A hand is pointing at a colorful periodic table of elements. The table is divided into sections of yellow, red, blue, and green. The text is overlaid on the top half of the image.

# The nuclear symmetry energy and the breaking of the isospin symmetry: how do they reconcile with each other?

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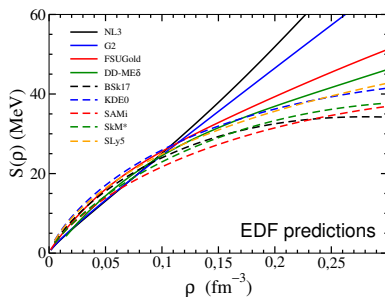
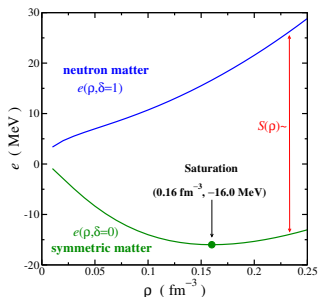
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# The Nuclear Equation of State: Infinite System



- Expansion for small asymmetries  $\delta = \frac{\rho_n - \rho_p}{\rho_n + \rho_p}$ :

$$e(\rho, \delta) = e(\rho, \delta = 0) + S(\rho)\delta^2 + \mathcal{O}[\delta^4]$$

- Expansion on the density around saturation  $x \equiv \frac{\rho - \rho_0}{3\rho}$ :

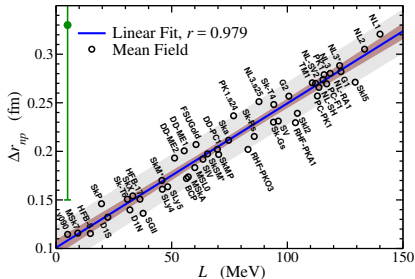
$$e(\rho, \delta) = \left[ e(\rho_0, \delta = 0) + \frac{1}{2}K_0x^2 \right] + \left[ J + Lx + \frac{1}{2}K_{\text{sym}}x^2 \right] \delta^2 + \mathcal{O}[\delta^2, \rho^3]$$

Uncertainties on  $S(\rho)$  around saturation (mainly due to  $L$ ) **impact** on many nuclear physics and astrophysics **observables**.

# Example: $L$ and the neutron skin in $^{208}\text{Pb}$

$$\Delta r_{np} \equiv \langle r_n^2 \rangle^{1/2} - \langle r_p^2 \rangle^{1/2}$$

$$\text{Macroscopic model: } \Delta r_{np} \sim \frac{1}{12} \frac{(N-Z)R}{A} \frac{R}{J} L \quad (L \propto p_0^{\text{neut}})$$

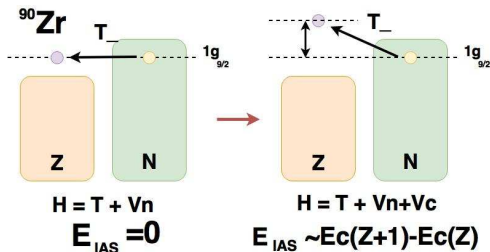


*Physical Review Letters* **106**, 252501 (2011)

The faster the symmetry energy increases with density ( $L$ ), the largest the size of the neutron skin in (heavy) nuclei.

[Exp. from strongly interacting probes:  $\sim 0.15 - 0.22$  fm (*Physical Review C* **86** 015803 (2012))].

# The isobaric analog state energy: $E_{IAS}$



- **Analog state** can be defined:  $|A\rangle = \frac{T_-|0\rangle}{\langle 0|T_+T_-|0\rangle}$

- **Displacement energy or  $E_{IAS}$**

$$E_{IAS} = E_A - E_0 = \langle A|\mathcal{H}|A\rangle - \langle 0|\mathcal{H}|0\rangle = \frac{\langle 0|T_+[\mathcal{H}, T_-]|0\rangle}{\langle 0|T_+T_-|0\rangle}$$

$E_{IAS} \neq 0$  only due to Isospin Symmetry Breaking terms  $\mathcal{H}$   
 $E_{IAS}^{\text{exp}}$  usually accurately measured !

# Contributions

$[\mathcal{H}, T_-] \neq 0$  ? essentially **Coulomb potential** but not only

Table: Estimate of the different effects on  $E_{IAS}$  in  $^{208}\text{Pb}$ .

	$E_{IAS}$ Correction
Coumb direct	$\sim 20$ MeV
Coulomb exchange	$\sim -300$ keV
n-p mass difference	$\sim$ tens keV
Electromagnetic spin-orbit	$\sim -$ tens keV
Finite size effects	$\sim -$ 100 keV
Short range correlations	$\sim$ 100 keV
sospin impurity	$\sim -$ 100 keV
Isospin symmetry breaking	$\sim -$ 250 keV
	$\sim 19$ MeV

Physical Review Letters **23**, 484 (1969).

$E_{IAS}^{\text{exp}} = 18.83 \pm 0.01$  MeV. *Nuclear Data Sheets 108*, 1583 (2007).

## Coulomb direct contribution: very simple model

- Assuming independent particle model and good isospin for  $|0\rangle$   
( $\langle 0|T_+T_-|0\rangle = 2T_0 = N - Z$ )

$$E_{IAS} \approx E_{IAS}^{C,direct} = \frac{1}{N - Z} \int [\rho_n(\vec{r}) - \rho_p(\vec{r})] U_C^{direct}(\vec{r}) d\vec{r}$$

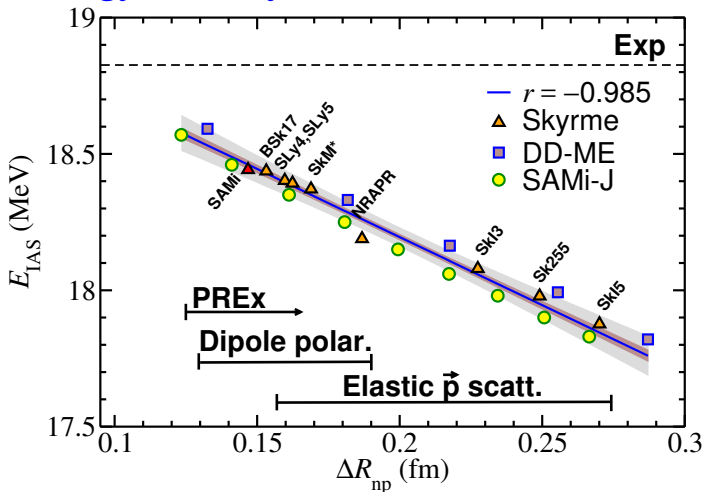
where  $U_C^{direct}(\vec{r}) = \int \frac{e^2}{|\vec{r}_1 - \vec{r}|} \rho_{ch}(\vec{r}_1) d\vec{r}_1$

- Assuming also a uniform neutron and proton distributions of radius  $R_n$  and  $R_p$  respectively, and  $\rho_{ch} \approx \rho_p$  one can find

$$E_{IAS} \approx E_{IAS}^{C,direct} \approx \frac{6}{5} \frac{Ze^2}{R_p} \left( 1 - \sqrt{\frac{5}{12}} \frac{N}{N - Z} \frac{\Delta r_{np}}{R_p} \right)$$

One may expect: **the larger the  $\Delta r_{np}$  the smallest  $E_{IAS}$**

# $E_{IAS}$ in Energy Density Functionals (No Corr.)



Phys. Rev. Lett. 120, 202501 (2018)

**Nuclear models (EDFs) where the nuclear part is isospin symmetric and  $U_{ch}$  is calculated from the  $\rho_p$**



## Corrections: within self-consistent HF+RPA

Within the **HF+RPA** one can **estimate** the  $E_{IAS}$  accounting (in an effective way) for **short-range correlations and effects of the continuum** (if a large sp base is adopted).

- **Coulomb exchange** exact (usually Slater approx.):

$$U_C^{x,\text{exact}} \varphi_i(\vec{r}) = -\frac{e^2}{2} \int d^3r' \frac{\varphi_j^*(\vec{r}') \varphi_j(\vec{r})}{|\vec{r} - \vec{r}'|} \varphi_i(\vec{r}')$$

- The **electromagnetic spin-orbit** correction to the nucleon single-particle energy (non-relativistic),

$$\varepsilon_i^{\text{emso}} = \frac{\hbar^2 c^2}{2m_i^2 c^4} \langle \vec{l}_i \cdot \vec{s}_i \rangle x_i \int \frac{1}{r} \frac{dU_C}{dr} |R_i(r)|^2$$

where  $x_i$ :  $g_p - 1$  for Z and  $g_n$  for N;  $g_n = -3.82608545(90)$  and  $g_p = 5.585694702(17)$ ,  $R_i \rightarrow R_{nl}$  radial wf.

## Corrections:

- **Finite size** effects (assuming spherical symmetry):

$$\begin{aligned}\rho_{\text{ch}}(q) &= \left(1 - \frac{q^2}{8m^2}\right) [G_{E,p}(q^2)\rho_p(q) + G_{E,n}(q^2)\rho_n(q)] \\ &- \frac{\pi q^2}{2m^2} \sum_{l,t} [2G_{M,t}(q^2) - G_{E,t}(q^2)] \langle \vec{l} \cdot \vec{s} \rangle \int_0^\infty dx \frac{j_1(qx)}{qx} |\mathcal{R}_{nl}(x)x^2|^2\end{aligned}$$

- **Vacuum polarization:** lowest order correction in the fine-structure constant to the Coulomb potential  $\frac{eZ}{r}$ :

$$V_{\text{vp}}(\vec{r}) = -\frac{2}{3} \frac{\alpha e^2}{\pi} \int d\vec{r}' \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} \mathcal{K}_1\left(\frac{2}{\lambda_e} |\vec{r} - \vec{r}'|\right)$$

where  $e$  is the fundamental electric charge,  $\alpha$  the fine-structure constant,  $\lambda_e$  the reduced Compton electron wavelength and

$$\mathcal{K}_1(x) \equiv \int_1^\infty dt e^{-xt} \left(\frac{1}{t^2} + \frac{1}{2t^4}\right) \sqrt{t^2 - 1}$$

# Corrections:

- **Isospin symmetry breaking** (Skyrme-like): **two parts** (contact interaction)

**charge symmetry breaking** +

$$V_{\text{CSB}} = V_{\text{nn}} - V_{\text{pp}}$$

$$V_{\text{CSB}}(\vec{r}_1, \vec{r}_2) \equiv \frac{1}{4} [\tau_z(1) + \tau_z(2)] s_0 (1 + y_0 P_\sigma)$$

$\tau_z$  Pauli in isospin space;  $P_\sigma$  are the usual projector operators in spin space.

**charge independence breaking\***

$$V_{\text{CIB}} = \frac{1}{2} (V_{\text{nn}} + V_{\text{pp}}) - V_{\text{pn}}$$

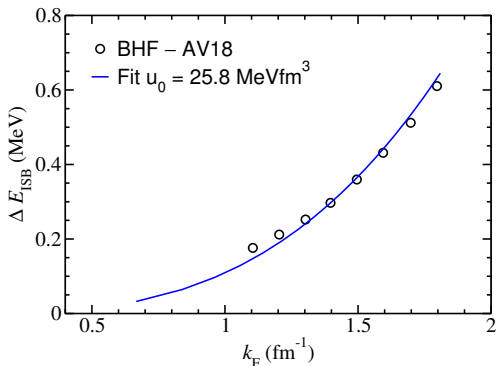
$$V_{\text{CIB}}(\vec{r}_1, \vec{r}_2) \equiv \frac{1}{2} \tau_z(1) \tau_z(2) u_0 (1 + z_0 P_\sigma)$$

\* general operator form  $\tau_z(1) \tau_z(2) - \frac{1}{3} \vec{\tau}(1) \cdot \vec{\tau}(2)$ . Our prescription  $\tau_z(1) \tau_z(2)$  not change structure of HF+RPA.

- Opposite to the other corrections, **ISB contributions depends on new parameters that need to be determined!**

## Isospin symmetry breaking in the medium:

- **keeping** things **simple**: **CSB** and **CIB** interaction just **delta function** depending on  $s_0$  and  $u_0$ . **Different possibilities**:
  - **Fitting** to (two) experimentally known **IAS energies**
  - **Derive from theory**
  - **our option**:  $u_0$  to reproduce **BHF** (symmetric nuclear matter) and  $s_0$  to reproduce  $E_{IAS}$  in  $^{208}\text{Pb}$



## Re-fit of SAMi: SAMi-ISB

- All these **corrections** are relatively **small** but **modify binding energies, neutron and proton distributions, etc.**  
⇒ a **re-fit of the interaction is needed.**
- Use **SAMi fitting protocol** (special care for spin-isospin resonances) including all corrections and **find SAMi-ISB**

Table: Saturation properties

	SAMi	SAMi-ISB	
$\rho_\infty$	0.159(1)	0.1613(6)	$\text{fm}^{-3}$
$e_\infty$	-15.93(9)	-16.03(2)	MeV
$m_{\text{IS}}^*$	0.6752(3)	0.730(19)	
$m_{\text{IV}}^*$	0.664(13)	0.667(120)	
J	28(1)	30.8(4)	MeV
L	44(7)	50(4)	MeV
$K_\infty$	245(1)	235(4)	MeV

## SAMi-ISB finite nuclei properties

El.	N	B [MeV]	B <sup>exp</sup> [MeV]	r <sub>c</sub> [fm]	r <sub>c</sub> <sup>exp</sup> [fm]	ΔR <sub>np</sub> [fm]
Ca	28	417.67	415.99	3.49	3.47	0.214
Zr	50	783.60	783.89	4.26	4.27	0.097
Sn	82	1102.75	1102.85	4.73	–	0.217
Pb	126	1635.78	1636.43	5.50	5.50	0.151

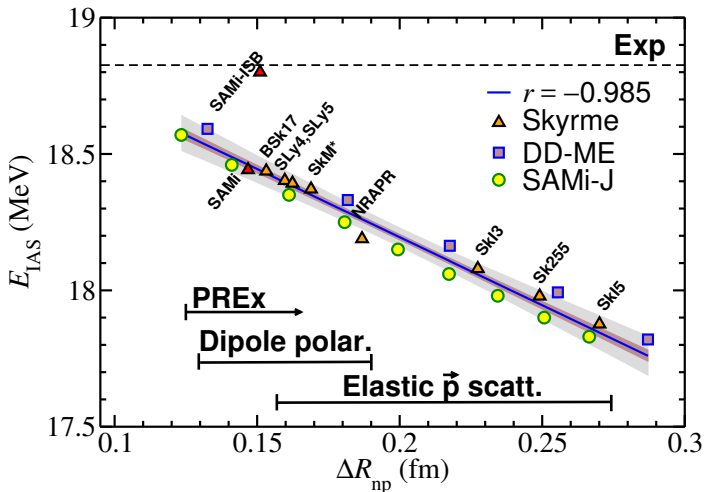
### Corrections on E<sub>IAS</sub> for <sup>208</sup>Pb one by one

	E <sub>IAS</sub> [MeV]	Correction [keV]
No corrections <sup>a</sup>	18.31	
Exact Coulomb exchange	18.41	+100
n/p mass difference	18.44	+30
Electromagnetic spin-orbit	18.45	+10
Finite size effects	18.40	-50
Vacuum polarization (V <sub>ch</sub> )	18.53	+130
Isospin symmetry breaking	<b>18.80(5)</b>	+270

<sup>a</sup>From Skyrme Hamiltonian where the nuclear part is isospin symmetric and V<sub>ch</sub> is calculated from the ρ<sub>p</sub>

$$E_{IAS}^{\text{exp}} = 18.83 \pm 0.01 \text{ MeV. } \textit{Nuclear Data Sheets 108, 1583 (2007).}$$

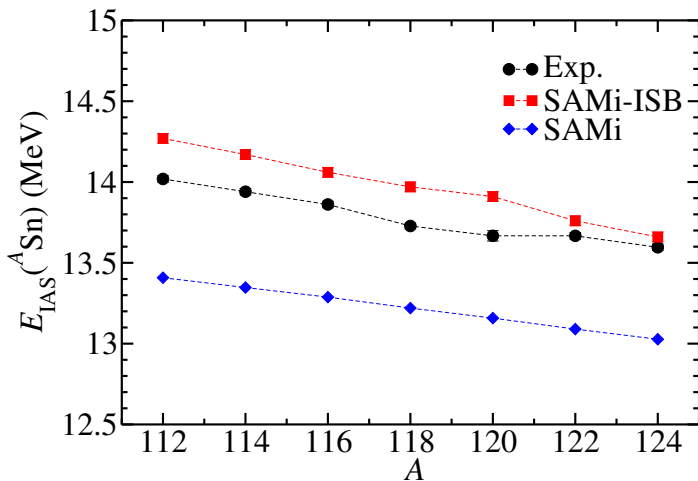
# $E_{IAS}$ with SAMi-ISB



Phys. Rev. Lett. 120, 202501 (2018)

**Measurement of  $\Delta r_{np}$   $\rightarrow$  determine ISB in the nuclear medium**

## Prediction: $E_{IAS}$ in the Sn isotopic chain



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# Conclusions

- The most **promising observables** to constraint the symmetry energy are those that can be **measured** via processes with little or **no direct influence from the strong force** (e.g. electromagnetic or weak probes).
- Alternatively,  $E_{IAS}$  has **no** dependence on the **isospin conserving** part (largest) of the **strong interaction**.
- **EDFs** of common use in nuclear physics show a **linear dependence between  $E_{IAS}$  and  $\Delta r_{np}$**
- **EDFs do not properly** describe the experimental  $E_{IAS}$
- **Modification of  $\mathcal{H}_{eff}$  requires a refit** of the interaction including **new ISB parameters**.
- **One can reconcile good reproduction of experimental charge radii, binding energies,  $E_{IAS}$ ...**

- **A better knowledge of ISB contributions in the medium may lead to an accurate determination of neutron skin thickness via  $E_{IAS}$  (or the other way around)**

**Thank you for your  
attention!**