{\odot}) = 190^{+390}_{-120}$$



$$P(2 \rho_{\text{nuc}}) = 3.5^{+2.7}_{-1.7} \times 10^{34} \text{ dyne/cm}^2$$

$$P(6 \rho_{\text{nuc}}) = 9.0^{+7.9}_{-2.6} \times 10^{35} \text{ dyne/cm}^2$$



Abbott et al. (LSC and Virgo), arxiv:1805.11581 (PRL accepted)

$$\rho_{\text{nuc}} = 2.8 \times 10^{14} \text{ g/cm}^3$$

A new constraint by GW Obs. (2)

EoS insensitive relations (Yagi&Yunes,PR2017)

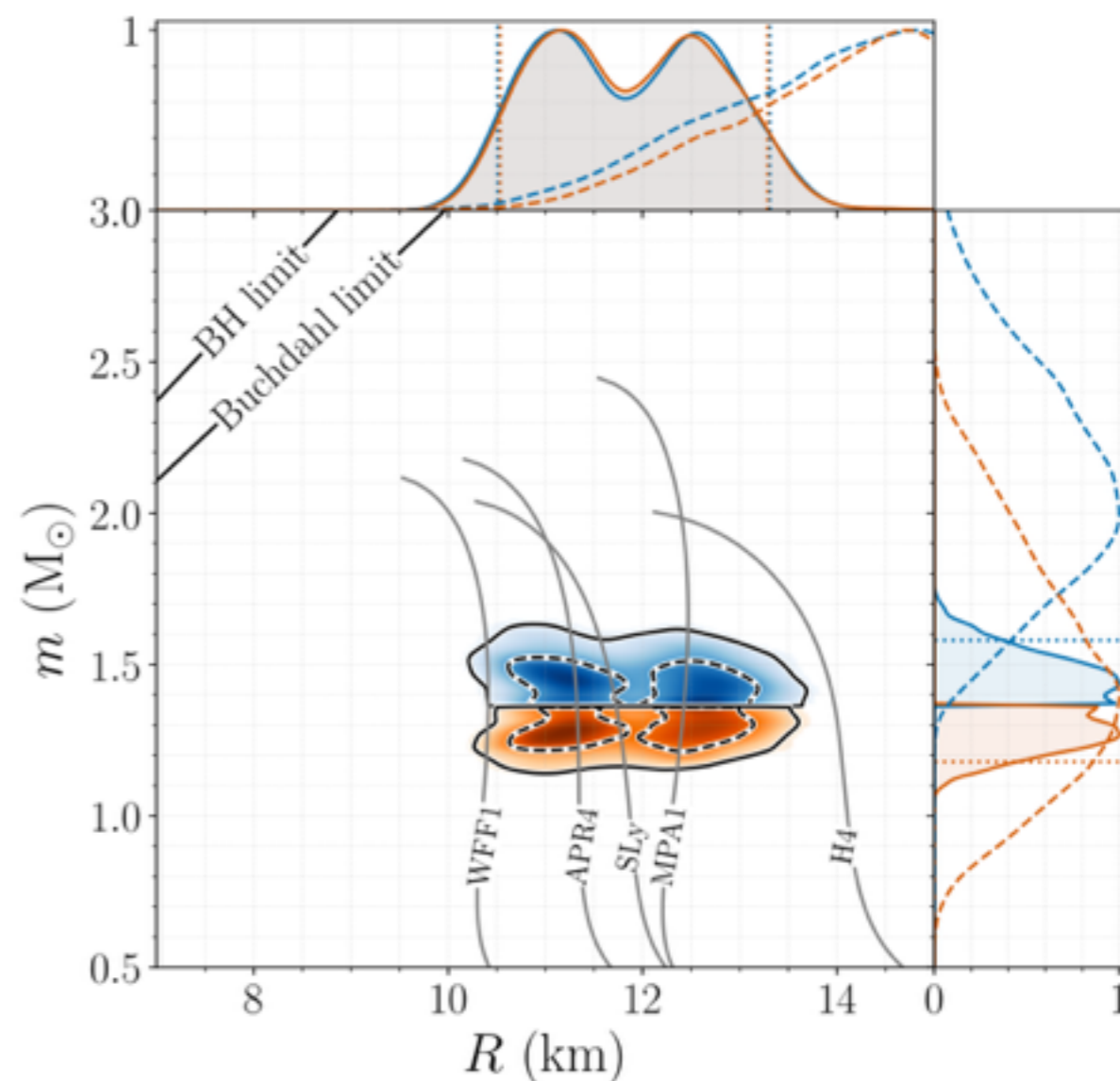
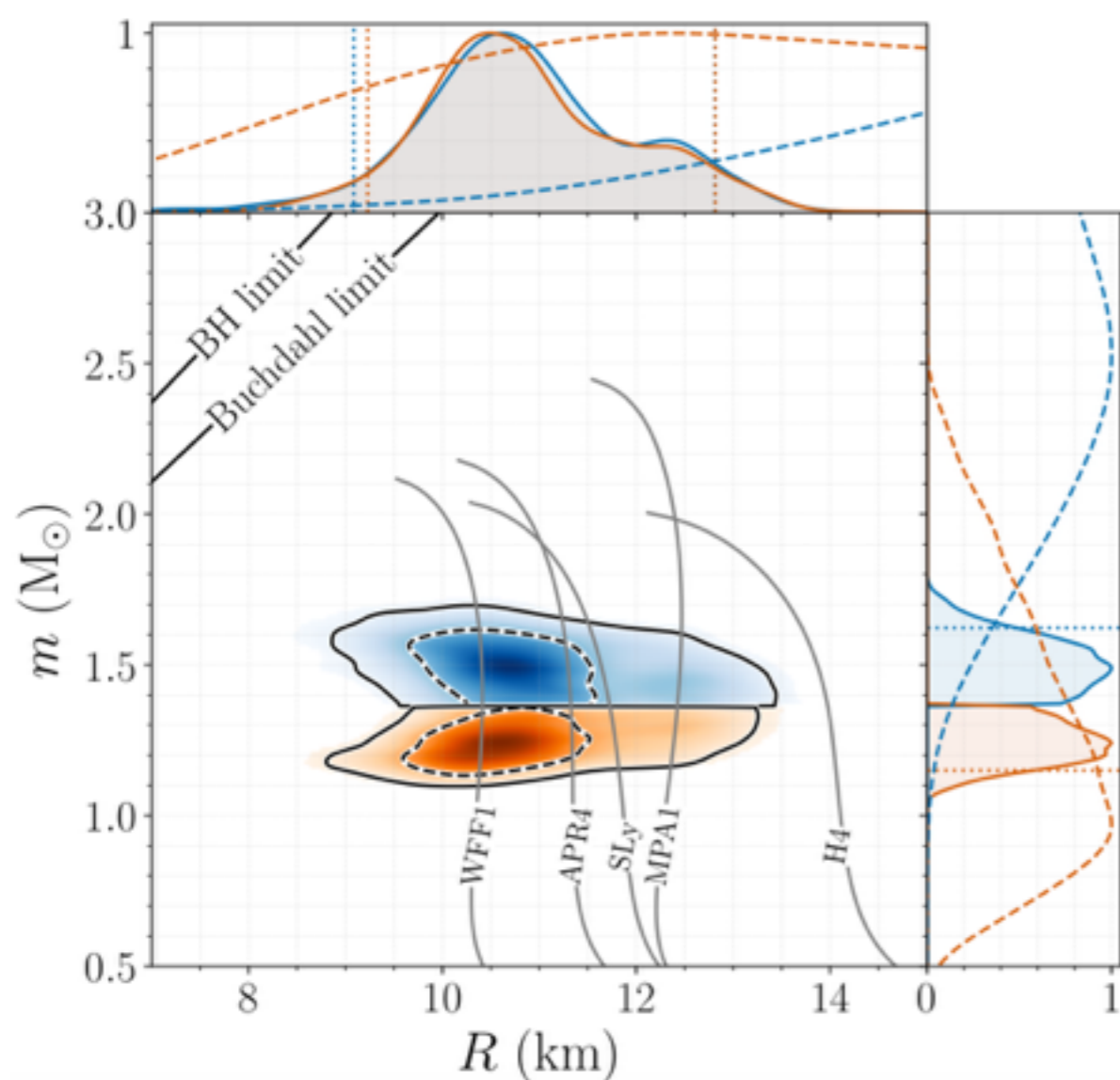
$$R_1 = 10.8^{+2.0}_{-1.7} \text{ km}$$

$$R_2 = 10.7^{+2.1}_{-1.5} \text{ km}$$

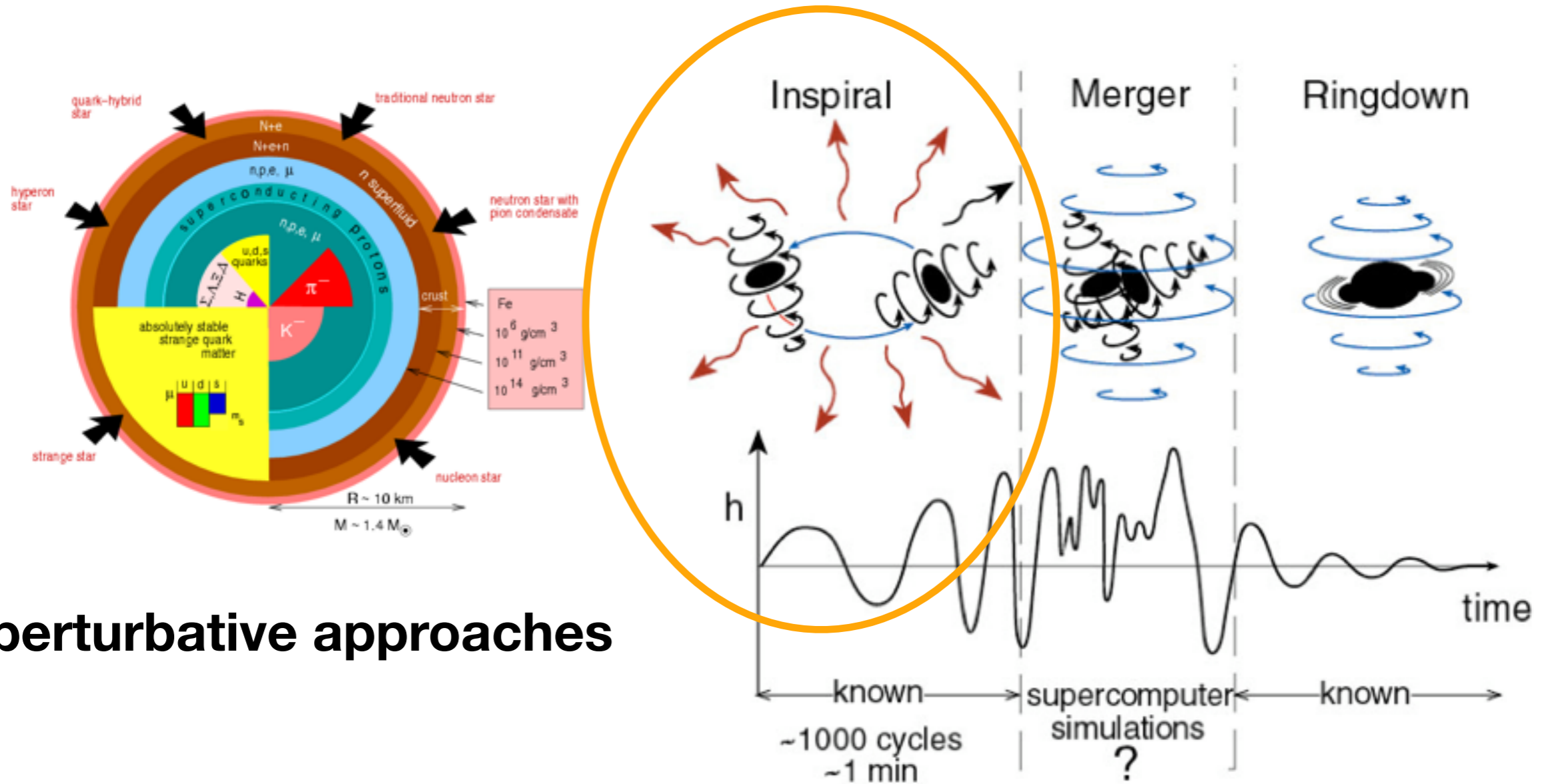
Parametrized EoS: $M_{\text{max}} \geq 1.97 M_{\odot}$

$$R_1 = 11.9^{+1.4}_{-1.4} \text{ km}$$

$$R_2 = 11.9^{+1.4}_{-1.4} \text{ km}$$



Response of NS to GW during inspiral



perturbative approaches

$$Q_{ij} = -\lambda \mathcal{E}_{ij}$$

$$\lambda = \frac{2}{3} \frac{R^5}{G} k_2$$

λ : Tidal deformability
 k_2 : Tidal Love number

GW waveform in Frequency Domain

$$\tilde{h}_T(f) = \mathcal{A} f^{-7/6} e^{i\Psi_T(f)}$$

M. Favata, PRL 112.101101 (2014)

$$\Psi_T(f) = \varphi_c + 2\pi f t_c + \frac{3}{128\eta v^5} (\Delta\Psi_{3.5\text{PN}}^{\text{pp}} + \Delta\Psi_{3\text{PN}}^{\text{spin}} + \Delta\Psi_{2\text{PN}}^{\text{ecc.}} + \Delta\Psi_{6\text{PN}}^{\text{tidal}} + \Delta\Psi_{6\text{PN}}^{\text{tm}}), \quad (1)$$

$$\Delta\Psi_{6\text{PN}}^{\text{tidal}} = -\frac{39}{2} \tilde{\Lambda} v^{10} + v^{12} \left(\frac{6595}{364} \delta\tilde{\Lambda} - \frac{3115}{64} \tilde{\Lambda} \right), \quad (4)$$

$$v = (\pi f M)^{1/3}$$

$$\tilde{\Lambda} = \frac{16}{13} \frac{(m_1 + 12m_2)m_1^4\Lambda_1 + (m_2 + 12m_1)m_2^4\Lambda_2}{(m_1 + m_2)^5}$$

$$\Lambda = \lambda/M^5 \rightarrow G \left(\frac{c^2}{GM} \right)^5 \lambda = \frac{2}{3} \left(\frac{Rc^2}{GM} \right)^5 k_2$$

Tidal Love number, k2

$$\frac{dH}{dr} = \beta \quad \frac{d\beta}{dr} = 2 \left(1 - 2\frac{M}{r}\right)^{-1} \times H \left\{ -2\pi [5\epsilon + 9P + (d\epsilon/dP)(\epsilon + P)] \right. \\ \left. + \frac{3}{r^2} + 2 \left(1 - 2\frac{M}{r}\right)^{-1} \left(\frac{M}{r^2} + 4\pi r P\right)^2 \right\} \\ + \frac{2\beta}{r} \left(1 - 2\frac{M}{r}\right)^{-1} \left\{ -1 + \frac{M}{r} + 2\pi r^2(\epsilon - P) \right\}$$

TOV

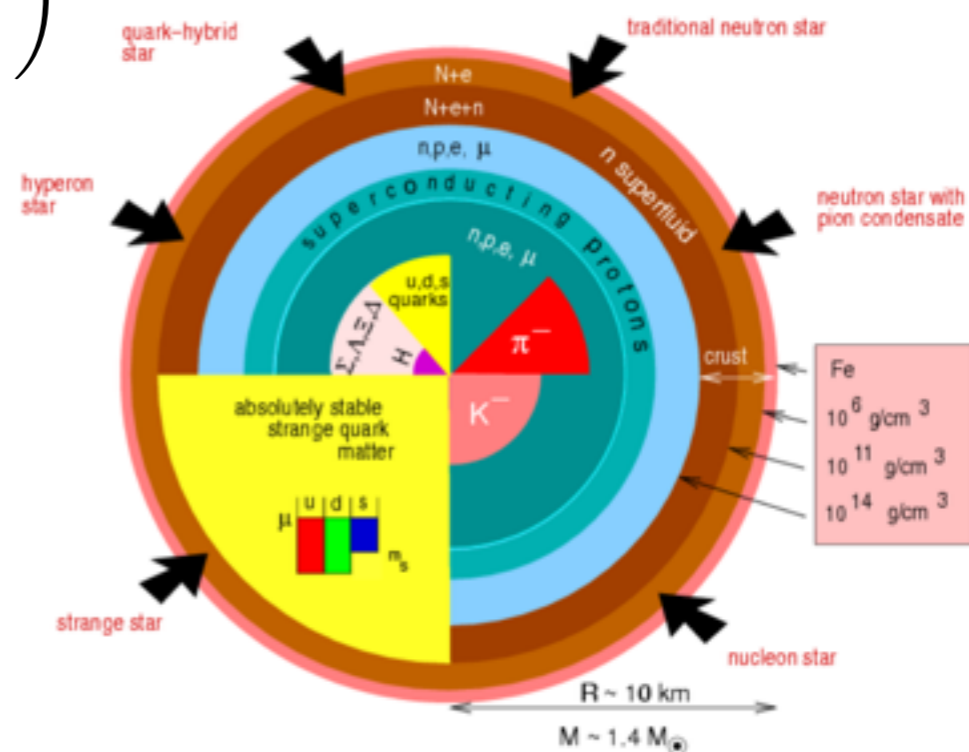
$$\frac{dP}{dr} = -\frac{GM\rho}{r^2} \left(1 + \frac{P}{\rho c^2}\right) \left(1 + \frac{4\pi P r^3}{M c^2}\right) \left(1 - \frac{2GM}{rc^2}\right)$$

$$\frac{dM}{dr} = 4\pi r^2 \left(\frac{\epsilon}{c^2}\right)$$

k2 (λ, Λ) depends on NS EoS !!

Compositions of a NS

F. Weber 2005



TOV Eq. vs. Diff. Eq. for Tidal deformability

a spherical symmetric star in hydrostatic equilibrium

$$G_{\mu\nu} = 8\pi T_{\mu\nu}$$



$$\begin{aligned} ds_0^2 &= g_{\alpha\beta}^{(0)} dx^\alpha dx^\beta \\ &= -e^{\nu(r)} dt^2 + e^{\lambda(r)} dr^2 + r^2 (d\theta^2 + \sin^2\theta d\phi^2). \end{aligned}$$

$$T_{\alpha\beta} = (\rho + p)u_\alpha u_\beta + pg_{\alpha\beta}^{(0)},$$

TOV eq.

$$\frac{dP}{dr} = -\frac{GM\rho}{r^2} \left(1 + \frac{P}{\rho c^2}\right) \left(1 + \frac{4\pi Pr^3}{Mc^2}\right) \left(1 - \frac{2GM}{rc^2}\right)$$

$$\frac{dM}{dr} = 4\pi r^2 \left(\frac{\epsilon}{c^2}\right)$$



Mass & Radius

TOV Eq. vs. Diff. Eq. for Tidal deformability

static linearized perturbations due to an external tidal field

$$\delta G_{\mu\nu} = 8\pi\delta T_{\mu\nu}$$



$$g_{\alpha\beta} = g_{\alpha\beta}^{(0)} + h_{\alpha\beta},$$

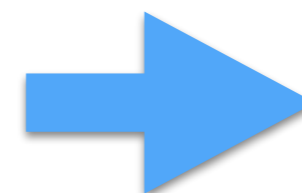
T. Hinderer (2008), K. Thorne and A. Campolattaro (1967)

$$h_{\alpha\beta} =$$

$$\text{diag}[-e^{\nu(r)}H_0(r), e^{\lambda(r)}H_2(r), r^2K(r), r^2\sin^2\theta K(r)]Y_{2m}(\theta, \varphi).$$

$$\delta T_0^0 = -\delta\rho = -(dp/d\rho)^{-1}\delta p \quad \delta T_i^i = \delta p$$

$$H'' + H' \left\{ \frac{2}{r} + e^\lambda \left[\frac{2m(r)}{r^2} + 4\pi r(p - \rho) \right] \right\} + H \left[-\frac{6e^\lambda}{r^2} + 4\pi e^\lambda \left(5\rho + 9p + \frac{\rho + p}{dp/d\rho} \right) - \nu'^2 \right] = 0,$$



k2 or λ

Electric-type tidal coefficients

Damour and Nagar , PRD 80, 084035 (2009)

Electric-type tidal coefficients

$$C_1 = \frac{2}{r} + \frac{1}{2}(\nu' - \lambda') = \frac{2}{r} + e^\lambda \left[\frac{2m}{r^2} + 4\pi r(p - e) \right], \quad (28)$$

$$G\mu_\ell = \frac{a_\ell}{(2\ell - 1)!!} \left(\frac{GM}{c_0^2} \right)^{2\ell+1} = \frac{2k_\ell}{(2\ell - 1)!!} R^{2\ell+1}. \quad (48)$$

$$\begin{aligned} C_0 &= e^\lambda \left[-\frac{\ell(\ell + 1)}{r^2} + 4\pi(e + p) \frac{de}{dp} + 4\pi(e + p) \right] \\ &\quad + \nu'' + (\nu')^2 + \frac{1}{2r}(2 - r\nu')(3\nu' + \lambda') \\ &= e^\lambda \left[-\frac{\ell(\ell + 1)}{r^2} + 4\pi(e + p) \frac{de}{dp} + 4\pi(5e + 9p) \right] \\ &\quad - (\nu')^2, \end{aligned} \quad (29)$$

$$H'' + C_1 H' + C_0 H = 0.$$

we calculated the case, $\ell=2$

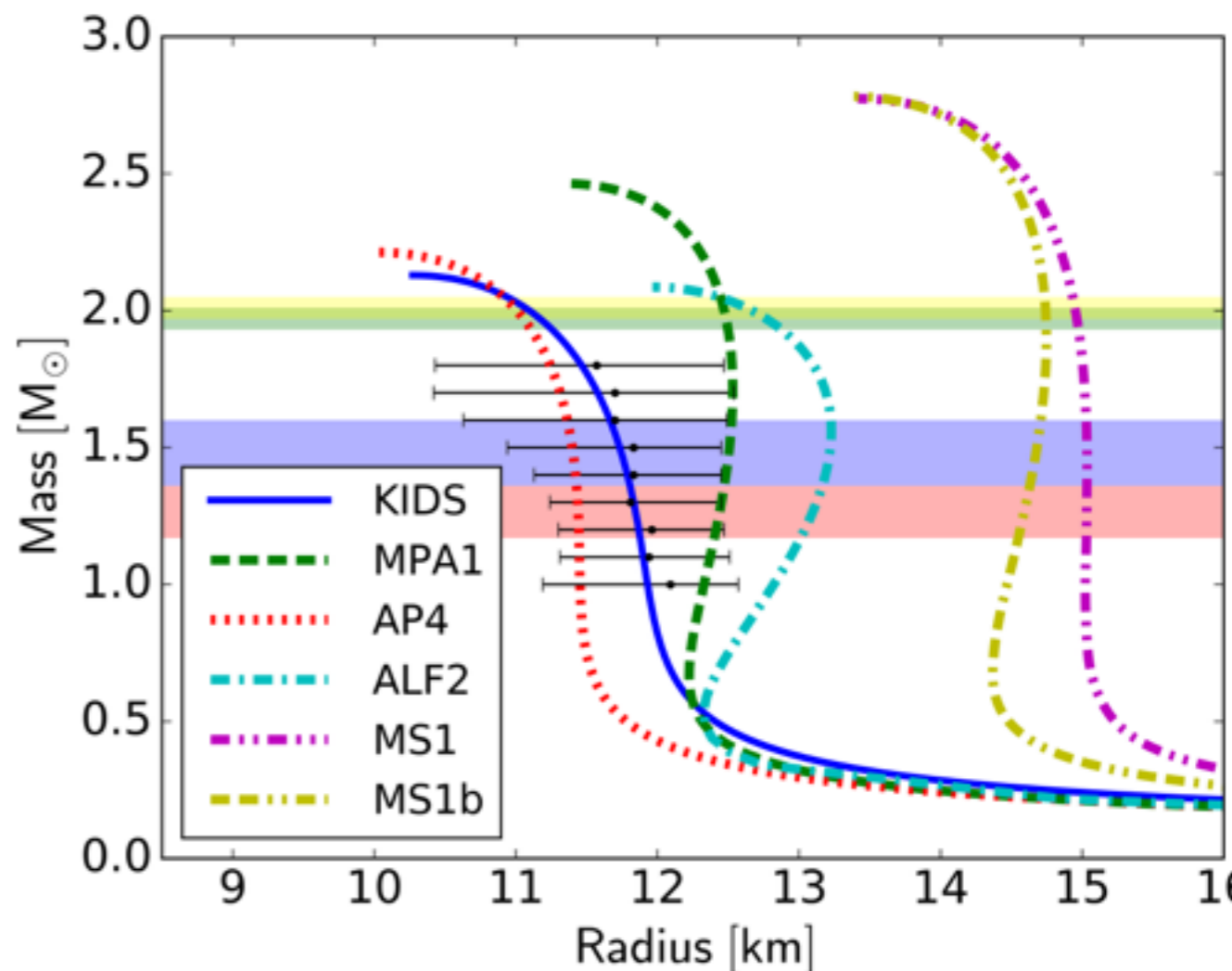
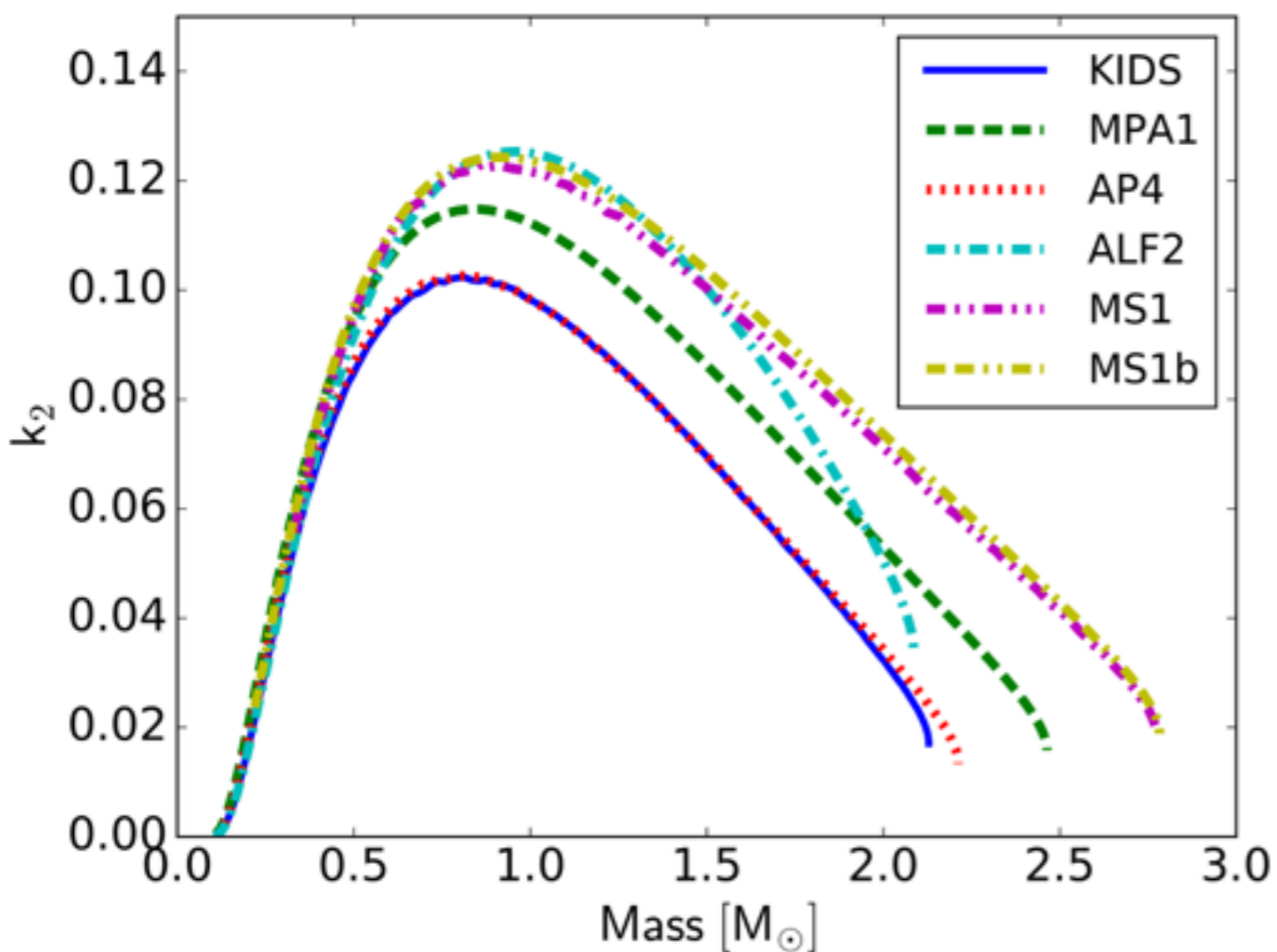
$$\begin{aligned} k_2 &= \frac{8}{5}(1 - 2c)^2 c^5 [2c(y - 1) - y + 2] \left\{ 2c(4(y + 1)c^4 + (6y - 4)c^3 + (26 - 22y)c^2 + 3(5y - 8)c - 3y + 6) \right. \\ &\quad \left. - 3(1 - 2c)^2 (2c(y - 1) - y + 2) \log\left(\frac{1}{1 - 2c}\right) \right\}^{-1}, \end{aligned} \quad (50)$$

$$y_R = \frac{rH'(r)}{H(r)} \Big|_{r=R}$$

$$\begin{aligned} k_3 &= \frac{8}{7}(1 - 2c)^2 c^7 [2(y - 1)c^2 - 3(y - 2)c + y - 3] \left\{ 2c[4(y + 1)c^5 + 2(9y - 2)c^4 - 20(7y - 9)c^3 + 5(37y - 72)c^2 \right. \\ &\quad \left. - 45(2y - 5)c + 15(y - 3)] - 15(1 - 2c)^2 (2(y - 1)c^2 - 3(y - 2)c + y - 3) \log\left(\frac{1}{1 - 2c}\right) \right\}^{-1}, \end{aligned} \quad (51)$$

Tidal deformability of a NS

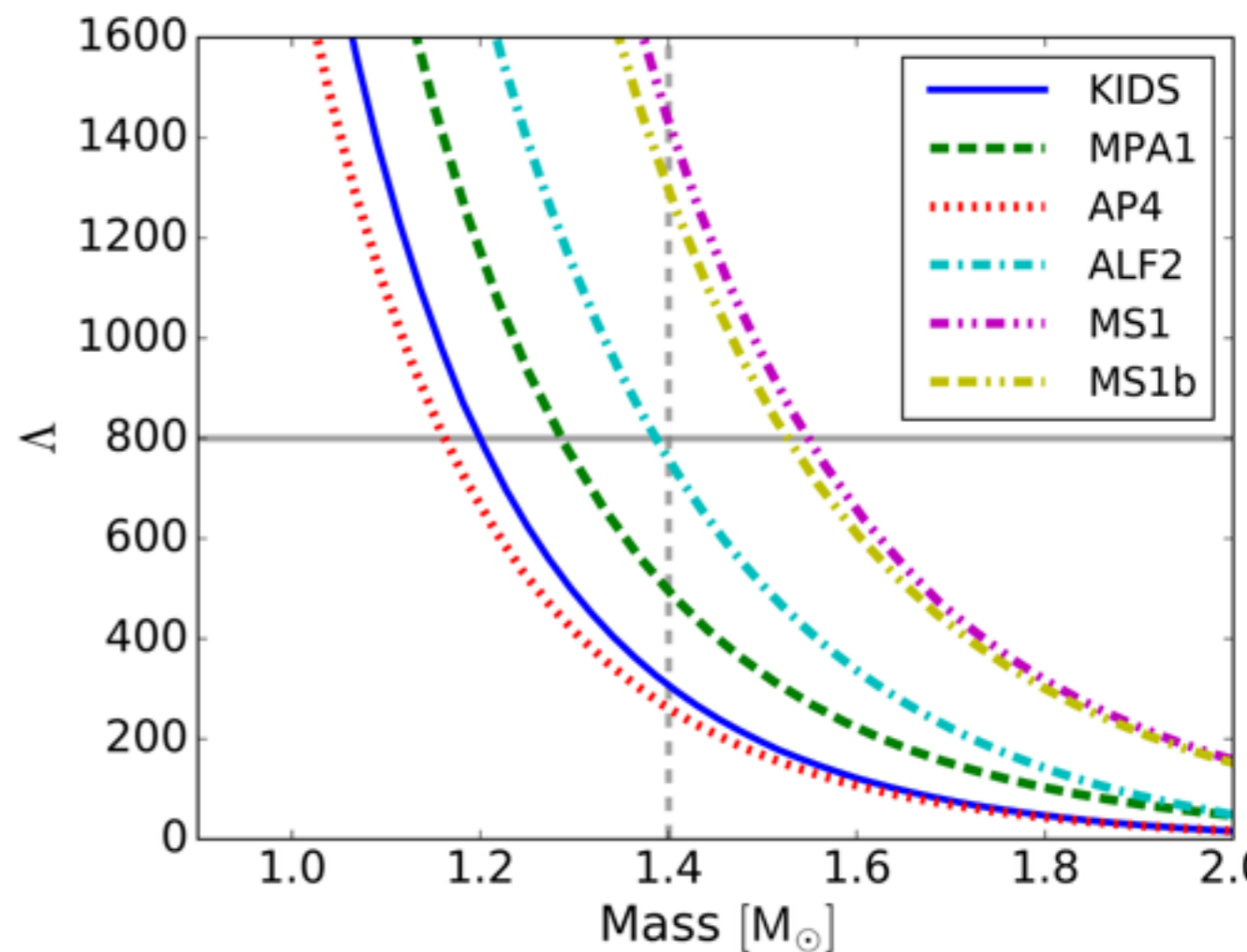
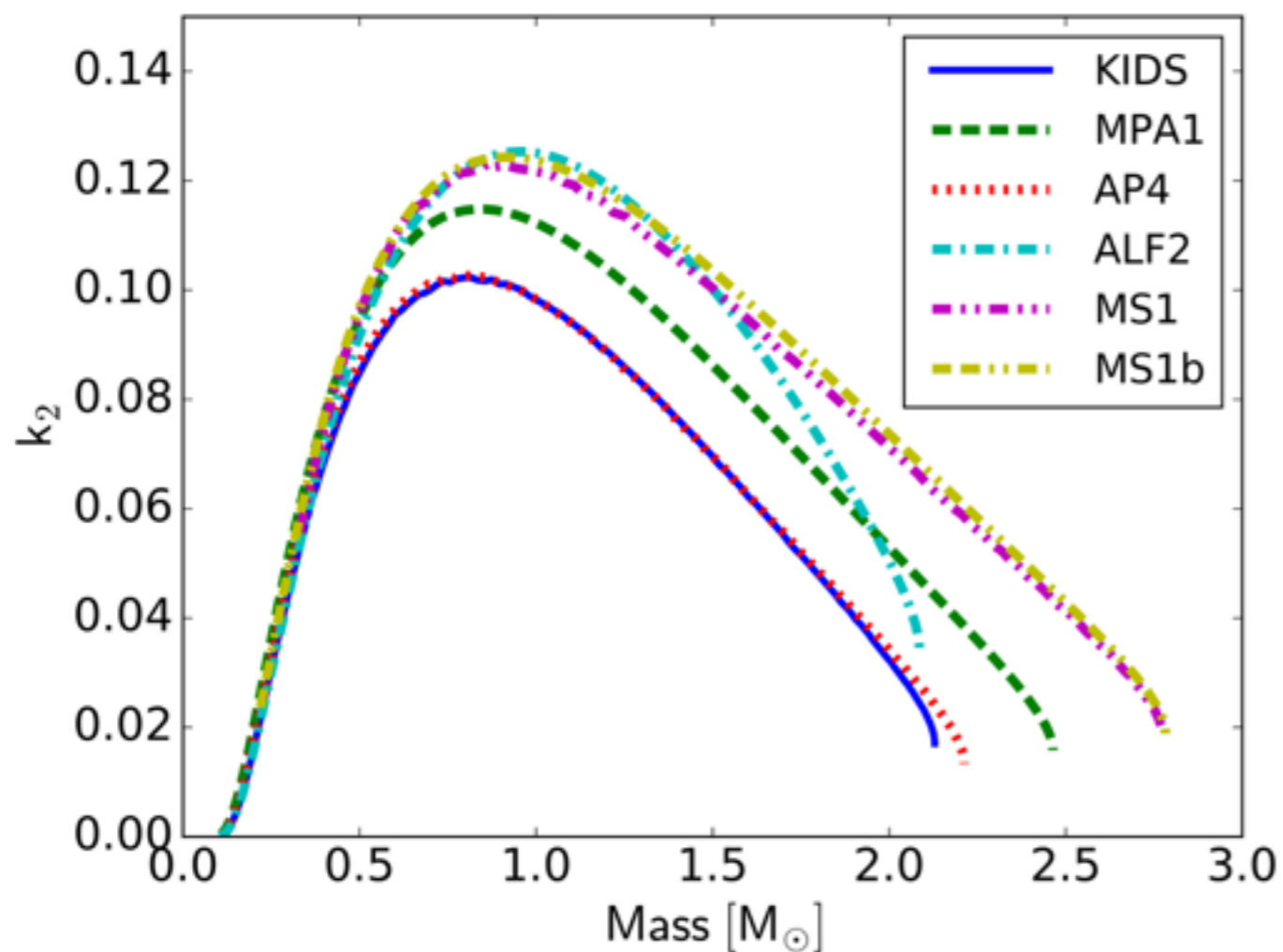
Kim et al., New Physics: Sae Mulli (2018)



GW170817 - Abbott et al. (LSC and Virgo), arxiv:1805.11579
 - Mchirp = 1.188 M_\odot
 - low spin prior : $\Lambda = 300^{+500}_{-190}$ (symmetric) / 300^{+420}_{-230} (HPD)
 - high spin prior : $\Lambda = 0 \sim 630$

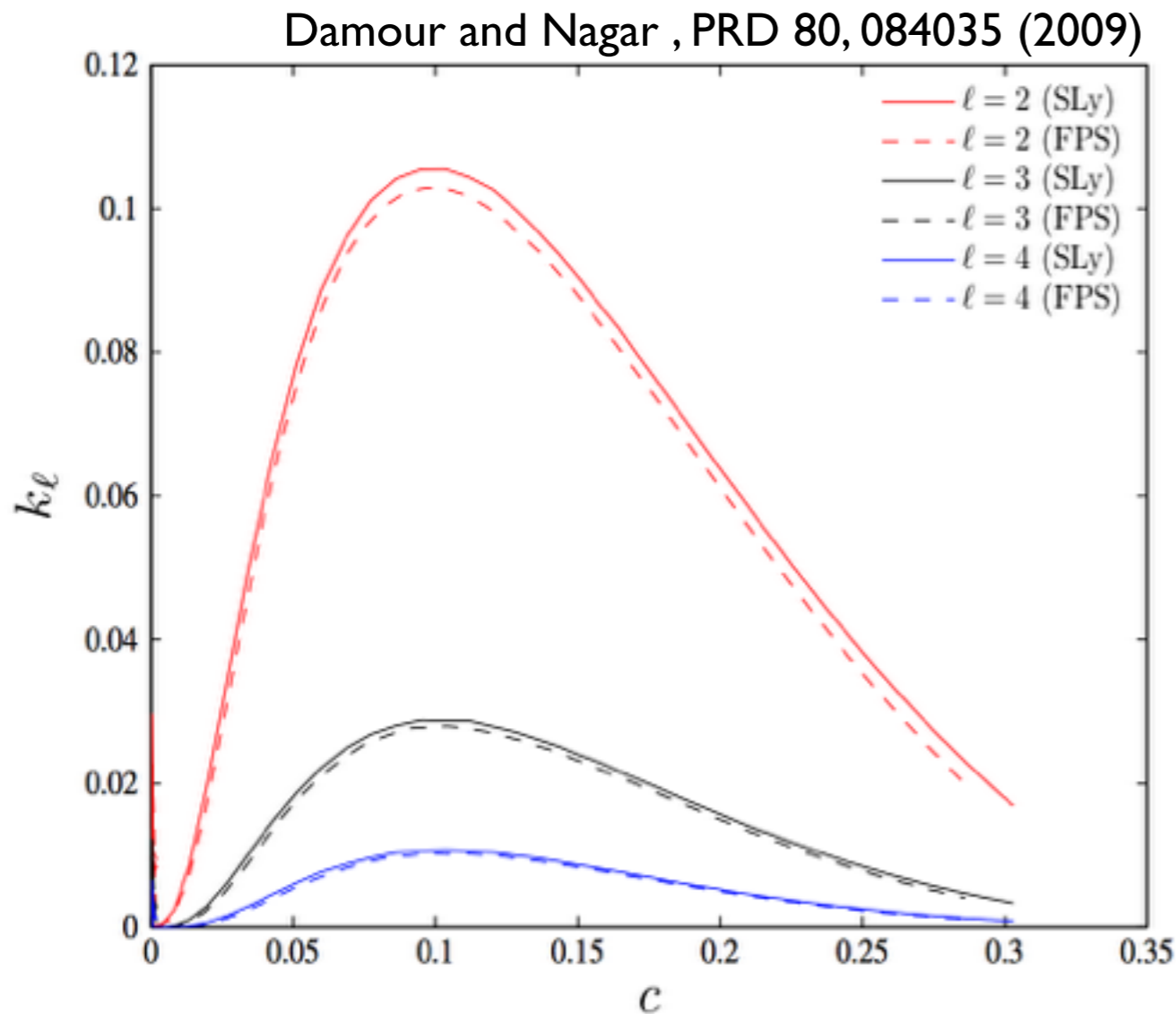
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Higher Tidal coefficients



$$x \equiv (M\omega)^{2/3}$$

$$\Delta\Psi_2^{tidal} \sim \lambda_2 x^{5/2}$$

$$\Delta\Psi_3^{tidal} \sim \lambda_3 x^{9/2}$$

$$|\Delta\Psi_3^{tidal} / \Delta\Psi_2^{tidal}| \sim \mathcal{O}(10^{-3})$$

We hardly expect to observe higher tidal coefficients in the waveform

$$\Lambda = \lambda/M^5 \rightarrow G \left(\frac{c^2}{GM} \right)^5 \lambda = \frac{2}{3} \left(\frac{Rc^2}{GM} \right)^5 k_2$$

Accumulated GW phase (I)

the number of wave cycles in frequency domain

$$\Delta N_{\text{cyc},\Psi} = \frac{1}{2\pi} \left[\Psi(f_2) - \Psi(f_1) + (f_1 - f_2) \frac{d\Psi}{df_1} \right], \quad (7.8)$$

$f_1 = 10$ Hz,
the low frequency cutoff for
Advanced LIGO
due to seismic noises

Waveform models:

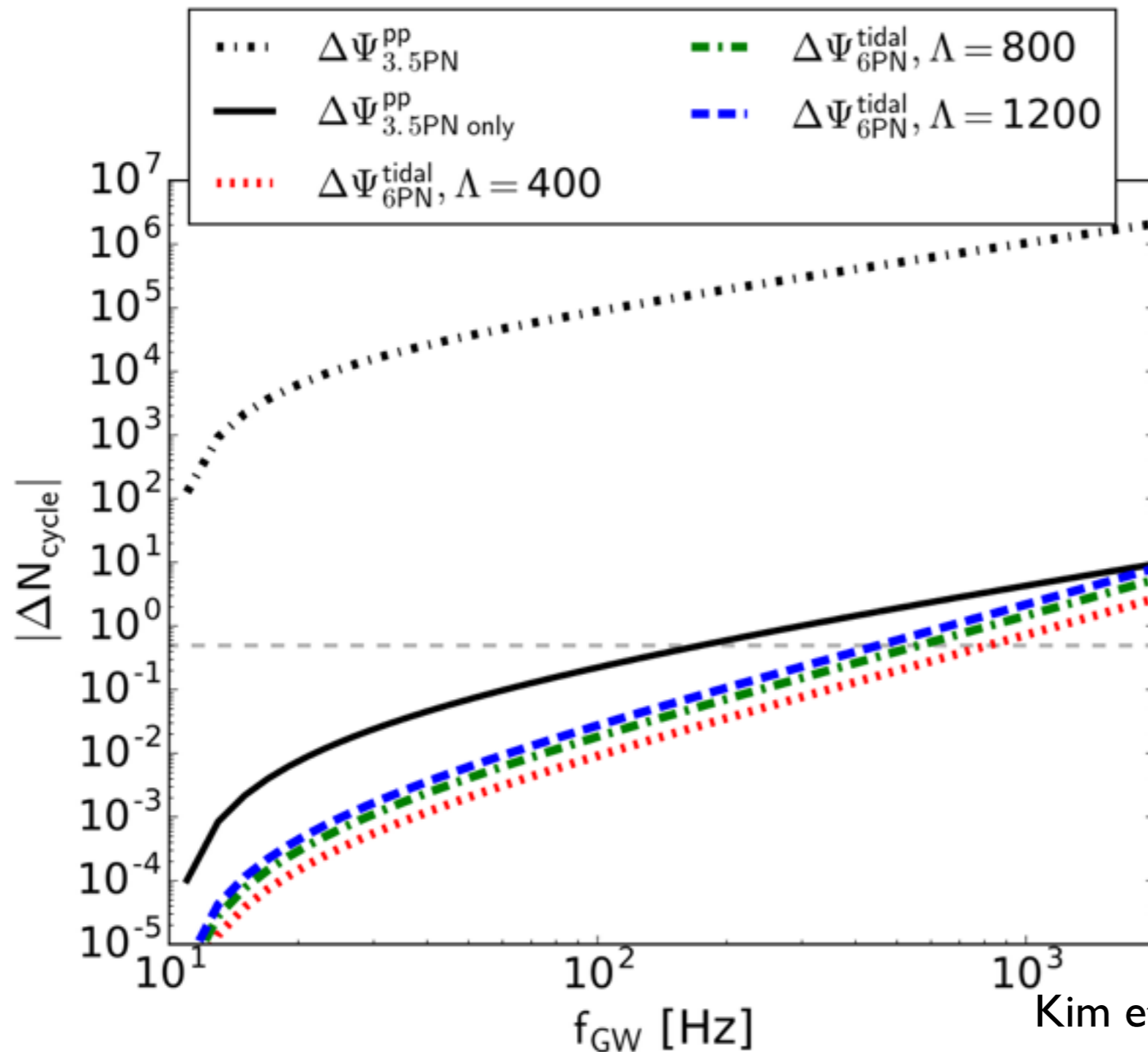
TaylorT2 for ΔN_{cyc}

TaylorF2(SPA) $\Delta N_{\text{cyc},\Psi}$

Moore et al., PRD.93.124061 (2016)

	$1.4M_{\odot} + 1.4M_{\odot}, f_2 = 1000$ Hz		
PN order	ΔN_{cyc}	$\Delta N_{\text{cyc},\Psi}$	$\Delta N_{\text{useful}}^{\text{norm}}$
0PN(circ)	16 031	986 372	1821
0PN(ecc)	-463	-36 137	-6.37
1PN(circ)	439	21 743	125
1PN(ecc)	-15.8	-1193	-0.332
1.5PN(circ)	-208	-8520	-94.8
1.5PN(ecc)	1.67	103	0.113
2PN(circ)	9.54	294	6.70
2PN(ecc)	-0.215	-15.4	-0.008 17
2.5PN(circ)	-10.6	-218	-10.6
2.5PN(ecc)	0.0443	2.61	0.004 73
3PN(circ)	2.02	18.2	2.80
3PN(ecc)	0.002 00	0.119	-0.000 238
3.5PN(circ)	-0.662	-4.39	-0.977
Total	15 785	962 445	1843

Accumulated GW phase (2)



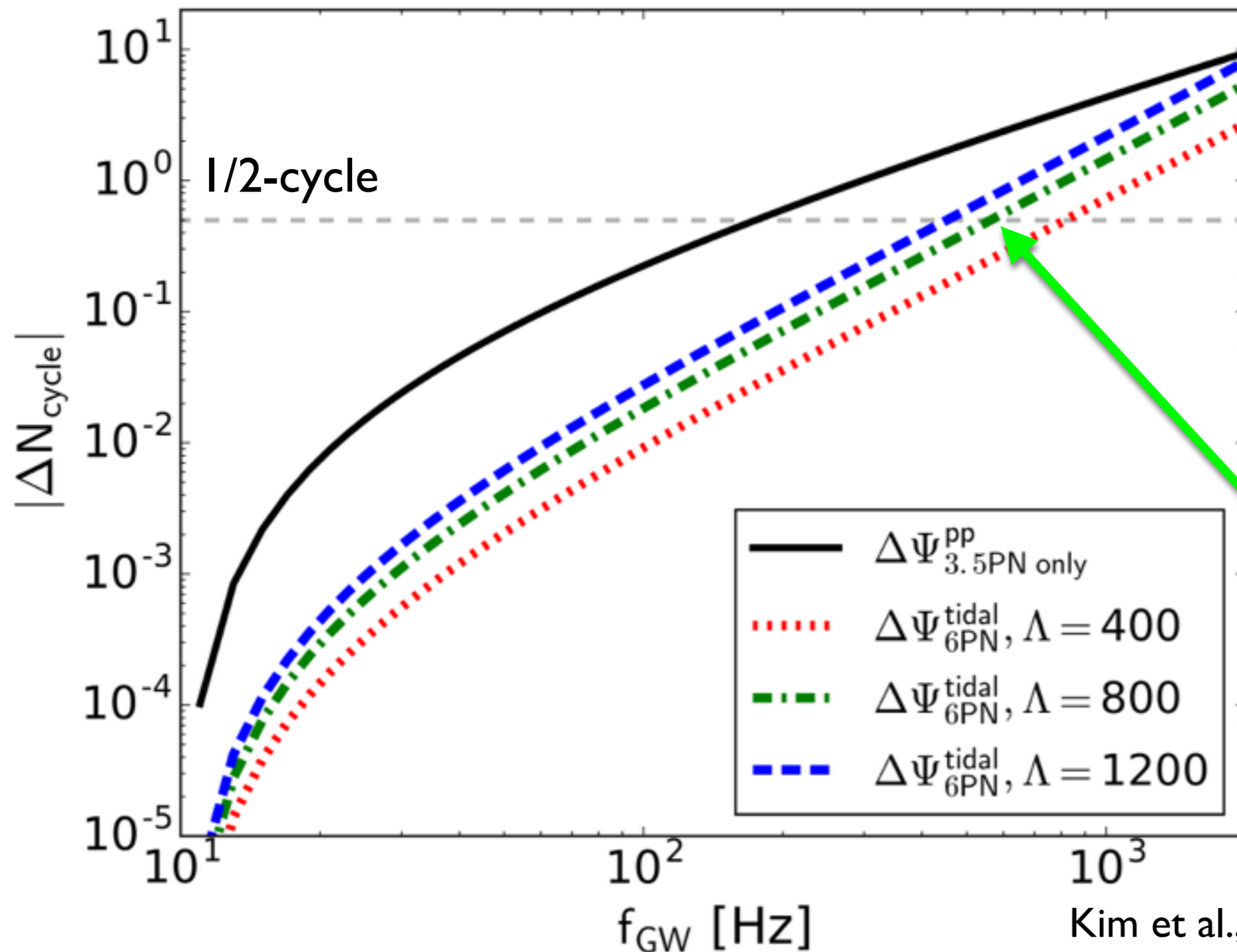
waveform model:
TaylorF2(SPA)

$M_{\text{ch}} = 1.188 M_{\odot}$

$M_1 = M_2 = 1.365 M_{\odot}$

Kim et al., New Physics: Sae Mulli (2018)

Accumulated GW phase (2)



waveform model:
TaylorF2(SPA)

$M_{\text{ch}} = 1.188 M_{\odot}$

$M_1 = M_2 = 1.365 M_{\odot}$

~ 600 Hz

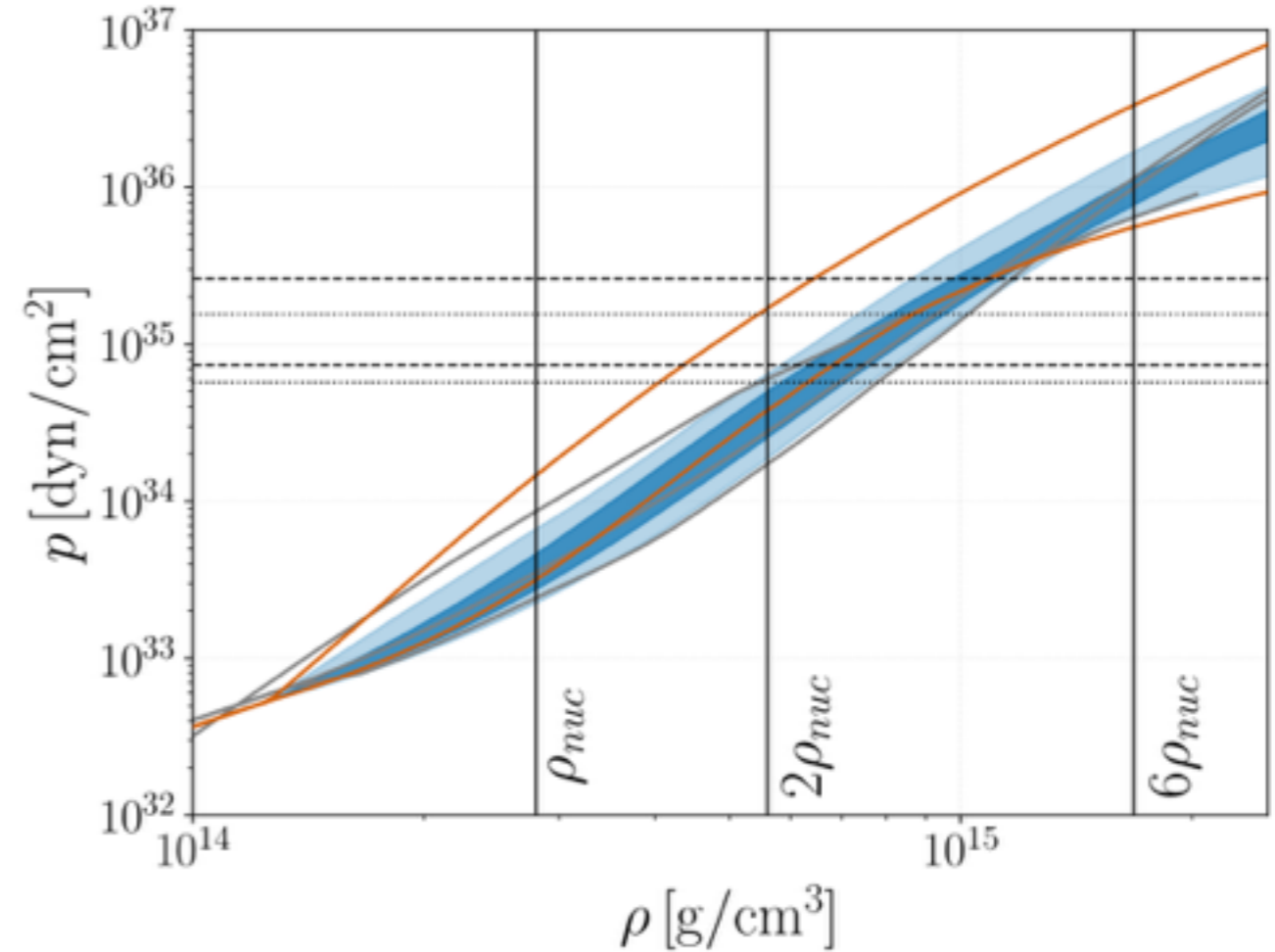
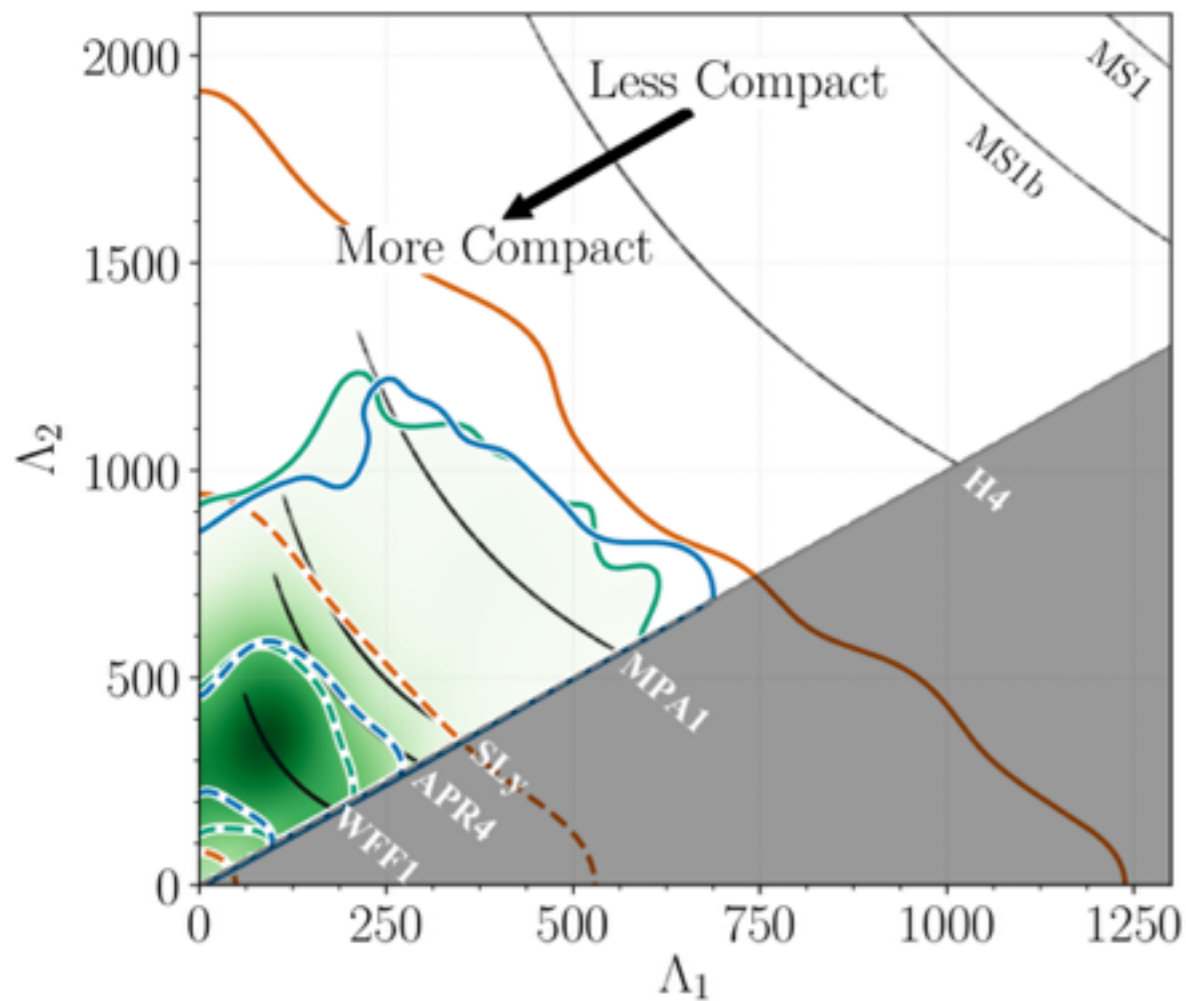
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A new constraint by GW Observation

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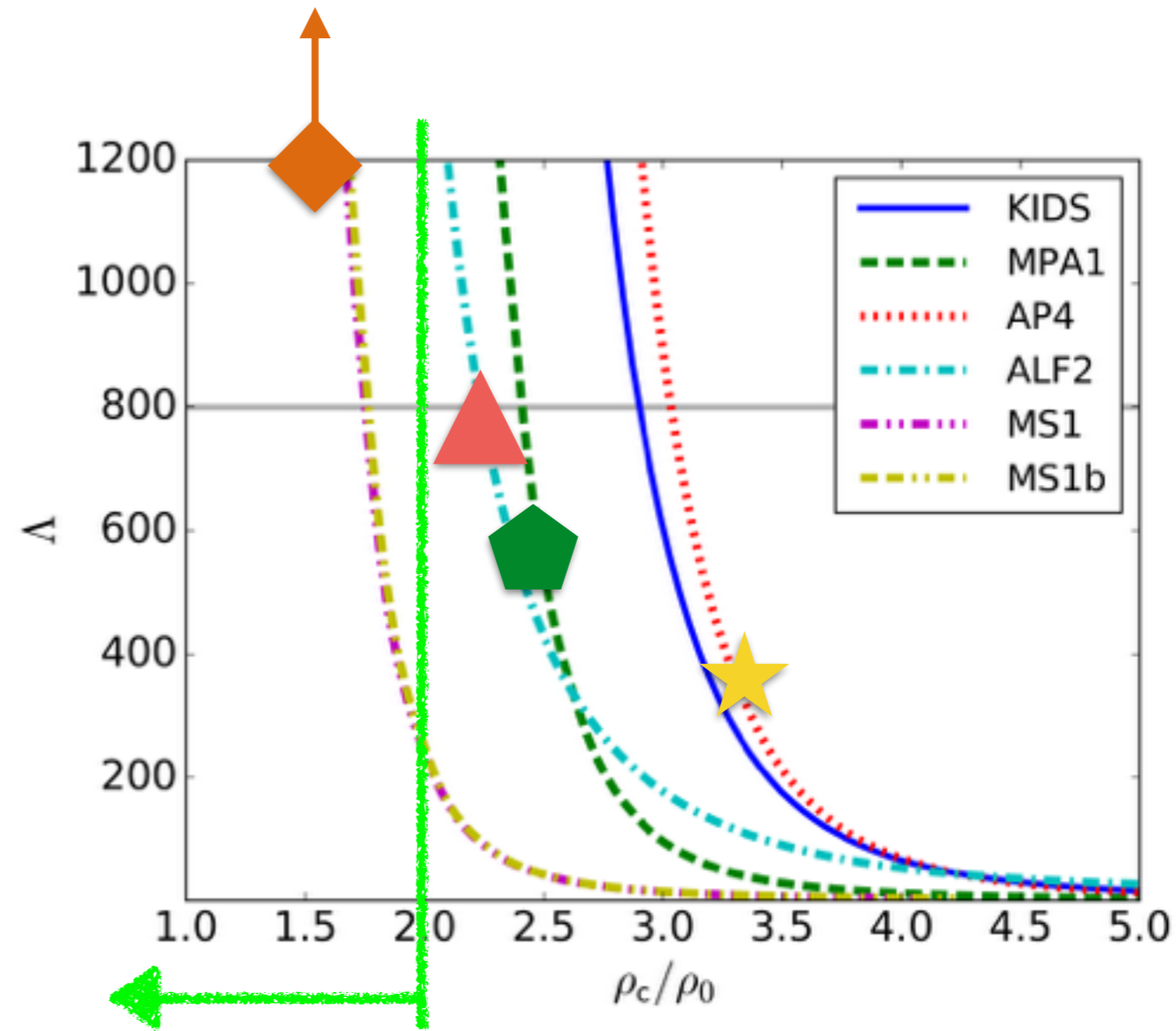
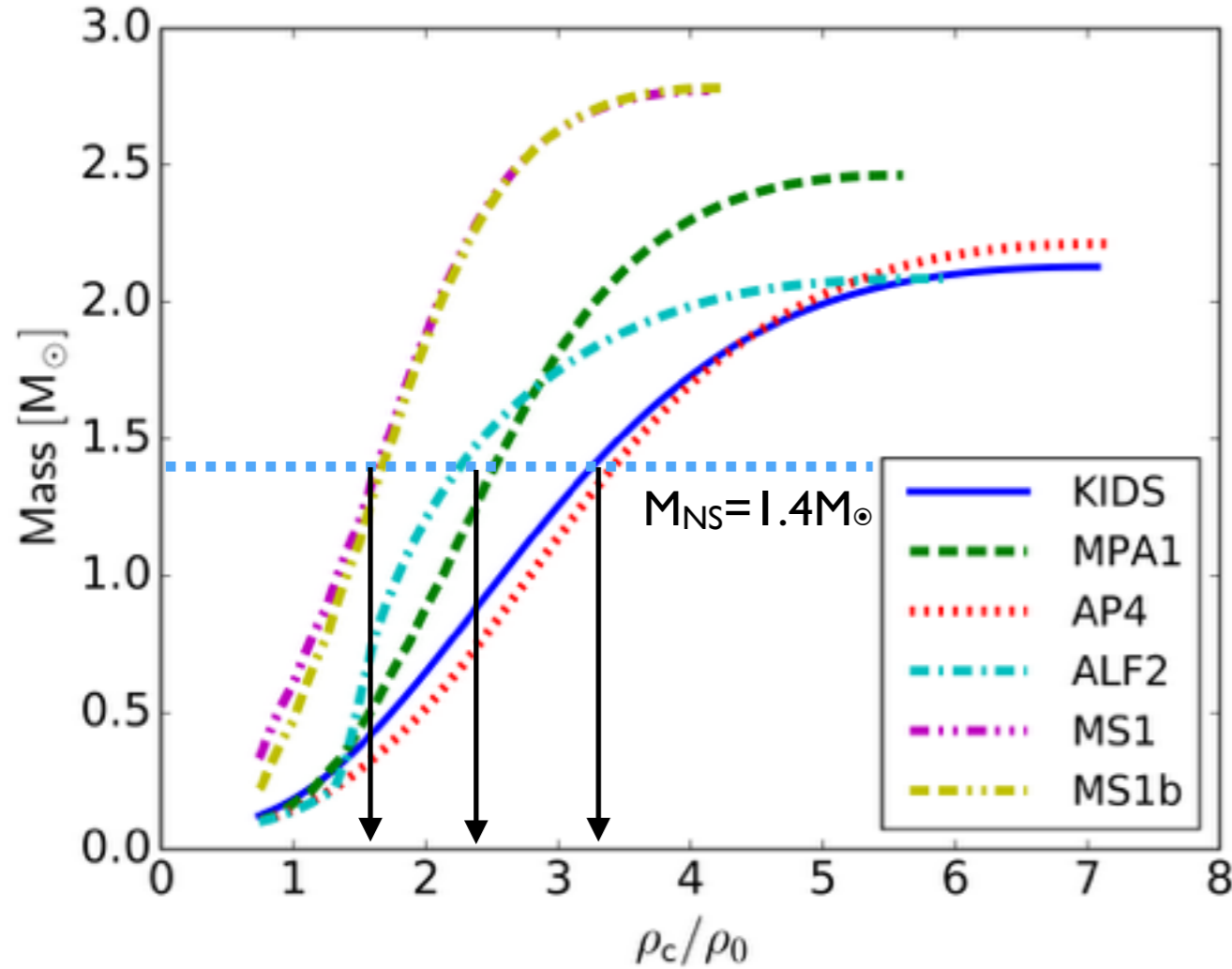


Abbott et al. (LSC and Virgo), arxiv:1805.11581 (PRL accepted)

$$\rho_{\text{nuc}} = 2.8 \times 10^{14} \text{ g/cm}^3$$

Central Density at $M_{\text{NS}}=1.4 M_{\odot}$

Kim et al., New Physics: Sae Mulli (2018)

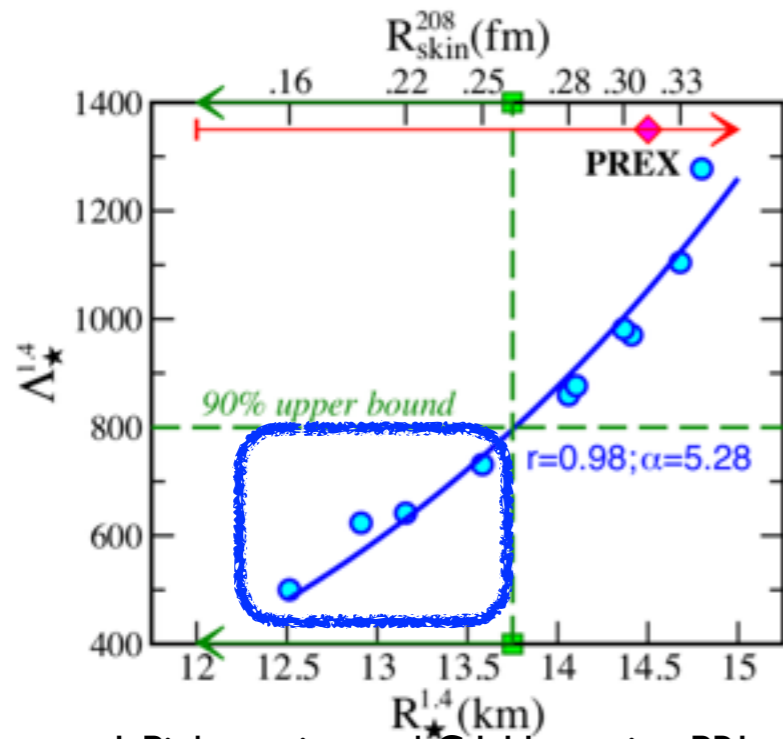


ruled out ?

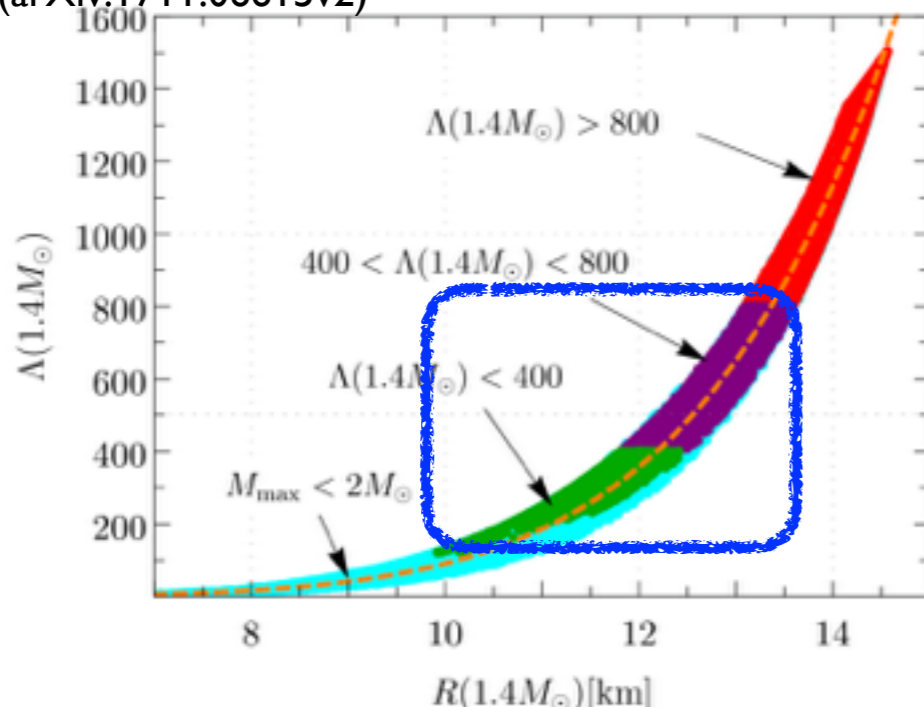
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Recent Researches



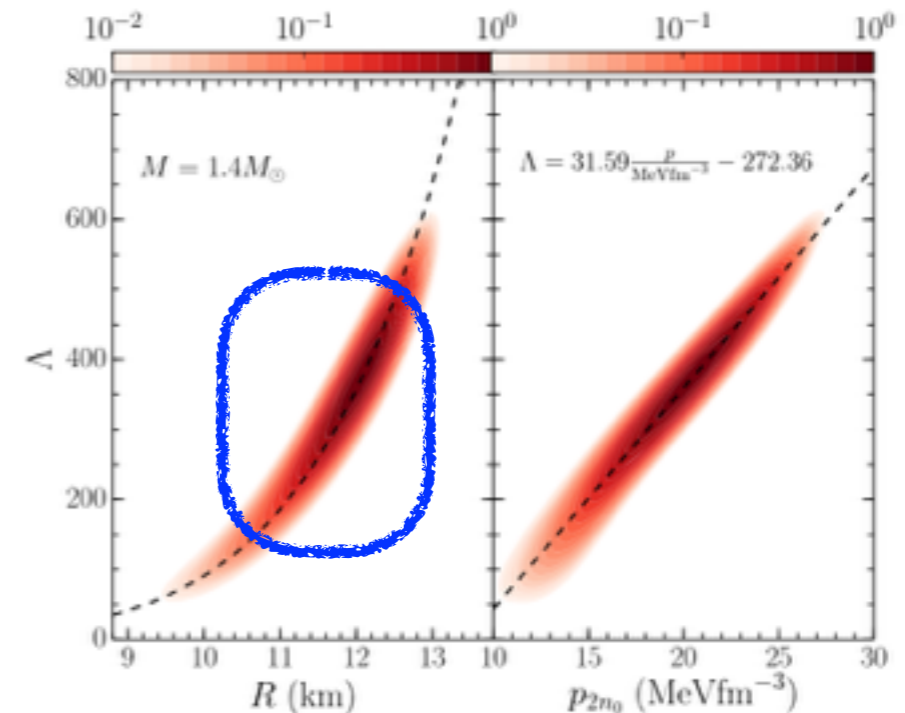
[A] F.J. Fattoyev, J. Piekarewicz, and C.J. Horowitz, PRL. 120.172702 (arXiv:1711.06615v2)



[C] E. Annala, T. Gorda, A. Kurkela, and A. Vuorinen, PRL. 120.172703 (arXiv:1711.02644v2)

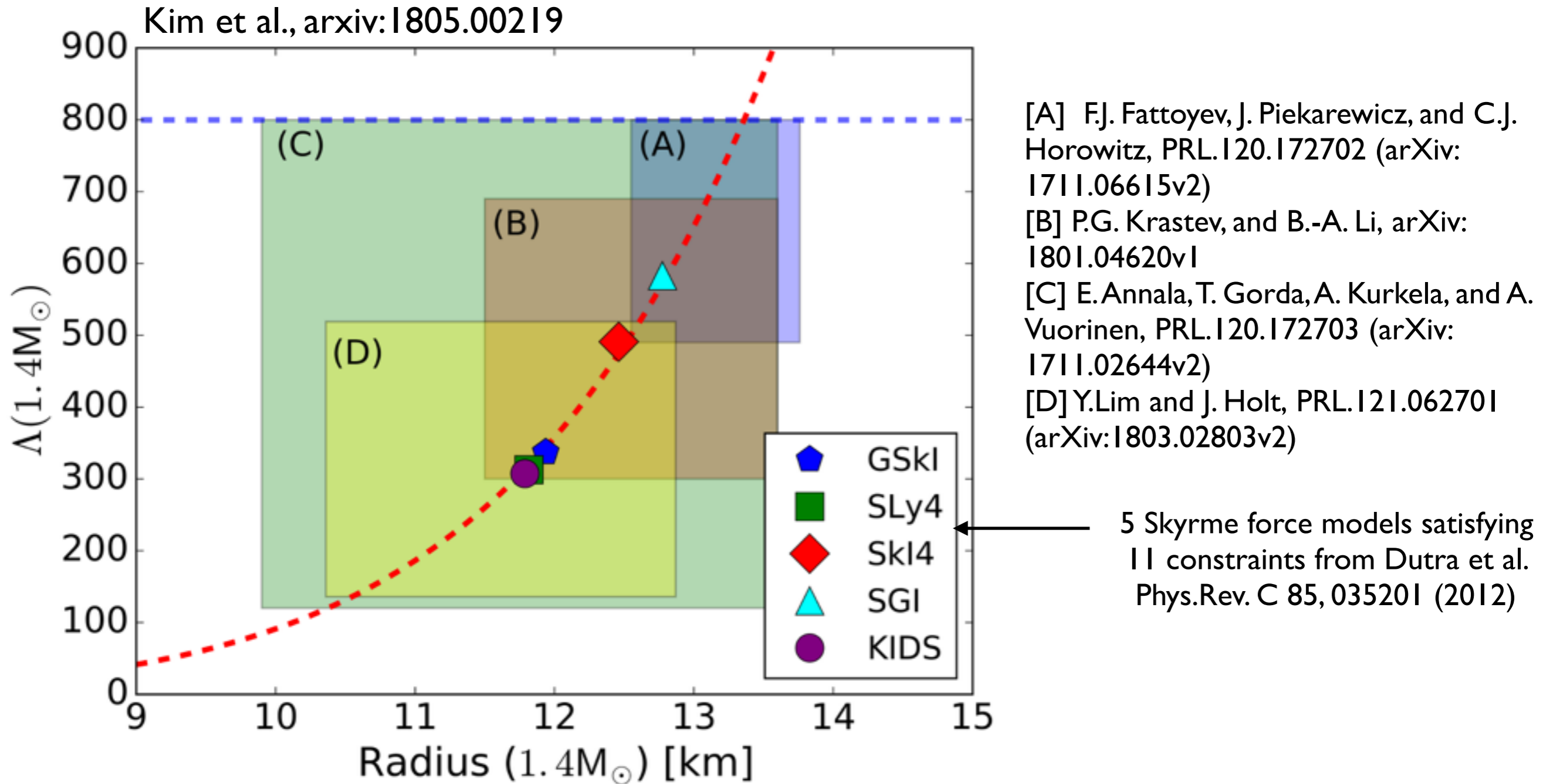
EOS	R	β	k_2	λ	L
APR	11.55	0.179	0.0721	1.48	62
MDI ($x = 0$)	11.85	0.174	0.0707	1.65	62
MDI ($x = -1$)	13.59	0.152	0.0831	3.85	107
DBHF+Bonn B	12.64	0.163	0.0946	3.06	69
FPS	10.84	0.191	0.0664	1.00	35
SLY4	11.72	0.176	0.0762	1.68	47

[B] P.G. Krastev, and B.-A. Li, arXiv:1801.04620v1



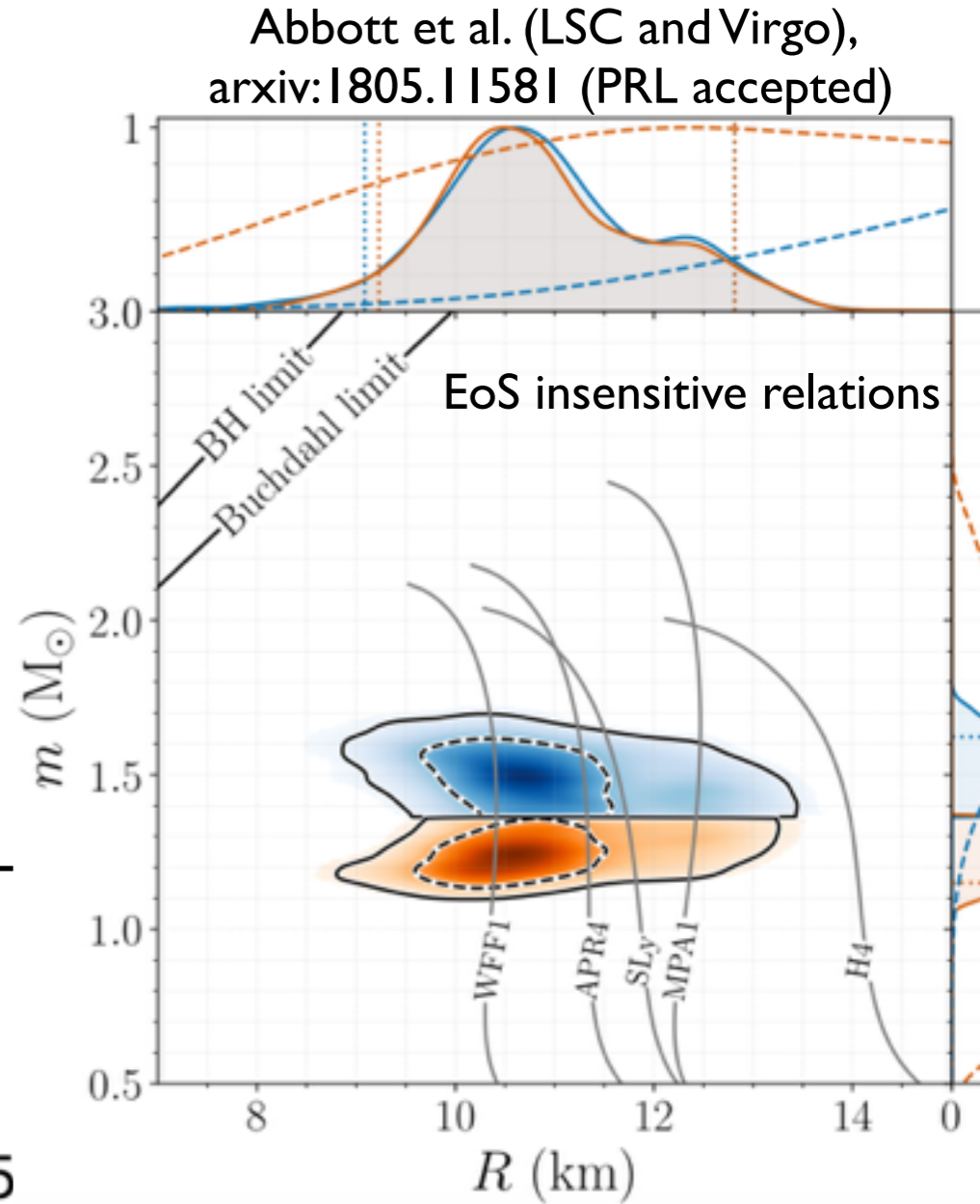
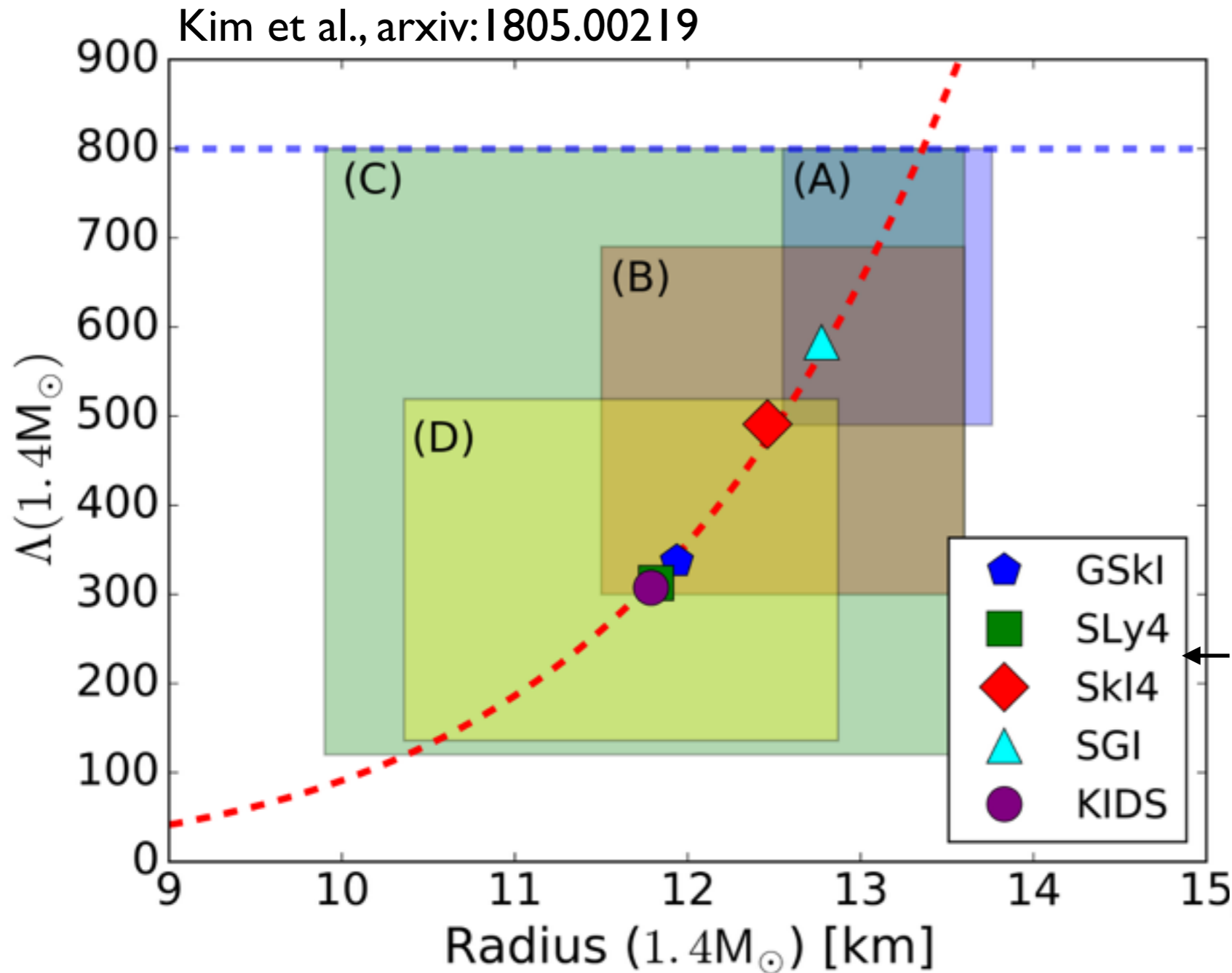
[D] Y.Lim and J. Holt, PRL. 121.062701 (arXiv:1803.02803v2)

Comparison with recent works



Red line: $\Lambda(1.4M_{\odot}) = 2.88 * 10^{-6} (R/\text{km})^{7.5}$ (fitting function in [C])

Comparison with recent works



Red line: $\Lambda(1.4M_{\odot}) = 2.88 * 10^{-6} (R/km)^{7.5}$ (fitting function in [C])

Prospects of the Observing Runs

“Prospects for Observing and Localizing Gravitational-Wave Transients with Advanced LIGO, Advanced Virgo and KAGRA”,
arXiv:1304.0670v4, LIGO-P1200087-v45, Living Rev. Relativity, 21, 3 (2018)

Epoch			2015–2016	2016–2017	2018–2019	2020+	2024+
Planned run duration			4 months	9 months	12 months	(per year)	(per year)
Expected burst range/Mpc	LIGO		40–60	60–75	75–90	105	105
	Virgo		—	20–40	40–50	40–70	80
	KAGRA		—	—	—	—	100
Expected BNS range/Mpc	LIGO		40–80	80–120	120–170	190	190
	Virgo		—	20–65	65–85	65–115	125
	KAGRA		—	—	—	—	140
Achieved BNS range/Mpc	LIGO		60–80	60–100	—	—	—
	Virgo		—	25–30	—	—	—
	KAGRA		—	—	—	—	—
Estimated BNS detections			0.05–1	0.2–4.5	1–50	4–80	11–180
Actual BNS detections			0	1	—	—	—
90% CR	% within	5 deg ²	< 1	1–5	1–4	3–7	23–30
		20 deg ²	< 1	7–14	12–21	14–22	65–73
		median/deg ²	460–530	230–320	120–180	110–180	9–12
Searched area	% within	5 deg ²	4–6	15–21	20–26	23–29	62–67
		20 deg ²	14–17	33–41	42–50	44–52	87–90

We expect to observe more BNS and/or NS-BH

Summary

1. Tidal deformability of a neutron star can be observed by gravitational-wave detection.
 - The most dominant tidal coefficient is $l=2$ electric-type coefficient λ_2 .
 - The weighted Λ in BNS was estimated by observation of GW170817
2. Tidal deformability is a new constraint on nuclear equation of states provided by GW observation.
 - more compact NS EoS is preferred.
3. In coming GW Obs., NS EoS can be studied more precisely.
4. Further investigation on Λ and NS EoS will be conducted by using Bayesian analyses on GW DA as well as HIC to study E_{sym}.

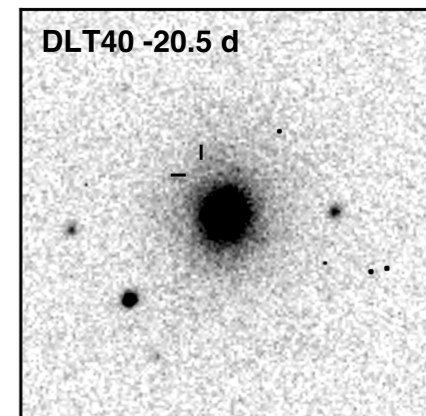
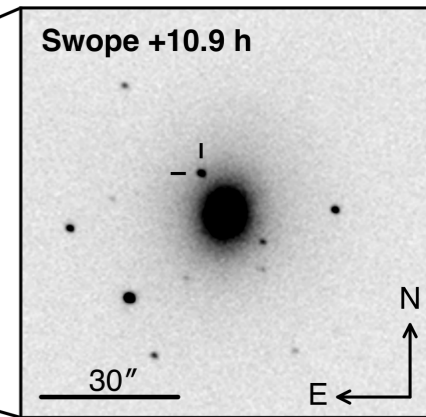
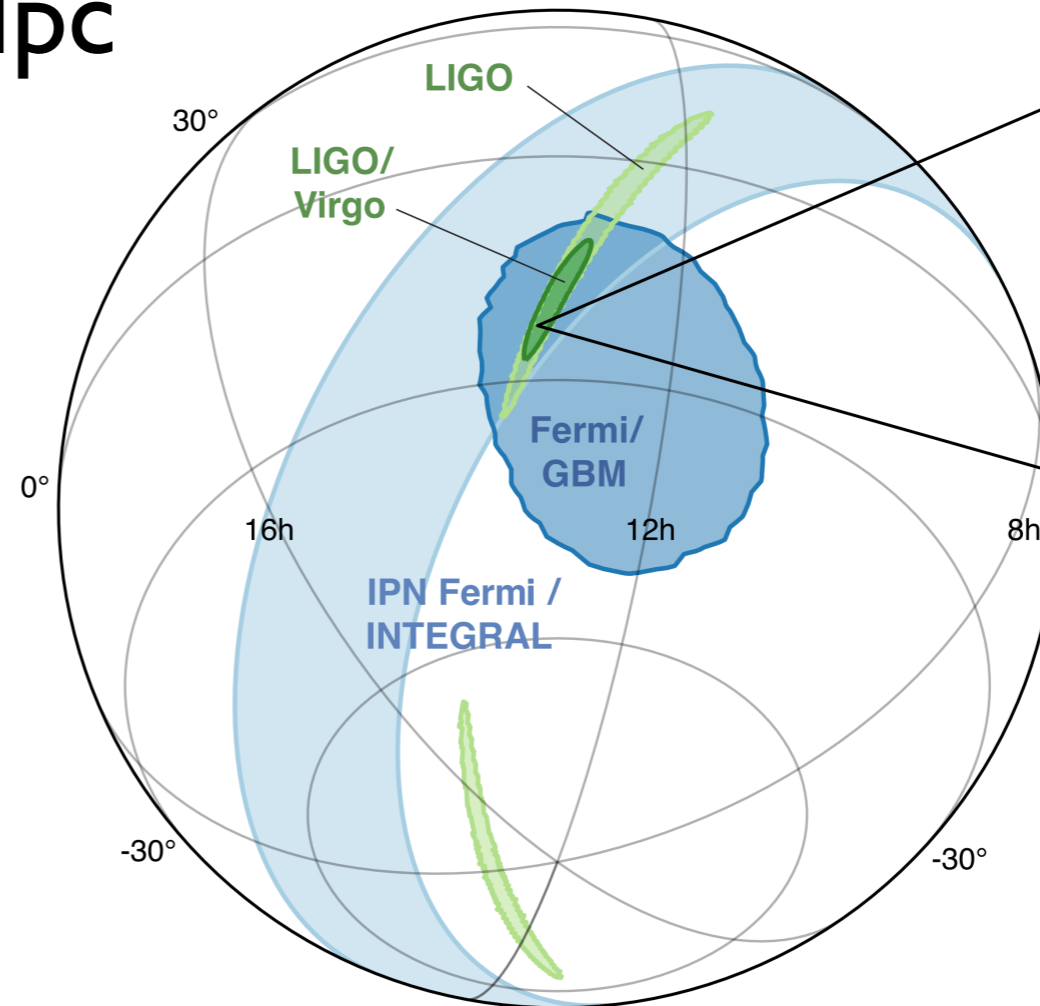
Thank you for your attention.

Extra Slides

Localization

$D_L \sim 40 \text{ Mpc}$

ApJL.848.L12(2017)



**AT2017gfo in
NGC4993**

NuSYM 2018 @ 2018.09.10

Magnetic-type tidal coefficients

Damour and Nagar , PRD 80, 084035 (2009)

Likewise,

Magnetic-type tidal coefficients

$$S_L^A = \sigma_\ell^A H_L^A.$$

$$\begin{aligned} \psi'' + \frac{e^\lambda}{r^2} [2m + 4\pi r^3(p - e)] \psi' \\ - e^\lambda \left[\frac{\ell(\ell + 1)}{r^2} - \frac{6m}{r^3} + 4\pi(e - p) \right] \psi = 0. \end{aligned} \quad (31)$$

$$\begin{aligned} G\sigma_\ell &= \frac{\ell - 1}{4(\ell + 2)} \frac{b_\ell}{(2\ell - 1)!!} \left(\frac{GM}{c_0^2} \right)^{2\ell+1} \\ &= \frac{\ell - 1}{4(\ell + 2)} \frac{j_\ell}{(2\ell - 1)!!} R^{2\ell+1}, \end{aligned}$$

$$j_\ell \equiv c^{2\ell+1} b_\ell = -c^{2\ell+1} \frac{\psi'_P(\hat{r}) - cy_{\text{odd}} \psi_P(\hat{r})}{\psi'_Q(\hat{r}) - cy_{\text{odd}} \psi_Q(\hat{r})} \Big|_{\hat{r}=1/c}.$$

$$G\sigma_2 = \frac{1}{48} j_2 R^5 = \frac{1}{48} b_2 \left(\frac{GM}{c_0^2} \right)^5,$$

$$j_2 = \frac{96c^5(2c - 1)(y - 3)}{5(2c(12(y + 1)c^4 + 2(y - 3)c^3 + 2(y - 3)c^2 + 3(y - 3)c - 3y + 9) + 3(2c - 1)(y - 3)\log(1 - 2c))}. \quad (73)$$

Magnetic-type tidal coefficients

Damour and Nagar , PRD 80, 084035 (2009)

Likewise,

Magnetic-type tidal coefficients

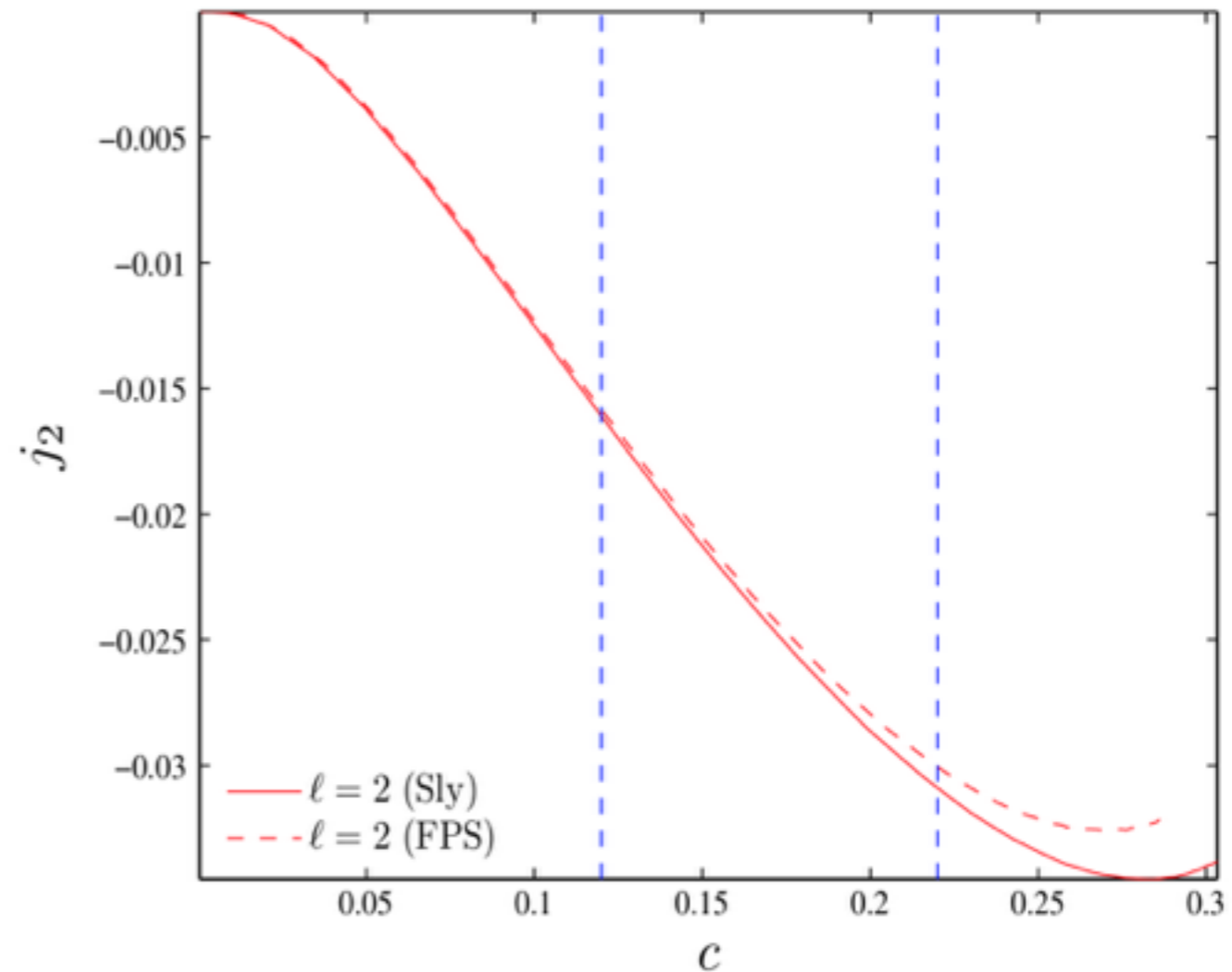
$$S_L^A = \sigma_\ell^A H_L^A.$$

$$G\sigma_\ell = \frac{\ell - 1}{4(\ell + 2)} \frac{b_\ell}{(2\ell - 1)!!} \left(\frac{GM}{c_0^2}\right)^{2\ell+1} \\ = \frac{\ell - 1}{4(\ell + 2)} \frac{j_\ell}{(2\ell - 1)!!} R^{2\ell+1},$$

$$\psi'' + \frac{\ell}{r} \psi'$$

$$j_\ell \equiv \dots$$

$$G\sigma_2 = \frac{1}{48} j_2 R^5 = \frac{1}{48} b_2 \left(\frac{GM}{c_0^2}\right)^5,$$



$$j_2 = \frac{96c^5(2c - 1)(y - 3)}{5(2c(12(y + 1)c^4 + 2(y - 3)c^3 + 2(y - 3)c^2 + 3(y - 3)c - 3y + 9) + 3(2c - 1)(y - 3)\log(1 - 2c))}. \quad (73)$$

Higher order terms in GW phase (I)

Rotational-tidal phasing of the binary neutron star waveform

- arxiv:1805.01882

$$\Psi = \frac{3M}{128\mu} x^{-2.5} \left[1 - \frac{39}{2} \tilde{\Lambda} x^5 + \tilde{\Sigma} x^6 - \tilde{X} x^{6.5} - \tilde{\Lambda}_3 x^7 + \tilde{\Sigma}_3 x^8 \right], \quad (8)$$

coupled to spins

$$\tilde{X} = \frac{1}{21M^6} c^{12} \left\{ \chi^{(1)} \left[36(35 + 614q) \hat{\lambda}_2^{(1)} - (7 - 4751q) \hat{\sigma}_2^{(1)} - 2316q \hat{\lambda}_3^{(1)} - 3474q \hat{\sigma}_3^{(1)} \right] + \chi^{(2)} \left[36(35 + 614/q) \hat{\lambda}_2^{(2)} - (7 - 4751/q) \hat{\sigma}_2^{(2)} - 2316 \hat{\lambda}_3^{(2)} / q - 3474 \hat{\sigma}_3^{(2)} / q \right] \right\},$$

$$\tilde{\Lambda}_3 = \frac{4000}{9M^7} c^{14} (q \lambda_3^{(1)} + \lambda_3^{(2)} / q),$$

$$\tilde{\Sigma}_3 = \frac{29925}{11M^7} c^{14} (q \sigma_3^{(1)} + \sigma_3^{(2)} / q).$$

Higher order terms in GW phase (2)

Post-Newtonian spin-tidal couplings for compact binaries -
arxiv:1805.01487

PN order	λ_2	σ_2	$\lambda_{23,32}, \sigma_{23,32}$	λ_3	σ_3
5	LO $\propto \Lambda$				
6	NLO $\propto \delta\Lambda$	LO $\propto \Sigma$			
6.5	NNLO $\propto \tilde{\Lambda}$	NLO $\propto \tilde{\Sigma}$	LO $\propto \tilde{\Gamma}$		
7	LO	
8	LO

$$\psi(x) = \frac{3}{128\nu x^{5/2}} \left\{ 1 + \left(\frac{3715}{756} + \frac{55}{9}\nu \right) x + \left(\frac{113}{3} \times \right. \right. \\ \left. \left. \times (\eta_1\chi_1 + \eta_2\chi_2) - \frac{38}{3}\nu(\chi_1 + \chi_2) \right) x^{1.5} + O(x^2) \right. \\ \left. + \Lambda x^5 + (\delta\Lambda + \Sigma)x^6 + (\tilde{\Lambda} + \tilde{\Sigma} + \tilde{\Gamma})x^{6.5} + O(x^7) \right\}, \quad (15)$$

$$\Lambda = \left(264 - \frac{288}{\eta_1} \right) \frac{c^{10}\lambda_2^{(1)}}{M^5} + (1 \leftrightarrow 2), \quad (16)$$

$$\delta\Lambda = \left(\frac{4595}{28} - \frac{15895}{28\eta_1} + \frac{5715\eta_1}{14} - \frac{325\eta_1^2}{7} \right) \frac{c^{10}\lambda_2^{(1)}}{M^5} \\ + (1 \leftrightarrow 2), \quad (17)$$

$$\Sigma = \left(\frac{6920}{7} - \frac{20740}{21\eta_1} \right) \frac{c^8\sigma_2^{(1)}}{M^5} + (1 \leftrightarrow 2).$$

$$\tilde{\Lambda} = \left[\left(\frac{593}{4} - \frac{1105}{8\eta_1} + \frac{567\eta_1}{8} - 81\eta_1^2 \right) \chi_2 \right. \\ \left. + \left(-\frac{6607}{8} + \frac{6639\eta_1}{8} - 81\eta_1^2 \right) \chi_1 \right] \frac{c^{10}\lambda_2^{(1)}}{M^5} \\ + (1 \leftrightarrow 2), \quad (19)$$

$$\tilde{\Sigma} = \left[\left(-\frac{9865}{3} + \frac{4933}{3\eta_1} + 1644\eta_1 \right) \chi_2 - \chi_1 \right] \frac{c^8\sigma_2^{(1)}}{M^5} \\ + (1 \leftrightarrow 2), \quad (20)$$

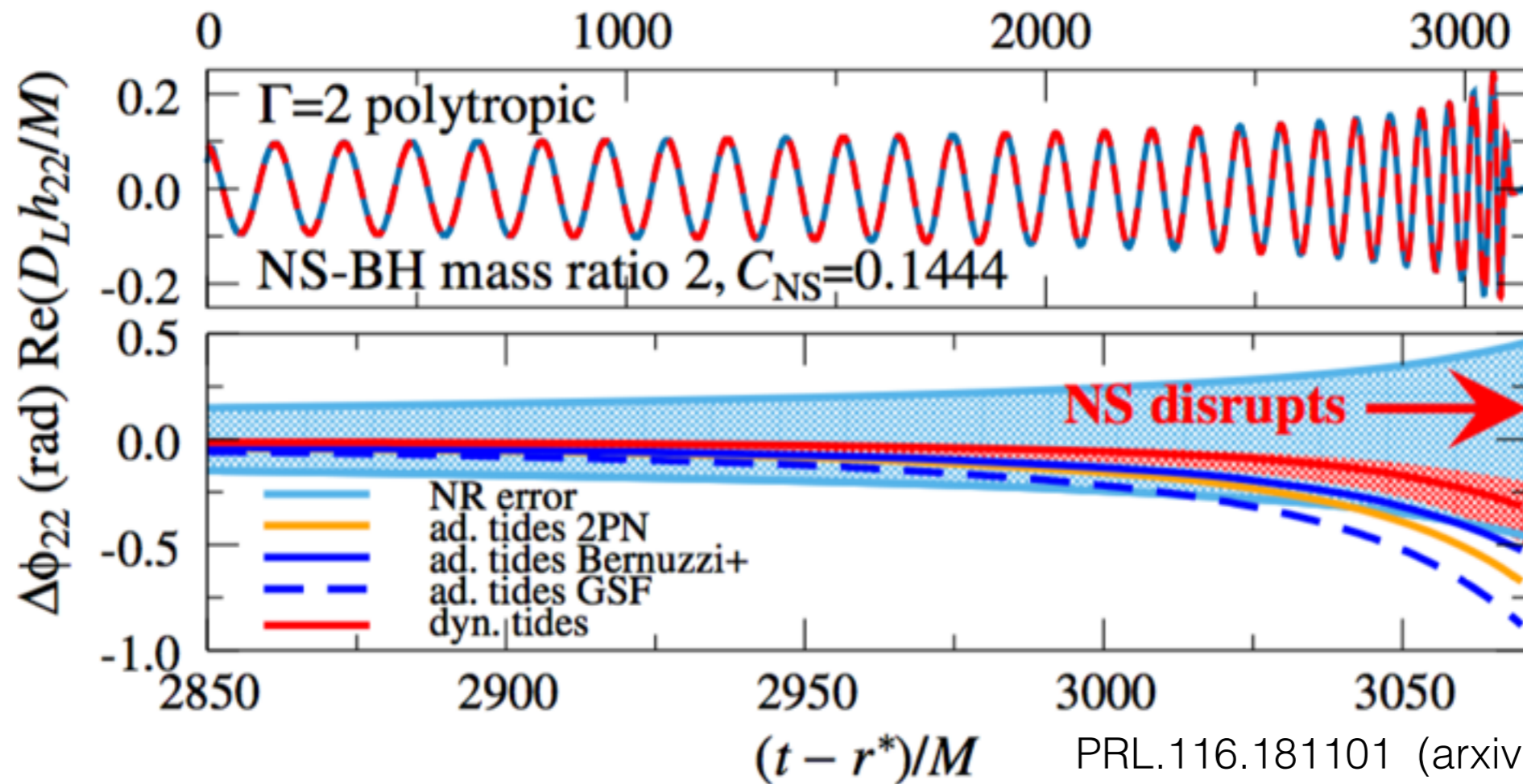
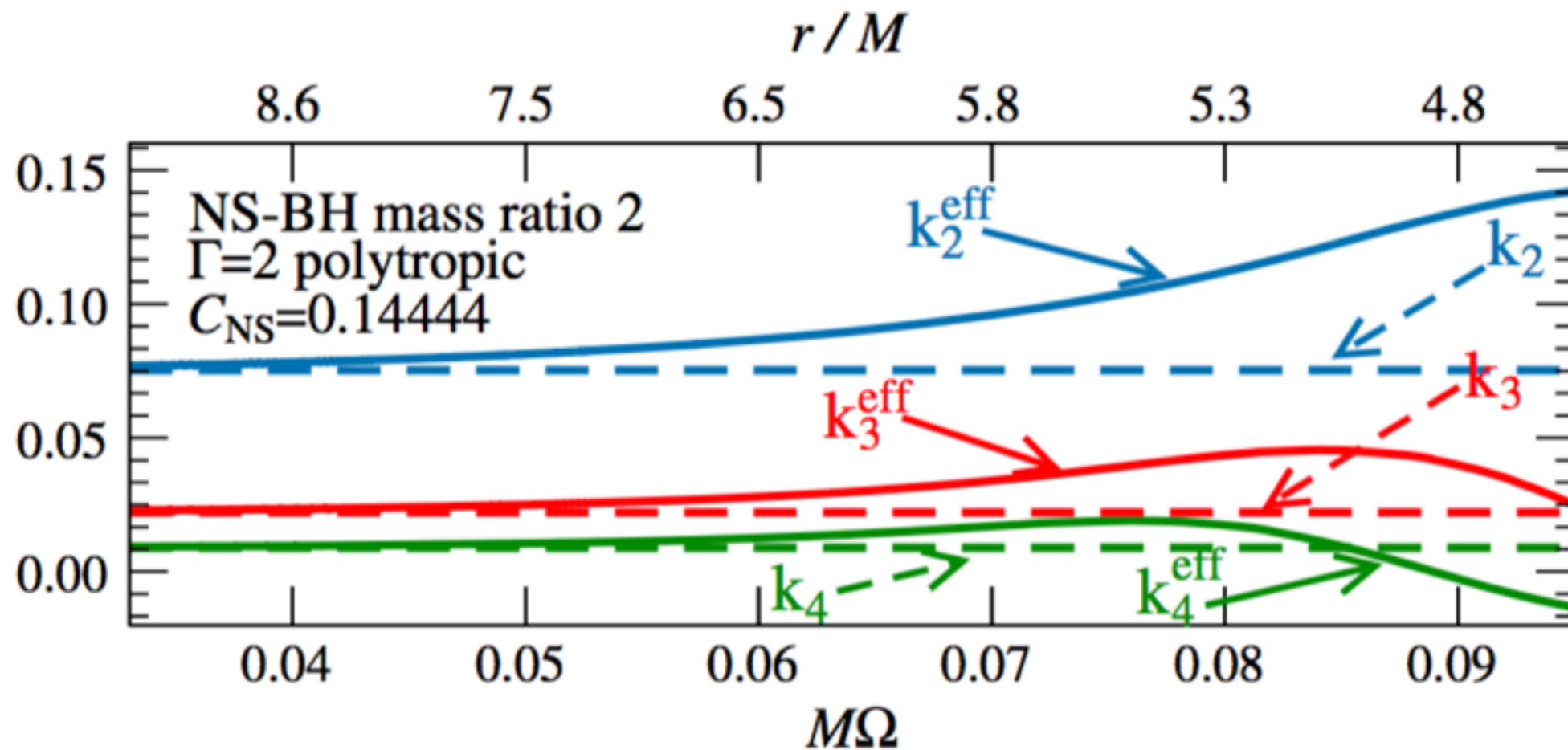
$$\tilde{\Gamma} = \frac{c^{10}\chi_1}{M^4} \left[(856\eta_1 - 816\eta_1^2) \lambda_{23}^{(1)} \right. \\ \left. - \left(\frac{4993\eta_1}{18} - \frac{2497\eta_1^2}{9} \right) \sigma_{23}^{(1)} \right. \\ \left. - \nu \left(272\lambda_{32}^{(1)} - 204\sigma_{32}^{(1)} \right) \right] + (1 \leftrightarrow 2). \quad (21)$$

Dynamic tide

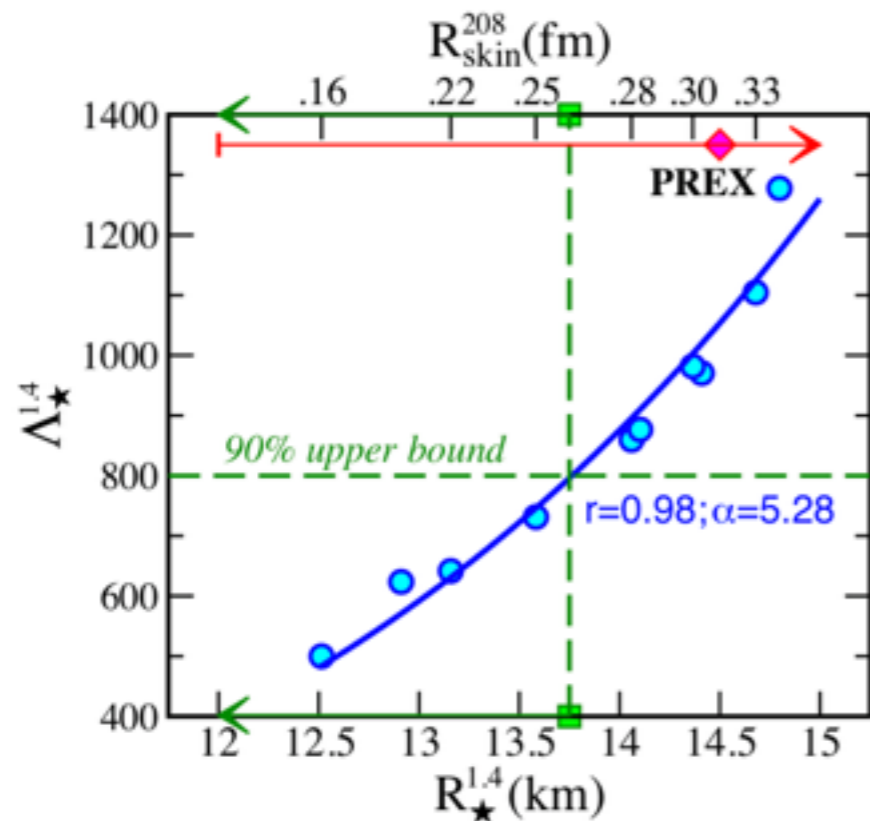
$$k_\ell^{\text{eff}} = k_\ell \left[a_\ell + \frac{b_\ell}{2} \left(\frac{Q_{m=\ell}^{\text{DT}}}{Q_{m=\ell}^{\text{AT}}} + \frac{Q_{m=-\ell}^{\text{DT}}}{Q_{m=-\ell}^{\text{AT}}} \right) \right],$$

$$\begin{aligned} \frac{Q_m^{\text{DT}}}{Q_m^{\text{AT}}} \approx & \frac{\omega_f^2}{\omega_f^2 - (m\Omega)^2} + \frac{\omega_f^2}{2(m\Omega)^2 \epsilon_f \Omega'_f (\phi - \phi_f)} \\ & \pm \frac{i\omega_f^2}{(m\Omega)^2 \sqrt{\epsilon_f}} e^{\pm i\Omega'_f \epsilon_f (\phi - \phi_f)^2} \int_{-\infty}^{\sqrt{\epsilon_f}(\phi - \phi_f)} e^{\mp i\Omega'_f s^2} ds, \end{aligned} \quad (2)$$

PRL.116.181101 (arxiv:1602.00599)



Recent Researches (I)



[A] F.J. Fattoyev, J. Piekarewicz, and C.J. Horowitz, arXiv:1711.06615v2

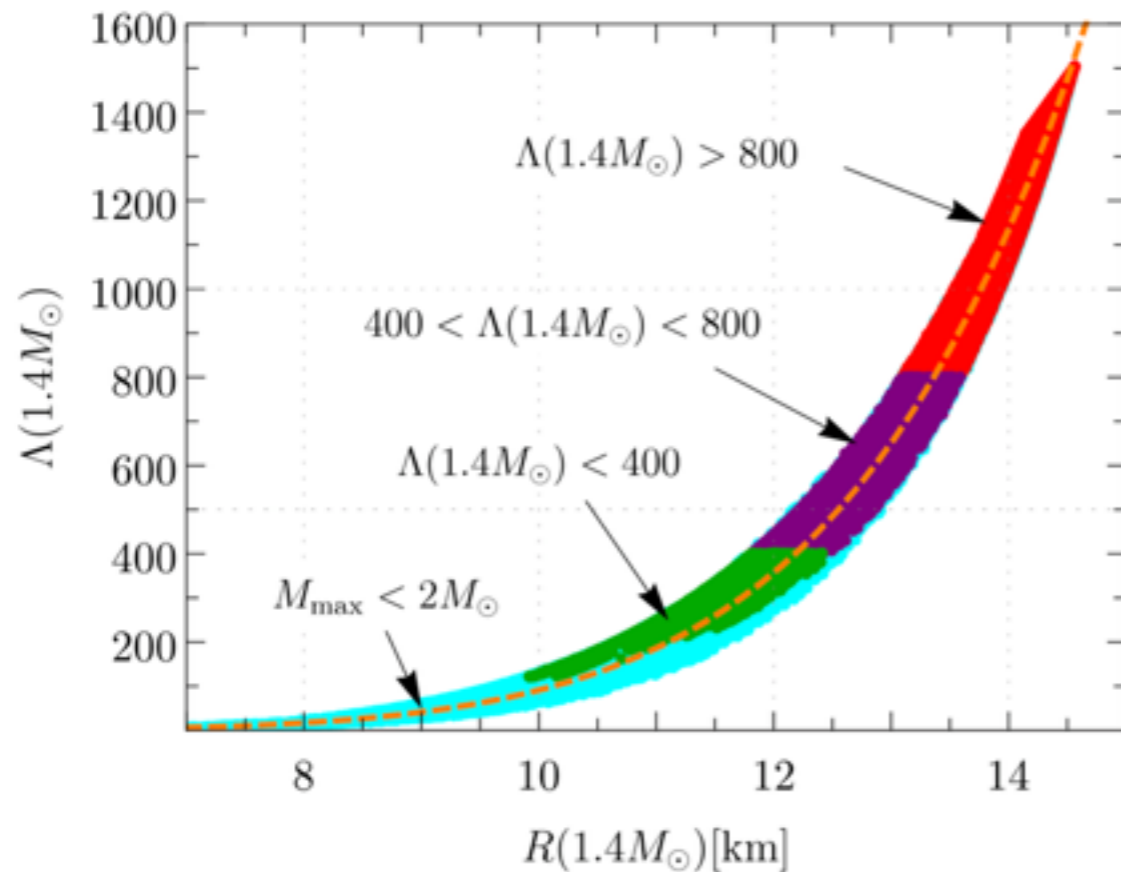
- RMF models
- Correlating neutron skin of ^{208}Pb , $\Lambda(1.4M_{\odot})$ and $R(1.4M_{\odot})$
- $490 < \Lambda(1.4M_{\odot}) < 800$
- $12.55 \text{ km} < R(1.4M_{\odot}) < 13.76 \text{ km}$

EOS	R	β	k_2	λ	L
APR	11.55	0.179	0.0721	1.48	62
MDI ($x = 0$)	11.85	0.174	0.0707	1.65	62
MDI ($x = -1$)	13.59	0.152	0.0831	3.85	107
DBHF+Bonn B	12.64	0.163	0.0946	3.06	69
FPS	10.84	0.191	0.0664	1.00	35
SLY4	11.72	0.176	0.0762	1.68	47

[B] P.G. Krastev, and B.-A. Li, arXiv:1801.04620v1

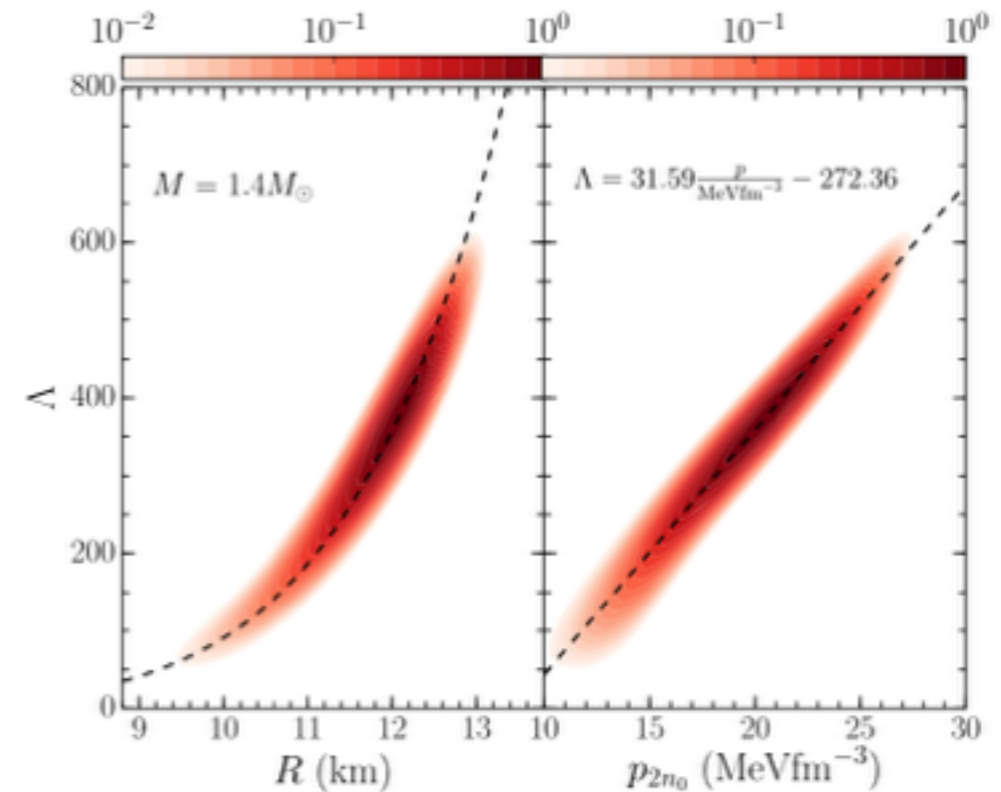
- MDI EoS
 - SNM part and symmetry energy constrained by heavy-ion reaction data up to $4.5 \rho_0$ and $1.2 \rho_0$, respectively
- $341 < \Lambda(1.4M_{\odot}) < 782$
- $11.5 \text{ km} < R(1.4M_{\odot}) < 13.6 \text{ km}$

Recent Researches (2)



[C] E. Annala, T. Gorda, A. Kurkela, and A. Vuorinen, arXiv:1711.02644v2

- Low density EoS from EFT
- High density EoS from pQCD
- Polytropic interpolation in the intermediate density
- $120 < \Lambda(1.4M_{\odot}) < 800$
- $9.9 \text{ km} < R(1.4M_{\odot}) < 13.6 \text{ km}$
- $\Lambda(1.4M_{\odot}) = 2.88 * 10^{-6} (R/\text{km})^{7.5}$



[D] Y. Lim and J. Holt, arXiv:1803.02803

- Prediction with uncertainties inherent EFT
- ~ 73000 energy density functionals
- $350 < \Lambda(1.4M_{\odot}) < 540$
- $11.65 \text{ km} < R(1.4M_{\odot}) < 12.84 \text{ km}$

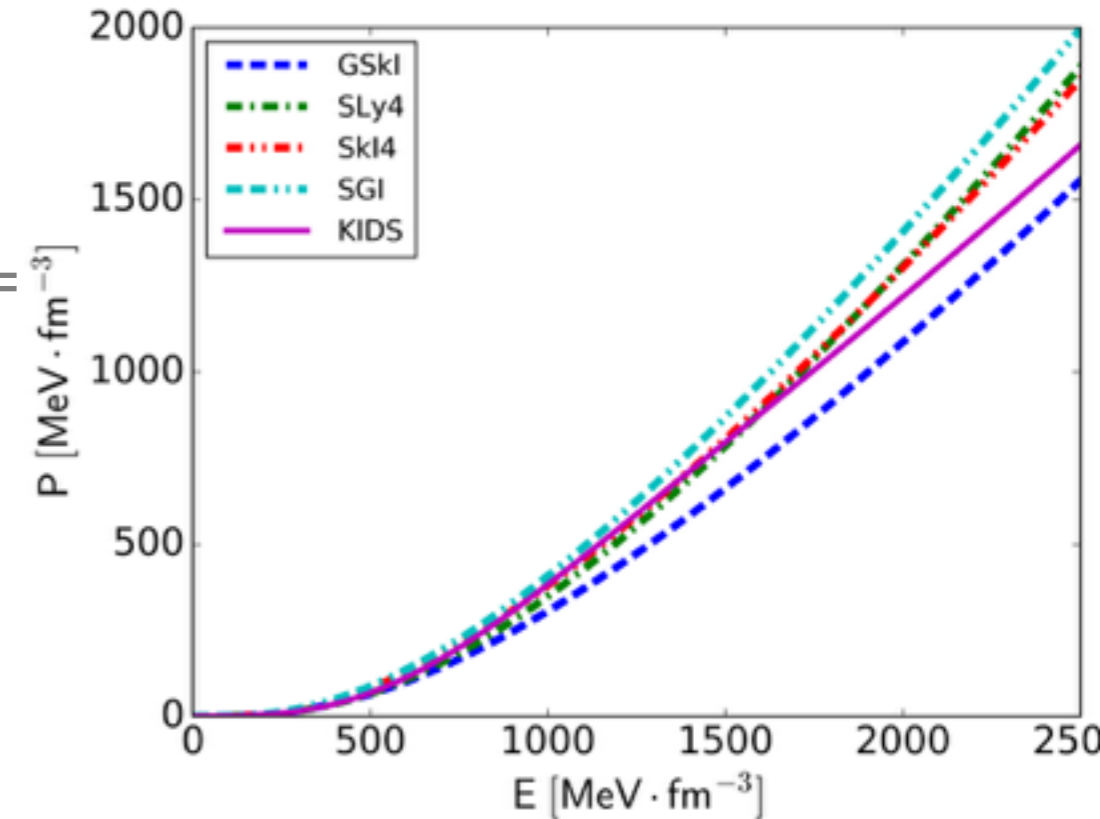
Constraints on Nuclear EoS

- Nuclear data: hundreds of models (Skyrme force, RMF, ...)
- Neutron star maximum mass
 $1.97 \pm 0.04 M_{\odot}$ [Nature 467, 1081 (2010)]
 $2.01 \pm 0.04 M_{\odot}$ [Science 340, 448 (2013)]
- II experimental/empirical data for nuclear matter around saturation density [Phys.Rev. C 85, 035201 (2012)]

Constraint	Quantity	Eq.	Density Region	Range of constraint exp/emp	Range of constraint from CSkP	Ref.
SM1	K_{\circ}	(7), (15)	ρ_{\circ} (fm^{-3})	200 – 260 MeV	202.0 – 240.3 MeV	[64]
SM2	$K' = -Q_{\circ}$	(8), (16)	ρ_{\circ} (fm^{-3})	200 – 1200 MeV	362.5 – 425.6 MeV	[65]
SM3	$P(\rho)$	(6)	$2 < \frac{\rho}{\rho_{\circ}} < 3$	Band Region	see Fig. 1	[78]
SM4	$P(\rho)$	(6)	$1.2 < \frac{\rho}{\rho_{\circ}} < 2.2$	Band Region	see Fig. 2	[80]
PNM1	$\frac{E_{PNM}}{E_{PNM}^{\circ}}$	(31)	$0.014 < \frac{\rho}{\rho_{\circ}} < 0.106$	Band Region	see Fig. 3	[39, 40]
PNM2	$P(\rho)$	(6)	$2 < \frac{\rho}{\rho_{\circ}} < 3$	Band Region	see Fig. 5	[78]
MIX1	J	(9)	ρ_{\circ} (fm^{-3})	30 – 35 MeV	30.0 – 35.5 MeV	[44]
MIX2	L	(10)	ρ_{\circ} (fm^{-3})	40 – 76 MeV	48.6 – 67.1 MeV	[101]
MIX3	$K_{\tau, \nu}$	(21)	ρ_{\circ} (fm^{-3})	-760 – -372 MeV	-407.1 – -360.1 MeV	[107]
MIX4	$\frac{S(\rho_{\circ}/2)}{J}$	-	ρ_{\circ} (fm^{-3})	0.57 – 0.86	0.61 – 0.67	[110]
MIX5	$\frac{3P_{PNM}}{L\rho_{\circ}}$	(41)	ρ_{\circ} (fm^{-3})	0.90 – 1.10	1.02 – 1.10	[112]

Selected EoSs

- Skyrme force models
- Basically fitted to properties of well-known nuclei
- Good saturation properties
- M_{\max} more than $2M_{\text{sun}}$



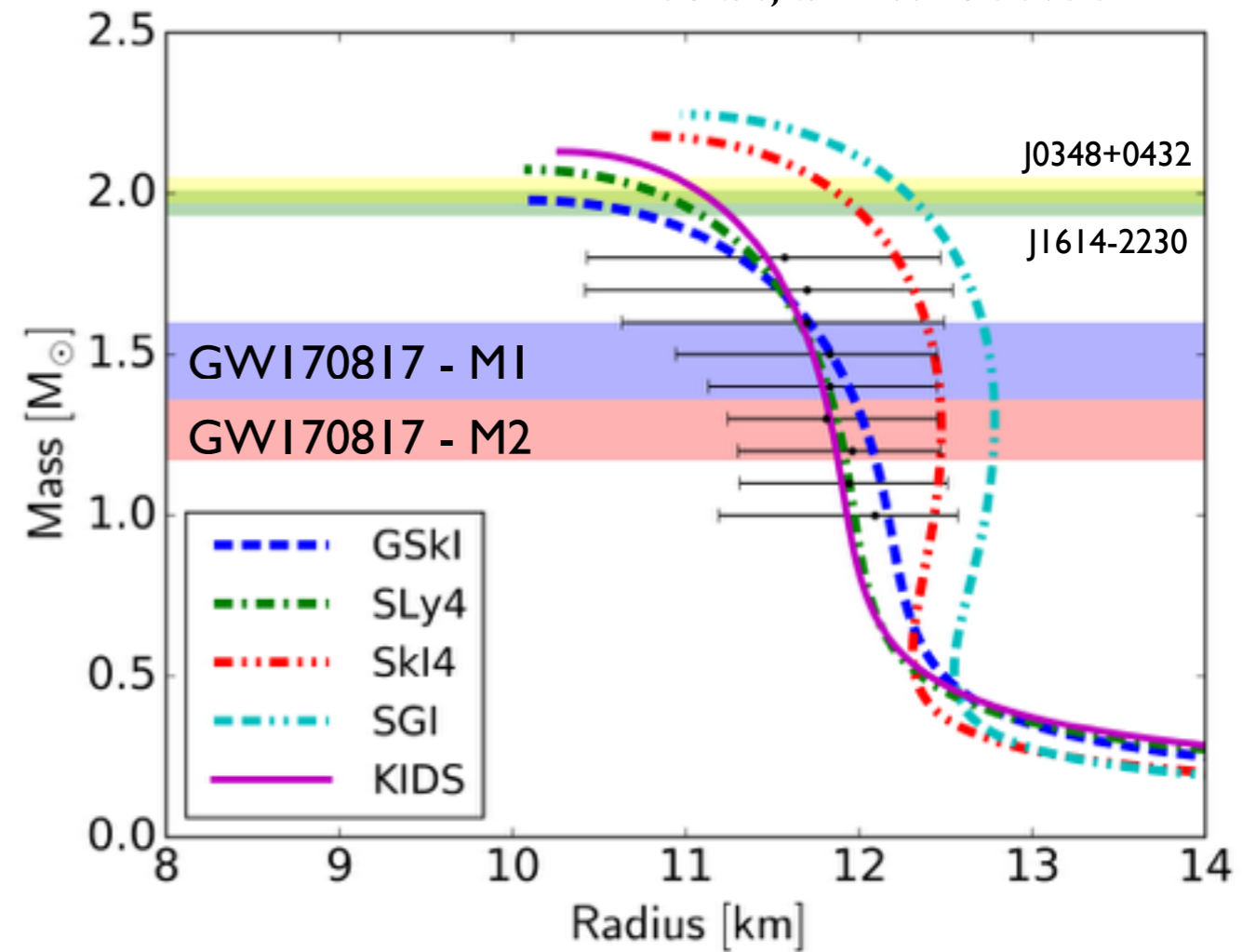
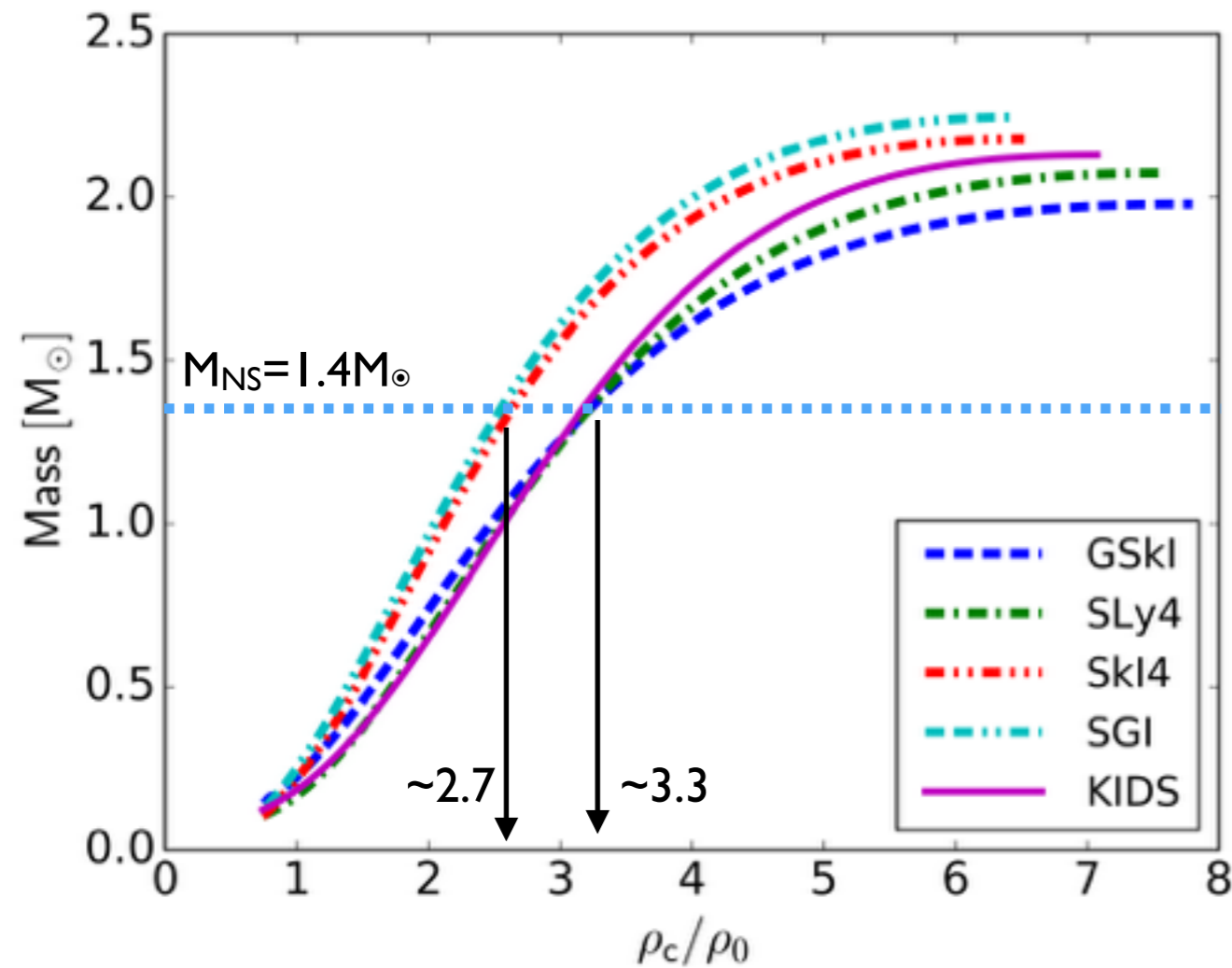
Model	ρ_0	E_0	K_0	$-Q_0$	J	L	$-K_\tau$	M_{\max}
Exp/Emp	$\simeq 0.16$	$\simeq 16.0$	200 ~ 260	200 ~ 1200	30 ~ 35	40 ~ 76	372 ~ 760	$\geq 1.93 \sim 2.05$
CSkP	-	-	202.0 ~ 240.3	362.5 ~ 425.6	30.0 ~ 35.5	48.6 ~ 67.1	360.1 ~ 407.1	-
GSkI	0.159	16.02	230.2	405.6	32.0	63.5	364.2	1.98
SLy4	0.160	15.97	229.9	363.1	32.0	45.9	322.8	2.07
SkI4	0.160	15.95	248.0	331.2	29.5	60.4	322.2	2.19
SGI	0.154	15.89	261.8	297.9	28.3	63.9	362.5	2.25
KIDS	0.160	16.00	240.0	372.7	32.8	49.1	375.1	2.14

Kim et al., arxiv:1805.00219

KIDS (Korea: IBS-Daegu-Sungkyunkwan): A new systematic expansion scheme for nuclear EDF
 [Phys. Rev. C 97, 014312 (2018)]

Mass-Radius relations

Kim et al., arxiv:1805.00219



GW170817 - Abbott et al. (LSC and Virgo), arxiv:1805.11579

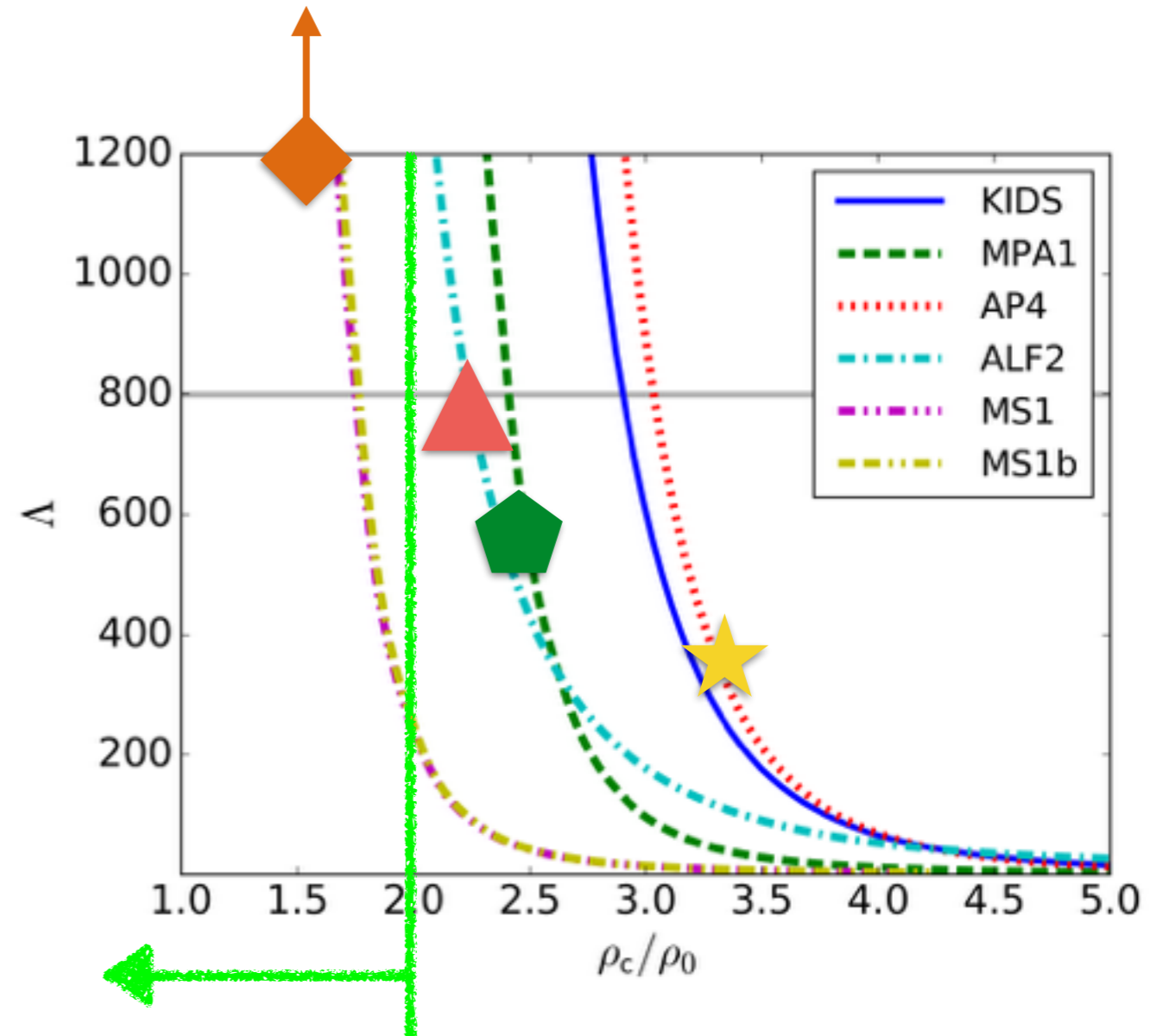
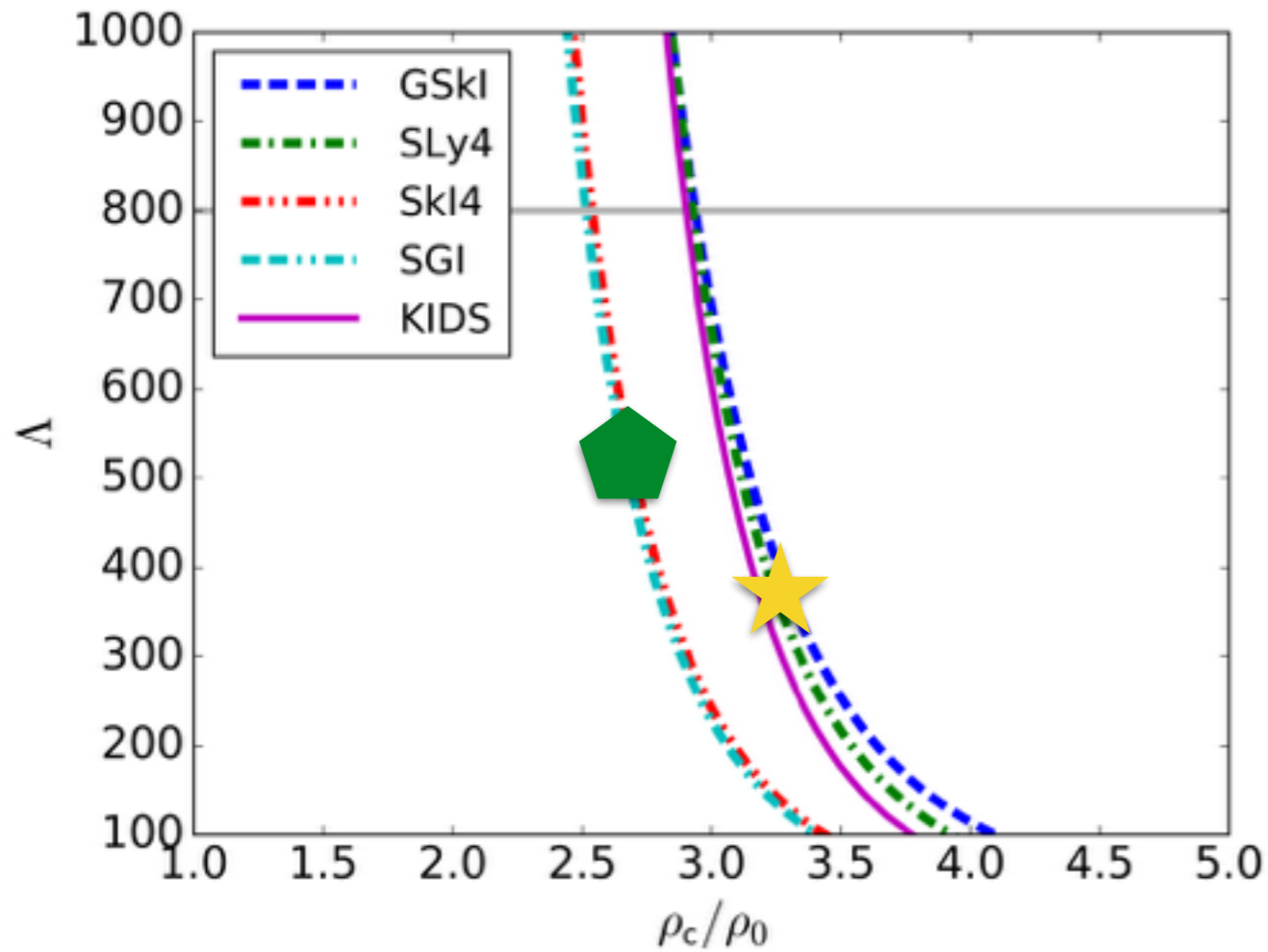
- Mchirp = $1.188 M_{\odot}$

- low spin prior : $M_1 = 1.36 \sim 1.60 M_{\odot}$, $M_2 = 1.16 \sim 1.36 M_{\odot}$

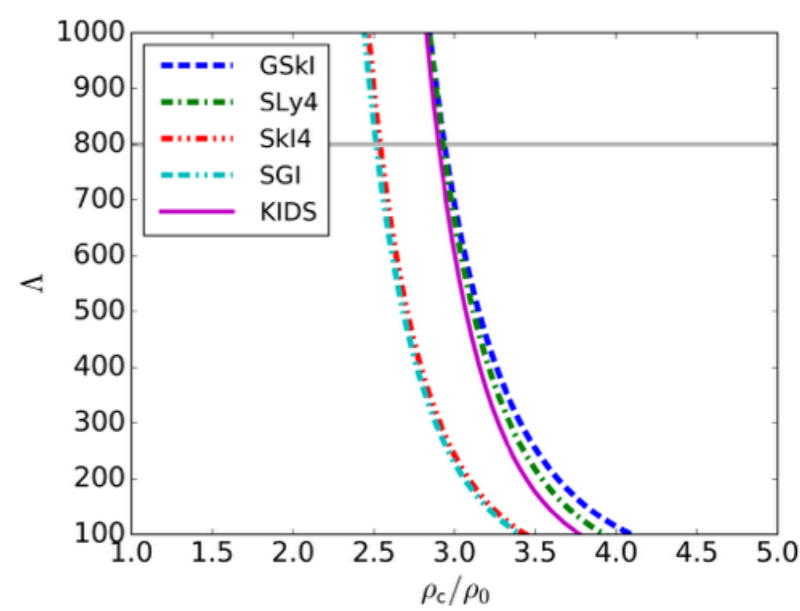
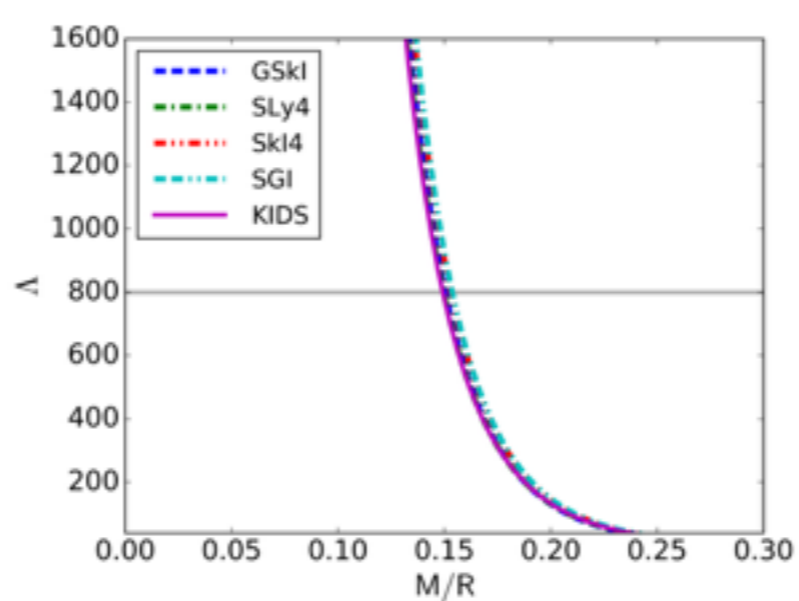
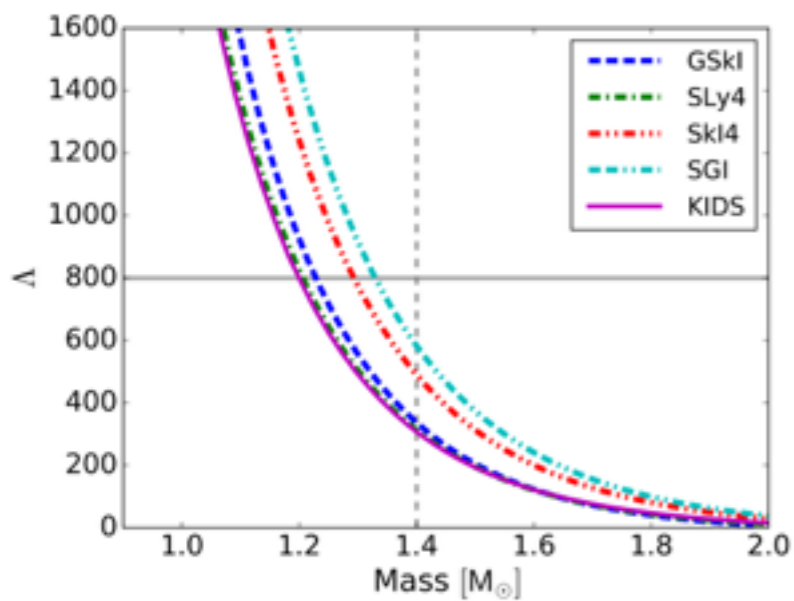
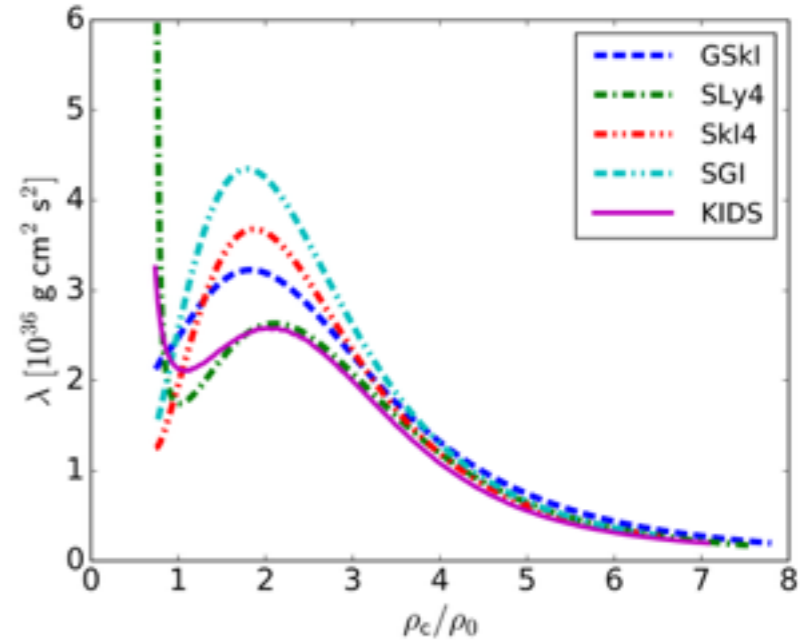
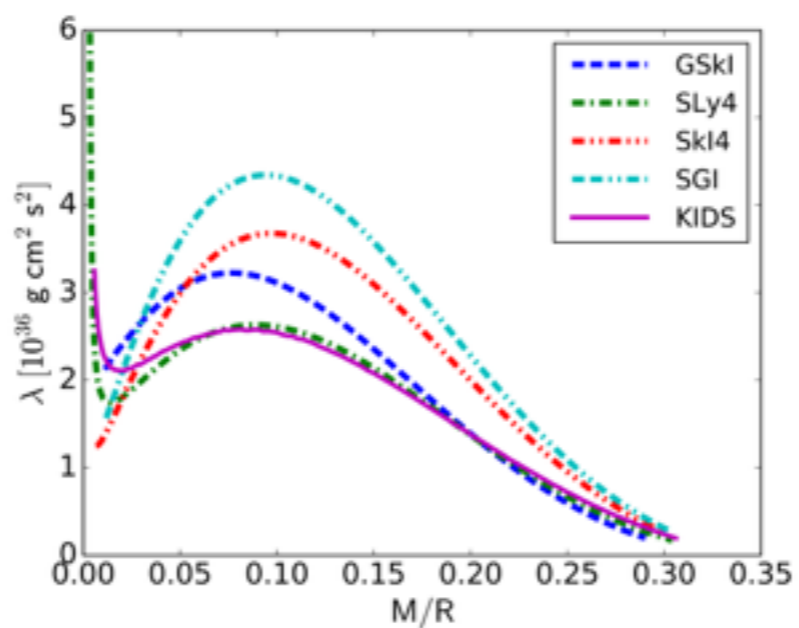
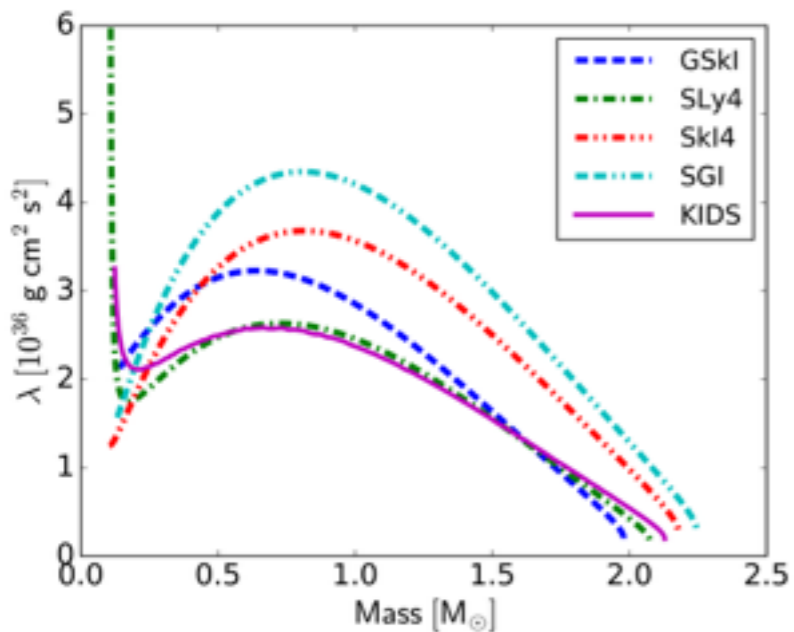
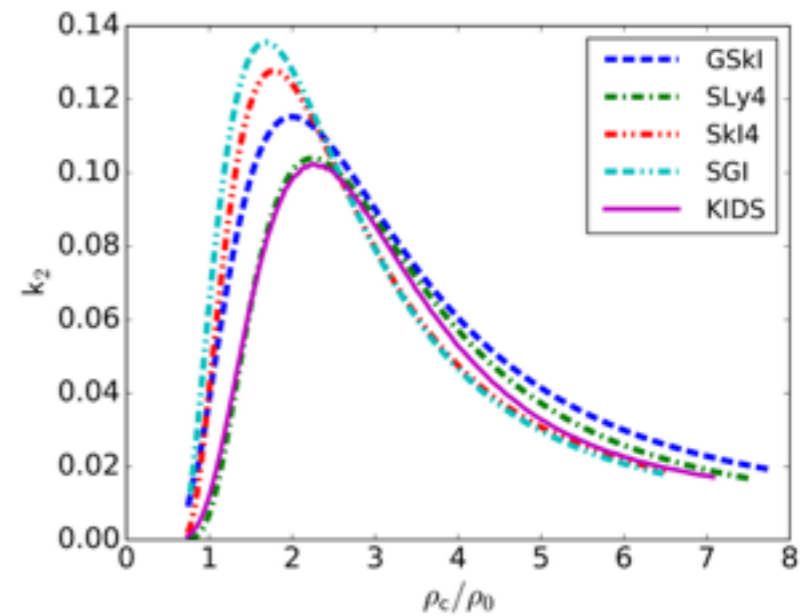
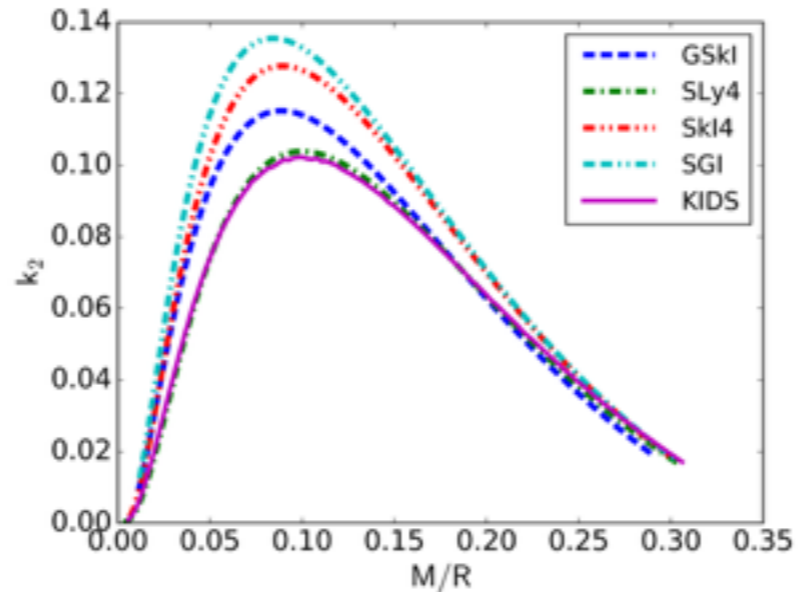
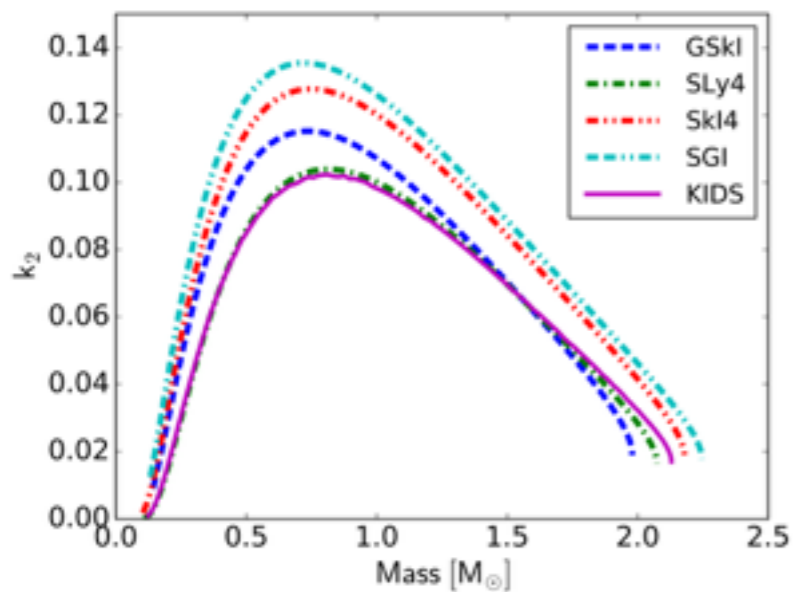
- high spin prior : $M_1 = 1.36 \sim 1.89 M_{\odot}$, $M_2 = 1.00 \sim 1.36 M_{\odot}$

Central Density at $M_{\text{NS}}=1.4 M_{\odot}$

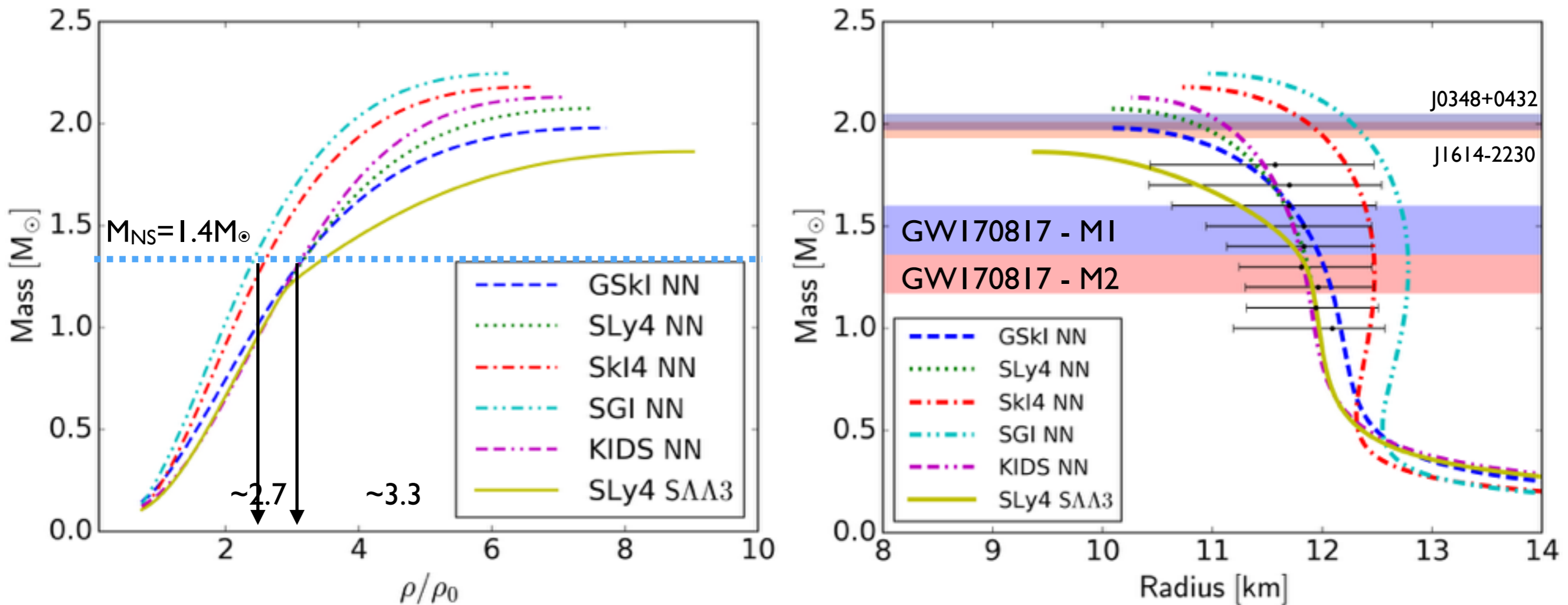
5 Skyrme force models



ruled out ?

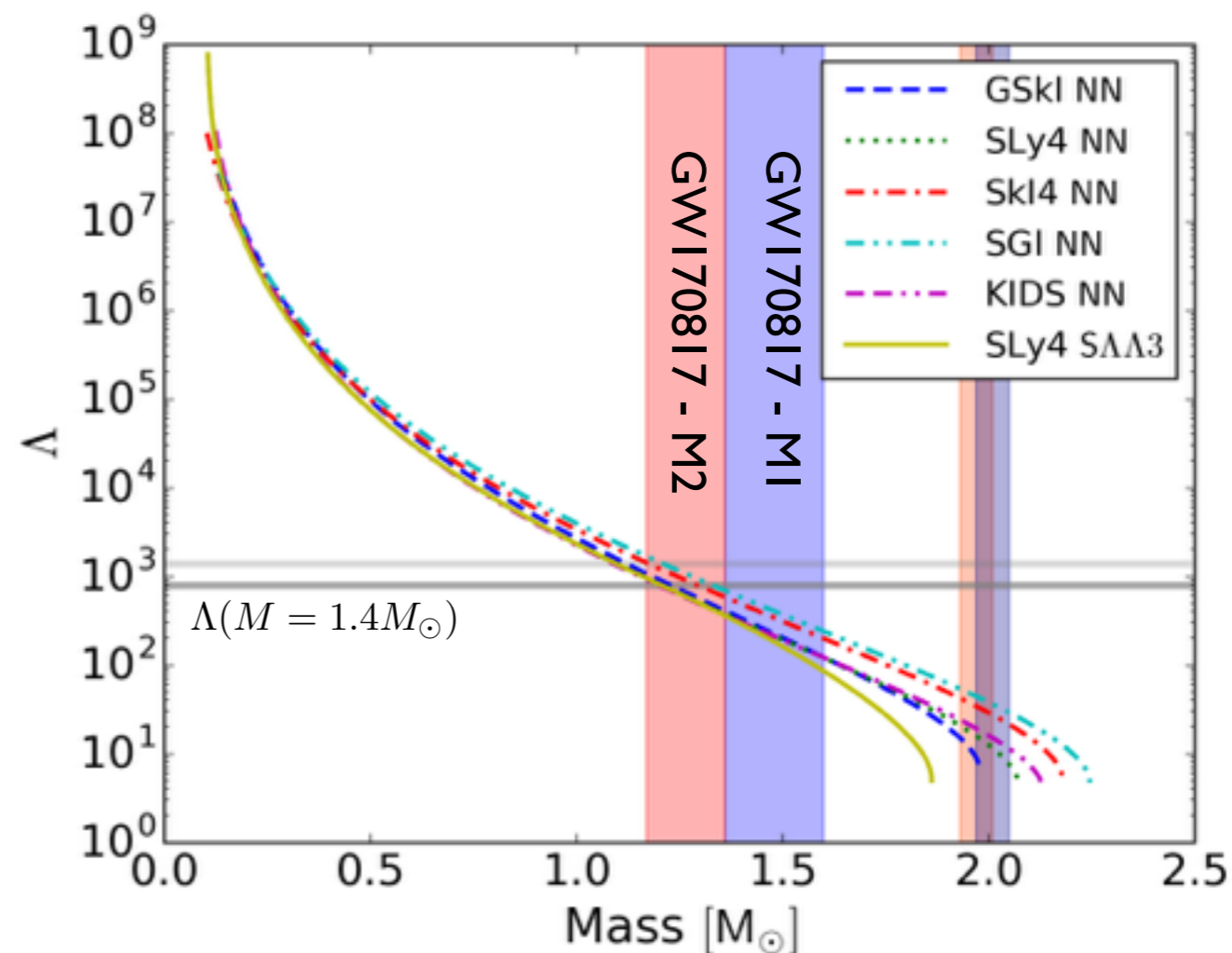
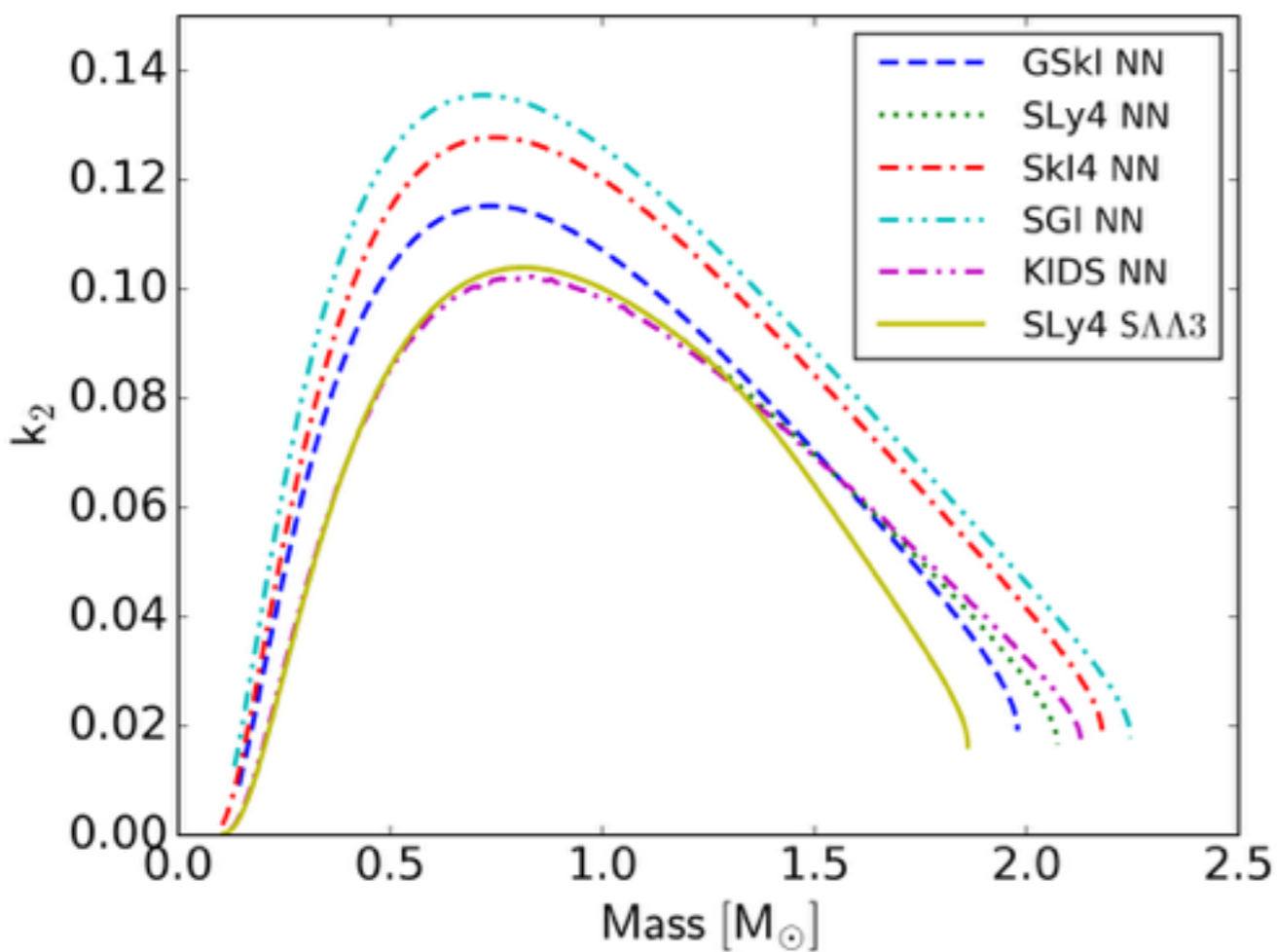


Mass-Radius relations



GW170817 - Abbott et al. (LSC and Virgo), arxiv:1805.11579
 - Mchirp = 1.188 M_{\odot}
 - low spin prior : M1 = 1.36 ~ 1.60 M_{\odot} , M2 = 1.17 ~ 1.36 M_{\odot}
 - high spin prior : M1 = 1.36 ~ 2.26 M_{\odot} , M2 = 0.86 ~ 1.36 M_{\odot}

Tidal deformability of a NS



Kim et al. in preparation

GW170817, $M_{\text{chirp}} = 1.188 M_\odot$

- low spin prior : $\Lambda(1.4 M_\odot) < 800$

- high spin prior : $\Lambda(1.4 M_\odot) < 1400$

O2 Summary

