

Energy Loss of (heavy)
charged particles by
atomic collision

Energy LOSS

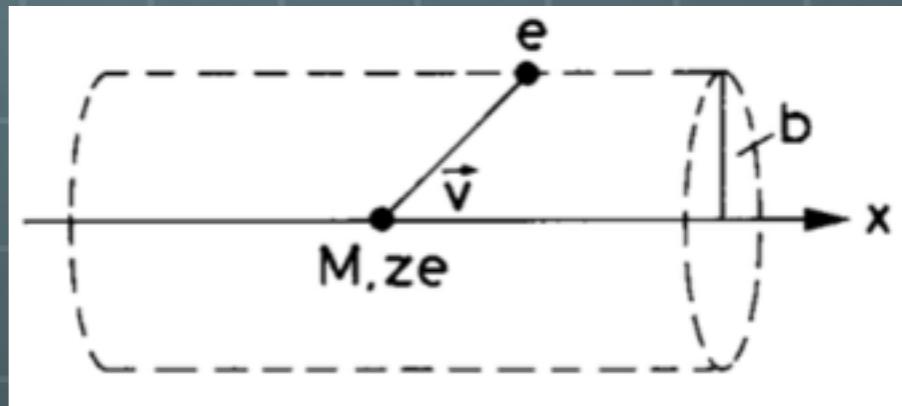


Charge ze , velocity v , mass M , Impact parameter b



Momentum Transfer

$$\begin{aligned}\Delta P_T &= \int F dt = \int F_T dt = \int F_T \frac{dx}{v} \\ &= \int \frac{ze^2}{(x^2 + b^2)} \cdot \frac{b}{\sqrt{x^2 + b^2}} \cdot \frac{1}{v} dx \\ &= \frac{2ze^2}{l}\end{aligned}$$



Consider a charged particle with a charge ze , velocity v , and mass M interacting with an atomic electron at distance b .

Energy Loss (cont.)

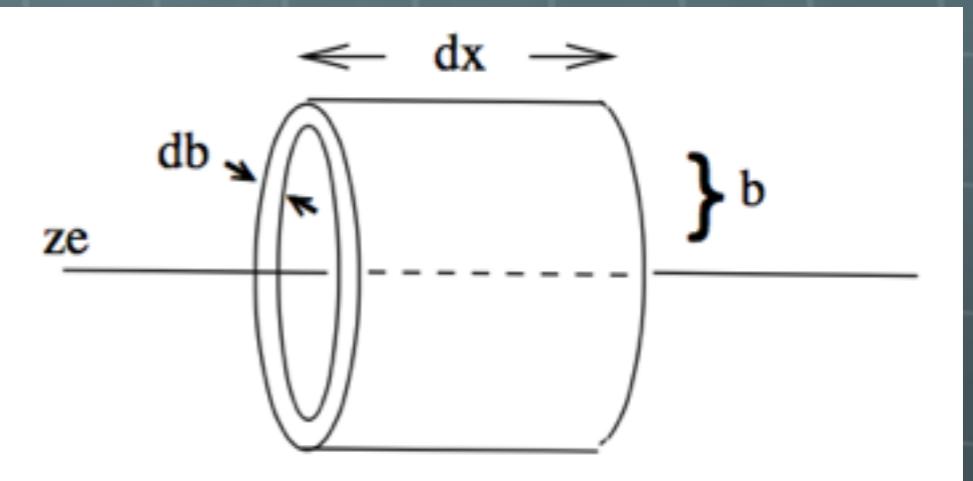


Consider a bulk of electrons



Number of electrons

$$N_e = \rho \cdot 2\pi b \cdot db \ dx$$



Energy loss becomes

$$dE(b) = N_e \cdot \Delta E(b) = \frac{4\pi z^2 e^4}{m_e v^2} \rho \frac{db}{b} dx$$

$$\frac{dE}{dx} = \frac{4\pi z^2 e^4}{m_e v^2} \rho \int \frac{db}{b} \rightarrow \frac{4\pi z^2 e^4}{m_e v^2} \rho \cdot \ln\left(\frac{b_{max}}{b_{min}}\right)$$

Bohr's Formula

b_{min} : head-on collision

$$2m_e v^2 = \frac{ze^2}{b_{min}} \rightarrow b_{min} = \frac{ze^2}{2m_e v^2}$$

b_{max} : adiabatic invariance

$$F \cdot \tau = \Delta P \quad \tau : \text{Period of bound electron's orbital motion}$$

$$\frac{ze^2 \tau}{b_{max}^2} = \frac{2ze^2}{b_{max} v} \rightarrow b_{max} = \frac{v \tau}{2} = \frac{v}{2\nu}$$

$$\left\langle -\frac{dE}{dx} \right\rangle = \frac{4\pi z^2 e^4}{m_e v^2} N_A \frac{Z}{A} \ln\left(\frac{mv^3}{\nu ze^2}\right)$$



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The Review of Particle Physics (2017)

C. Patrignani *et al.* (Particle Data Group), Chin. Phys. C, **40**, 100001 (2016) and 2017 update.



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33. PASSAGE OF PARTICLES THROUGH MATTER

Revised August 2015 by H. Bichsel (University of Washington), D.E. Groom (LBNL), and S.R. Klein (LBNL).

Bethe Equation

$$\left\langle -\frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$

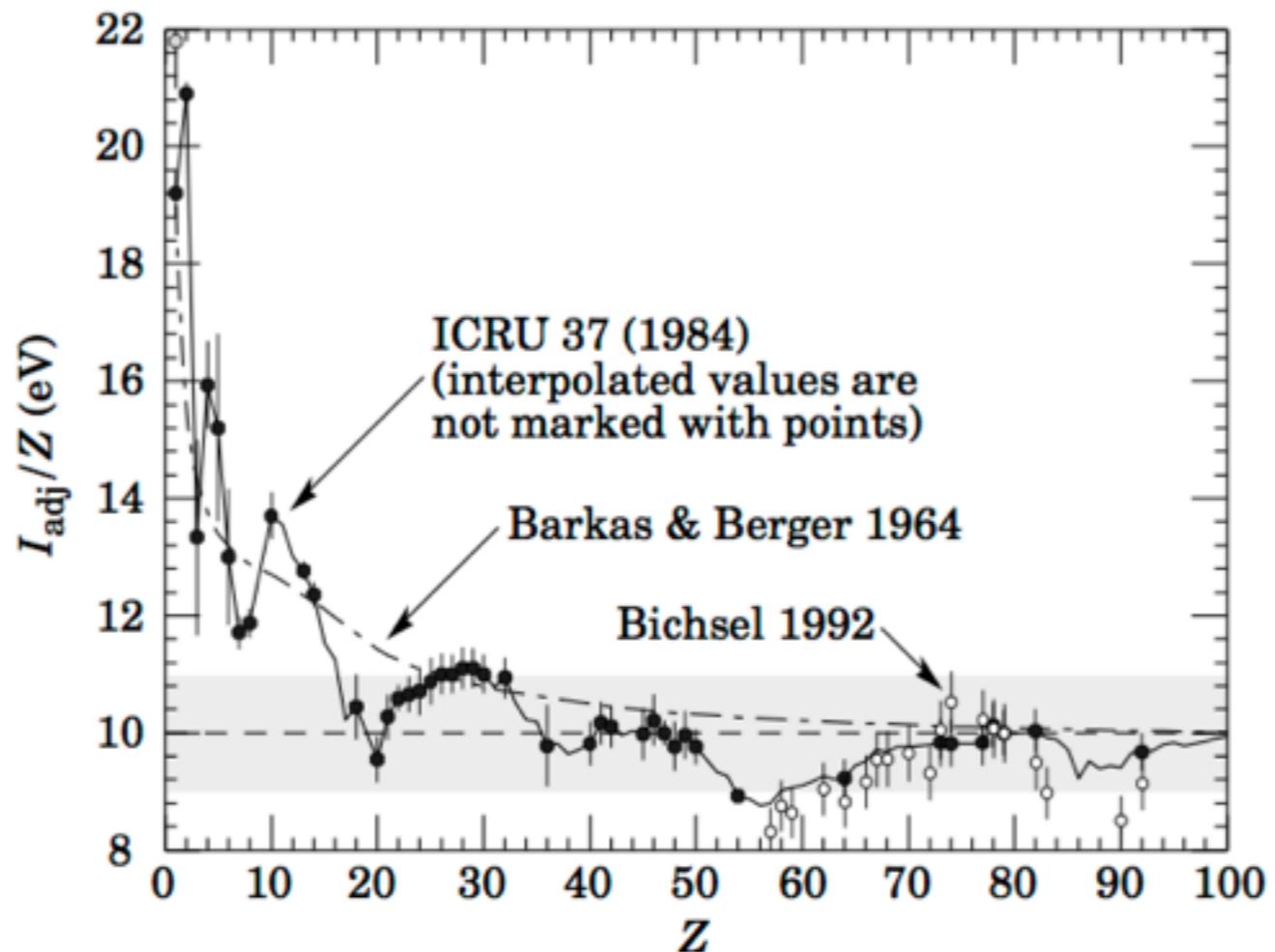
$$K = 4\pi N_A r_e^2 m_e c^2 = 0.307075 \text{ MeV mol}^{-1} \text{ cm}^2$$

$$r_e = e^2 / (4\pi\epsilon_0 m_e c^2) = 2.8179403267(27) \text{ fm}$$

$$W_{\max} = \frac{2m_e c^2 \beta^2 \gamma^2}{1 + 2\gamma m_e/M + (m_e/M)^2}$$

$$\left\langle -\frac{dE}{dx} \right\rangle = \frac{4\pi z^2 e^4}{m_e v^2} N_A \frac{Z}{A} \ln \left(\frac{mv^3}{4\nu z e^2} \right)$$

Mean Excitation Energy



ICRU : International Commission of Radiation Units & Measurements

I : average excitation potential of atoms

Bloch's law : $I = 10^*Z$

실제로 에너지를 전달한다는 것은 물질의 원자의 excitation energy에 관련하는 것이 되고, 초기의 실험에서는 이 Stopping power를 측정하는 실험이 많이 이루어졌으며, 이 결과로부터, 원자의 평균적인 Excitation potential 즉 mean excitation energy를 유도해 낸다.

먼저는 $I/Z=10$ Bloch's law 라고 하는 데, 내 눈으로 확인한 것이 아니라 불안함. 이 relation이 LOW-Z material에서 벗어나고, 이 correction을 Barkas correction. second-order correction이 Bloch임.

Atomic and nuclear properties of copper (Cu)

Quantity	Value	Units	Value	Units
Atomic number	29			
Atomic mass	63.546(3)	g mole ⁻¹		
Specific gravity	8.960	g cm ⁻³		
Mean excitation energy	322.0	eV		
Minimum ionization	1.403	MeV g ⁻¹ cm ²	12.57	MeV cm ⁻¹
Nuclear collision length	84.2	g cm ⁻²	9.393	cm
Nuclear interaction length	137.3	g cm ⁻²	15.32	cm
Pion collision length	109.3	g cm ⁻²	12.20	cm
Pion interaction length	165.9	g cm ⁻²	18.51	cm
Radiation length	12.86	g cm ⁻²	1.436	cm
<u>Critical energy</u>	19.42	MeV (for e^-)	18.79	MeV (for e^+)
Molière radius	14.05	g cm ⁻²	1.568	cm
Plasma energy $\hbar\omega_p$	58.27	eV		
Muon critical energy	317.	GeV		
Melting point	1358.	K	1085.	C
Boiling point @ 1 atm	2835.	K	2562.	C

For muons, $dE/dx = a(E) + b(E) E$. Tables of $b(E)$: [PDF](#) [TEXT](#)

Table of muon dE/dx and Range: [PDF](#) [TEXT](#)

Explanation of some entries

[Table of isotopes](#) Warning: may not be current

x ray mass attenuation coefficients

Materials

>) tables including radiative losses for muons, nuclear and pion
ius, plasma energy, and links to isotope and x-ray mass attenuation

Materials

33. PASSAGE OF PARTICLES THROUGH MATTER

Revised August 2015 by H. Bichsel (University of Washington), D.E. Groom (LBNL), and S.R. Klein (LBNL).

Bethe Equation

$$\left\langle -\frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$



charge square of incident particle (z^2)



charge of matter (Z)

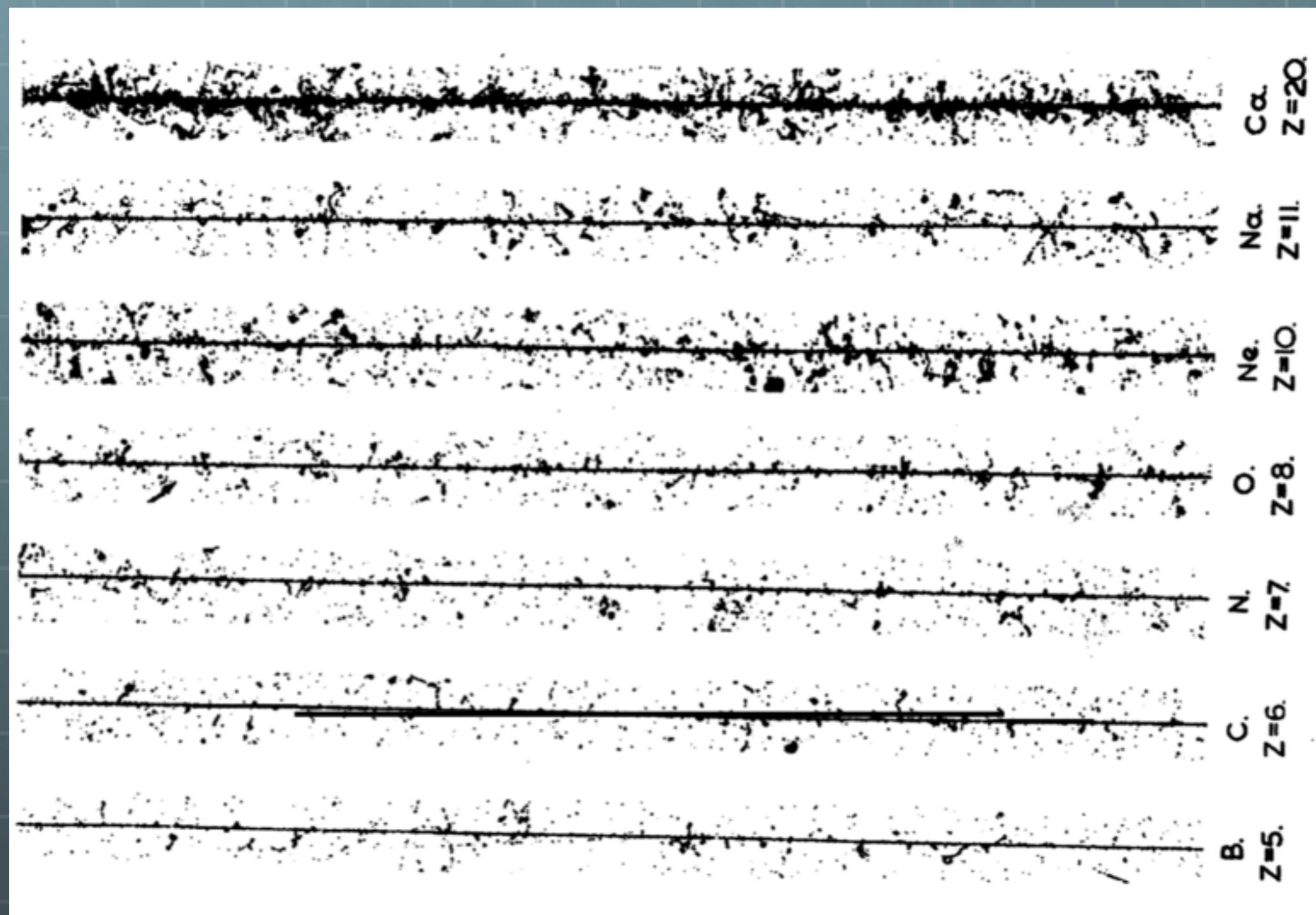


Inverse of mass of matter ($1/A$)



Unit of $\text{MeV} \cdot \text{cm}^2/\text{g}$

일반적인 방정식에서와 같이, 하나의 sacle을 나타내주는 Constant K로 표현하고 있고, 단위 길이당 평균적으로 잃어 버리는 에너지는 입사하는 입자의 전하의 제곱과 물질의 전하 Z에 비례하는 값이 됨을 알 수 있습니다. 물질의 Mass Number에 반비례하는 표현으로 나타나고 있는데, 이는 이 표현식이 $\text{MeV} \cdot \text{cm}^{-2}/\text{g}$ 의 단위를 갖게 함으로써, 좀 더 보편적인 표현식이 되도록 한 것입니다.



Cecil F. Powell Nobel Lecture 1950

33. PASSAGE OF PARTICLES THROUGH MATTER

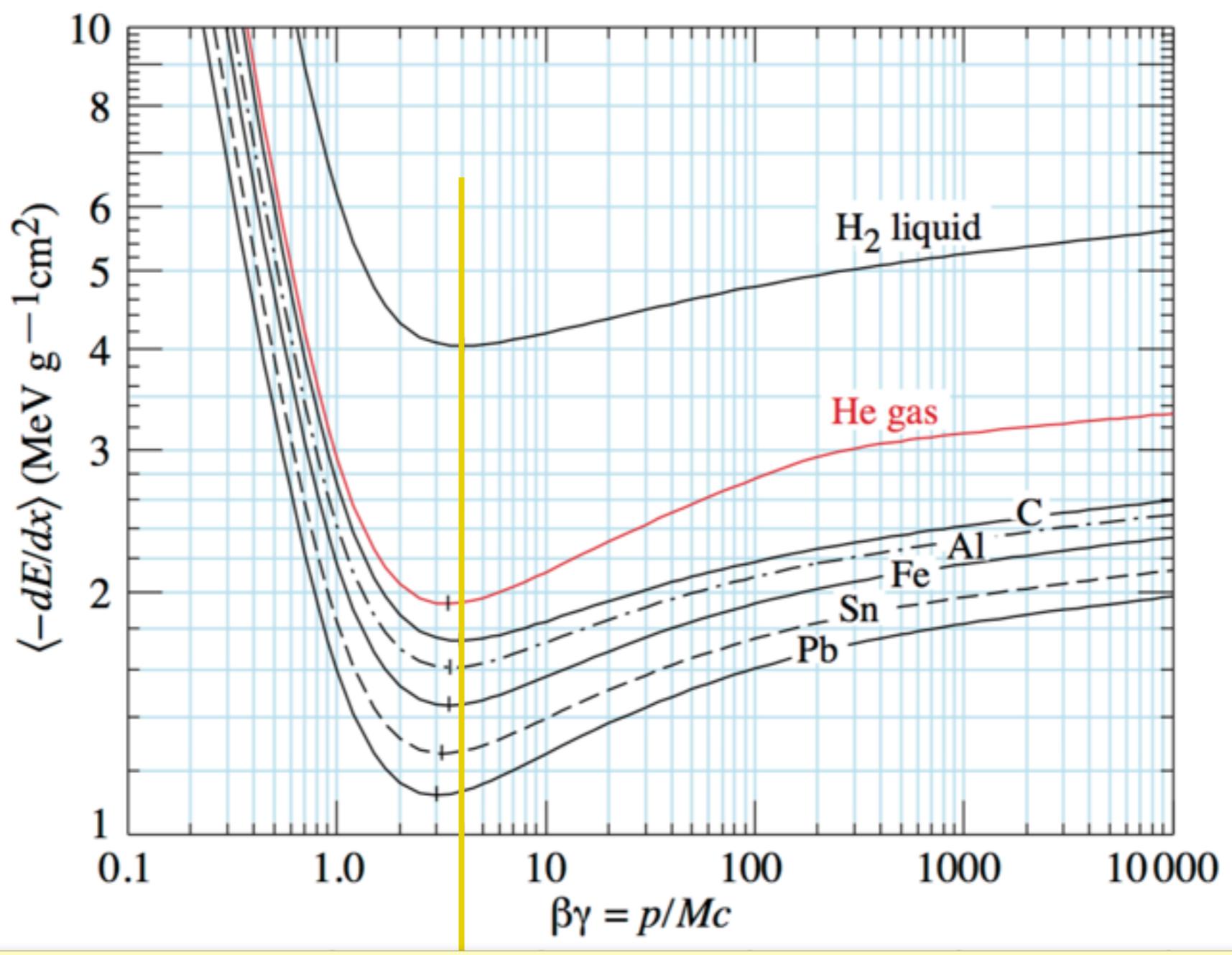
Revised August 2015 by H. Bichsel (University of Washington), D.E. Groom (LBNL), and S.R. Klein (LBNL).

Bethe Equation

$$\left\langle -\frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$

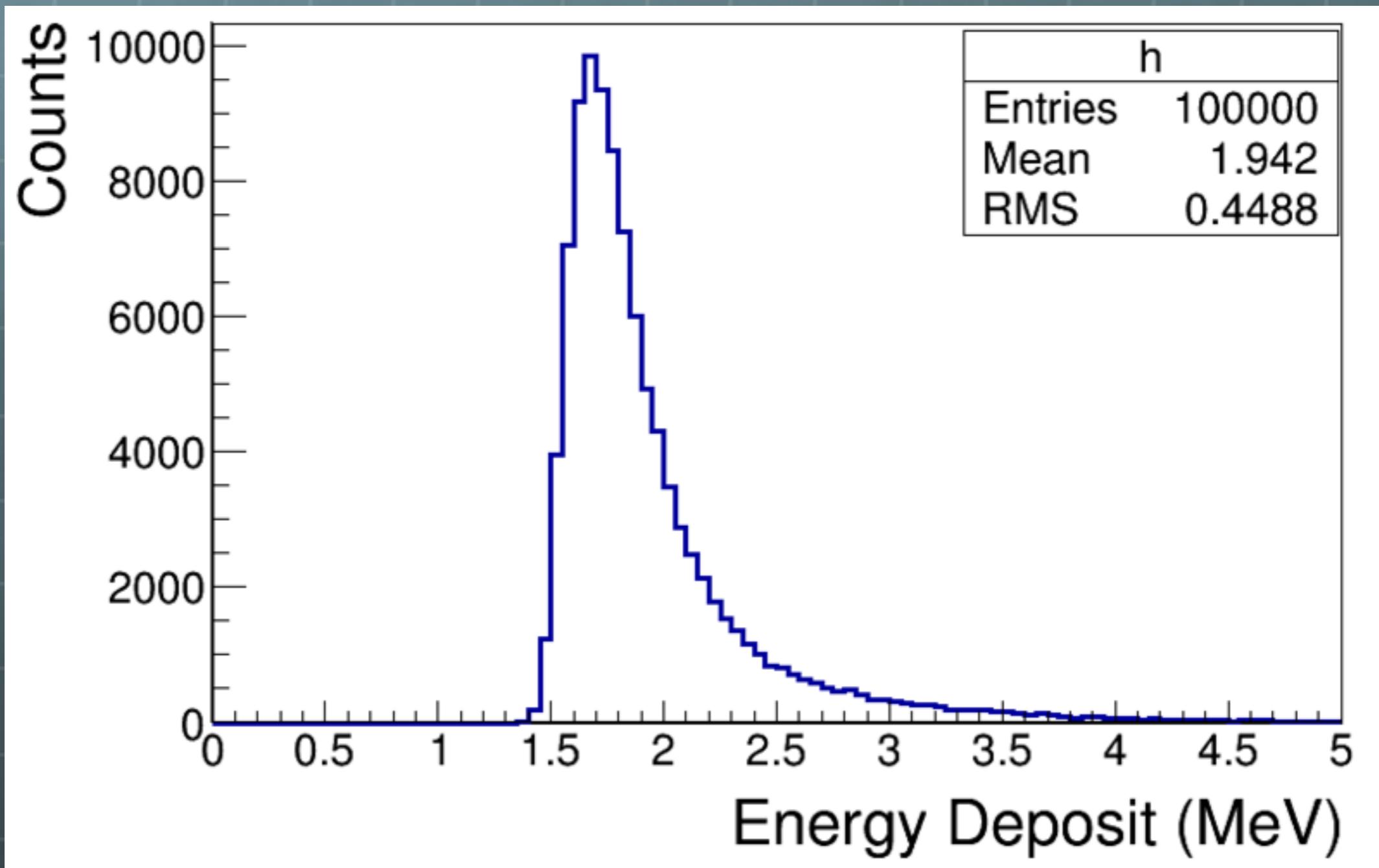
- charge square of incident particle (z^2)
- charge of matter (Z)
- Inverse of mass of matter ($1/A$)
- Mixture of different dependence on the velocity (β)
- Maximum energy transfer in a single collision

$$W_{\max} = \frac{2m_e c^2 \beta^2 \gamma^2}{1 + 2\gamma m_e/M + (m_e/M)^2}$$



이 그림에서 보는 바와 같이, 넓은 영역에 있어서, Mass stopping power는 Mass number로 normalize한 energy loss식으로 나타낼 수 있게 되는 것입니다. 이 그림에서 보면 $\beta\gamma$ 가 3에서 3.5사이의 영역에서 가장 작은 값을 갖게 되는데, 이 운동량을 갖는 입자를 minimum ionizing particle이라고 합니다. 이 값이 $\beta\gamma$ 가 갖는 값을 갖는 것이 되니까, 입사하는 입자의 운동량은 질량에 따라서 큰 차이를 나타내게 되는데.... 가령 muon의 경우에는 약 400 MeV/c의 운동량을 갖는 경우가 MIP이 되지만, proton의 경우가 되면 약 3GeV/c의 운동량을 갖는 경우가 MIP가 된다. 실제 실험실에서 Cosmic muon (엄밀히 말하면 각 지점에 따라서 약간의 차이가 있고 (대기중에서 만들어 져서 실험실까지 들어오니까....) (여러분의 실험실에 들어오는 Cosmic muon의 energy spectrum을 재보는 것도 재미있는 놀이가 될지도 모르겠네요.....). 에너지 분포를 갖고 있으니까, 이 cosmic muon을 MIP이라고 하는 것은 약간은 다른 의미가 되겠지만, 300 MeV/c을 넘어서면 그다지 크게 변하지 않기 때문에, 광범위한 의미에서 Cosmic muon을 MIP으로 취급하고, cosmic muon이 남기는 에너지를 하나의 값으로 나타내기도 합니다.

MIP energy deposit

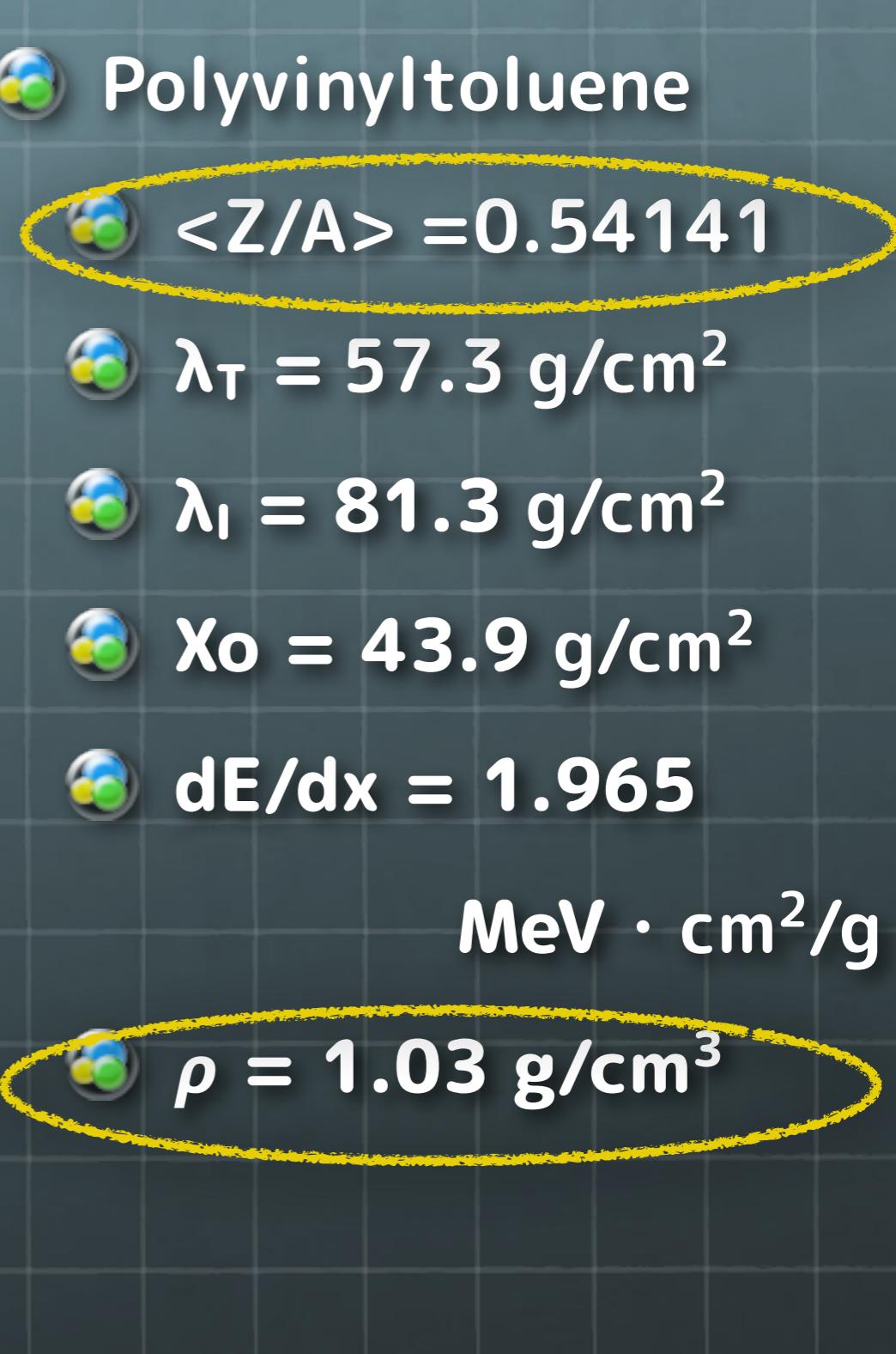


중심으로하는 gaussian 으로 나타나겠지만, 지금 이 그림에서 보는 바와같이 비대칭적인 분포를 갖게된다.

6. ATOMIC AND NUCLEAR PROPERTIES OF MATERIALS

Table 6.1 Abridged from pdg.lbl.gov/AtomicNuclearProperties by D.E. Groom (2015). See web pages for more detail about entries in this table and for several hundred other substances. Parentheses in the dE/dx and density columns indicate gases at 20°C and 1 atm. Boiling points are at 1 atm. Refractive indices n are evaluated at the sodium D line blend (589.2 nm); values $\gg 1$ in brackets indicate $(n - 1) \times 10^6$ for gases at 0°C and 1 atm.

Material	Z	A	$\langle Z/A \rangle$	Nucl.coll.	Nucl.inter.	Rad.len.	$dE/dx _{\text{min}}$	Density	Melting	Boiling	Refract.
				length λ_T [g cm $^{-2}$]	length λ_I [g cm $^{-2}$]	X $_0$ [g cm $^{-2}$]	{ MeV $g^{-1} \text{cm}^2$ }	{ g cm $^{-3}$ ({g t $^{-1}$ })	point (K)	point (K)	@ Na D
H ₂	1	1.008(7)	0.99212	42.8	52.0	63.04	(4.103)	0.071(0.084)	13.81	20.28	1.11[132.]
D ₂	1	2.01410177803(8)	0.49650	51.3	71.8	125.97	(2.053)	0.169(0.168)	18.7	23.65	1.11[138.]
He	2	4.002602(2)	0.49967	51.8	71.0	94.32	(1.937)	0.125(0.166)		4.220	1.02[35.0]
Li	3	6.94(2)	0.43221	52.2	71.3	82.78	1.639	0.534	453.6	1615.	
Be	4	9.0121831(5)	0.44384	55.3	77.8	65.19	1.595	1.848	1560.	2744.	
C diamond	6	12.0107(8)	0.49955	59.2	85.8	42.70	1.725	3.520			2.42
C graphite	6	12.0107(8)	0.49955	59.2	85.8	42.70	1.742	2.210			
N ₂	7	14.007(2)	0.49976	61.1	89.7	37.99	(1.825)	0.807(1.165)	63.15	77.29	1.20[298.]
O ₂	8	15.999(3)	0.50002	61.3	90.2	34.24	(1.801)	1.141(1.332)	54.36	90.20	1.22[271.]
F ₂	9	18.998403163(6)	0.47372	65.0	97.4	32.93	(1.676)	1.507(1.580)	53.53	85.03	[195.]
Ne	10	20.1797(6)	0.49555	65.7	99.0	28.93	(1.724)	1.204(0.839)	24.56	27.07	1.09[67.1]
Al	13	26.9815385(7)	0.48181	69.7	107.2	24.01	1.615	2.699	933.5	2792.	
Si	14	28.0855(3)	0.49848	70.2	108.4	21.82	1.664	2.329	1687.	3538.	3.95
Cl ₂	17	35.453(2)	0.47951	73.8	115.7	19.28	(1.630)	1.574(2.980)	171.6	239.1	[773.]
Ar	18	39.948(1)	0.45059	75.7	119.7	19.55	(1.519)	1.396(1.662)	83.81	87.26	1.23[281.]
Ti	22	47.867(1)	0.45961	78.8	126.2	16.16	1.477	4.540	1941.	3560.	
Fe	26	55.845(2)	0.46557	81.7	132.1	13.84	1.451	7.874	1811.	3134.	
Cu	29	63.546(3)	0.45636	84.2	137.3	12.86	1.403	8.960	1358.	2835.	
Ge	32	72.630(1)	0.44053	86.9	143.0	12.25	1.370	5.323	1211.	3106.	
Sn	50	118.710(7)	0.42119	98.2	166.7	8.82	1.263	7.310	505.1	2875.	
Xe	54	131.293(6)	0.41129	100.8	172.1	8.48	(1.255)	2.953(5.483)	161.4	165.1	1.39[701.]
W	74	183.84(1)	0.40252	110.4	191.9	6.76	1.145	19.300	3695.	5828.	
Pt	78	195.084(9)	0.39983	112.2	195.7	6.54	1.128	21.450	2042.	4098.	
Au	79	196.966569(5)	0.40108	112.5	196.3	6.46	1.134	19.320	1337.	3129.	
Pb	82	207.2(1)	0.39575	114.1	199.6	6.37	1.122	11.350	600.6	2022.	
U	92	[238.02891(3)]	0.38651	118.6	209.0	6.00	1.081	18.950	1408.	4404.	
Air (dry, 1 atm)			0.49919	61.3	90.1	36.62	(1.815)	(1.205)		78.80	[289]
Shielding concrete			0.50274	65.1	97.5	26.57	1.711	2.300			
Borosilicate glass (Pyrex)			0.49707	64.6	96.5	28.17	1.696	2.230			
Lead glass			0.42101	95.9	158.0	7.87	1.255	6.220			
Standard rock			0.50000	66.8	101.3	26.54	1.688	2.650			
Methane (CH ₄)			0.62334	54.0	73.8	46.47	(2.417)	(0.667)	90.68	111.7	[444.]
Ethane (C ₂ H ₆)			0.59861	55.0	75.9	45.66	(2.304)	(1.263)	90.36	184.5	
Propane (C ₃ H ₈)			0.58962	55.3	76.7	45.37	(2.262)	0.493(1.868)	85.52	231.0	
Butane (C ₄ H ₁₀)			0.59497	55.5	77.1	45.23	(2.278)	(2.489)	134.9	272.6	
Octane (C ₈ H ₁₈)			0.57778	55.8	77.8	45.00	2.123	0.703	214.4	398.8	
Paraffin (CH ₃ (CH ₂) _{n=23} CH ₃)			0.57275	56.0	78.3	44.85	2.088	0.930			
Nylon (type 6, 6/6)			0.54790	57.5	81.6	41.92	1.973	1.18			
Polycarbonate (Lexan)			0.52697	58.3	83.6	41.50	1.886	1.20			
Polyethylene ([CH ₂ CH ₂] _n)			0.57034	56.1	78.5	44.77	2.079	0.89			
Polyethylene terephthalate (Mylar)			0.52037	58.9	84.9	39.95	1.848	1.40			
Polyimide film (Kapton)			0.51264	59.2	85.5	40.58	1.820	1.42			
Polymethylmethacrylate (acrylic)			0.53937	58.1	82.8	40.55	1.929	1.19			1.49
Polypropylene			0.55998	56.1	78.5	44.77	2.041	0.90			1.59
Polystyrene ([C ₆ H ₅ CHCH ₂] _n)			0.53768	57.5	81.7	43.79	1.936	1.06			
Polytetrafluoroethylene (Teflon)			0.47992	63.5	94.4	34.84	1.671	2.20			
Polyvinyltoluene			0.54141	57.3	81.3	43.90	1.956	1.03			1.58
Aluminum oxide (sapphire)			0.49038	65.5	98.4	27.94	1.647	3.970	2327.	3273.	1.77
Barium fluoride (BaF ₂)			0.42207	90.8	149.0	9.91	1.303	4.893	1641.	2533.	1.47
Bismuth germanate (BGO)			0.42065	96.2	159.1	7.97	1.251	7.130	1317.		2.15
Carbon dioxide gas (CO ₂)			0.49989	60.7	88.9	36.20	1.819	(1.842)			[449.]
Solid carbon dioxide (dry ice)			0.49989	60.7	88.9	36.20	1.787	1.563	Sublimes at 194.7 K		
Cesium iodide (CsI)			0.41569	100.6	171.5	8.39	1.243	4.510	894.2	1553.	1.79
Lithium fluoride (LiF)			0.46262	61.0	88.7	39.26	1.614	2.635	1121.	1946.	1.39
Lithium hydride (LiH)			0.50321	50.8	68.1	79.62	1.897	0.820	965.		
Lead tungstate (PbWO ₄)			0.41315	100.6	168.3	7.39	1.229	8.300	1403.		2.20
Silicon dioxide (SiO ₂ , fused quartz)			0.49930	65.2	97.8	27.05	1.699	2.200	1986.	3223.	1.46
Sodium chloride (NaCl)			0.47910	71.2	110.1	21.91	1.847	2.170	1075.	1738.	1.54
Sodium iodide (NaI)			0.42697	93.1	154.6	9.49	1.305	3.667	933.2	1577.	1.77
Water (H ₂ O)			0.55509	58.5	83.3	36.08	1.992	1.000	273.1	373.1	1.33
Silica aerogel			0.50093	65.0	97.3	27.25	1.740	0.200	(0.03 H ₂ O, 0.97 SiO ₂)		



In case of 400 MeV/c muon entering the plastic scintillator

$$\left\langle -\frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$

$K = 0.307075 \text{ MeV} \cdot \text{cm}^2/\text{mol}$

$\langle Z/A \rangle = 0.54141$

$\beta = 0.9668, \beta\gamma = 3.7858$

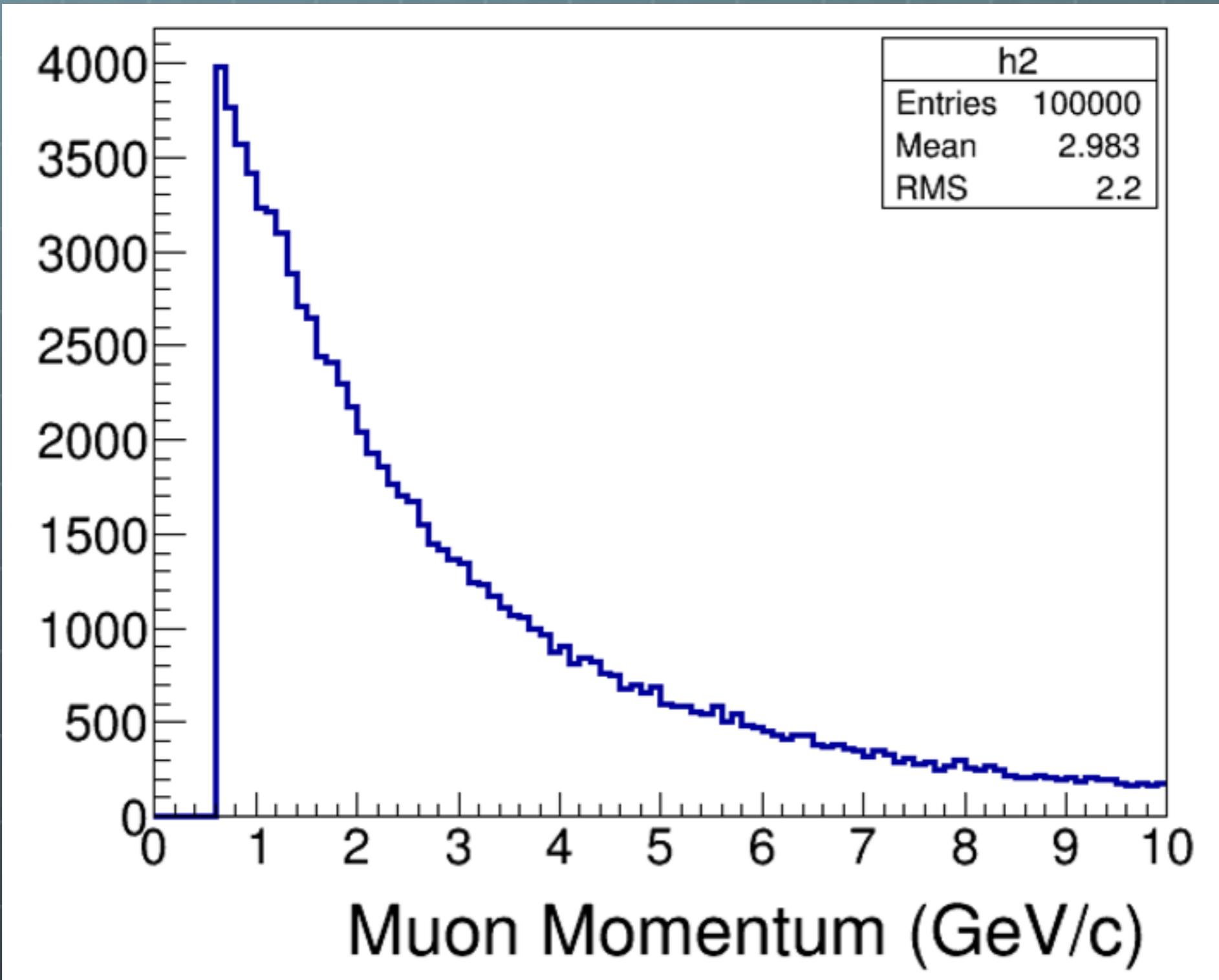
$I = 64.7 \text{ eV}$

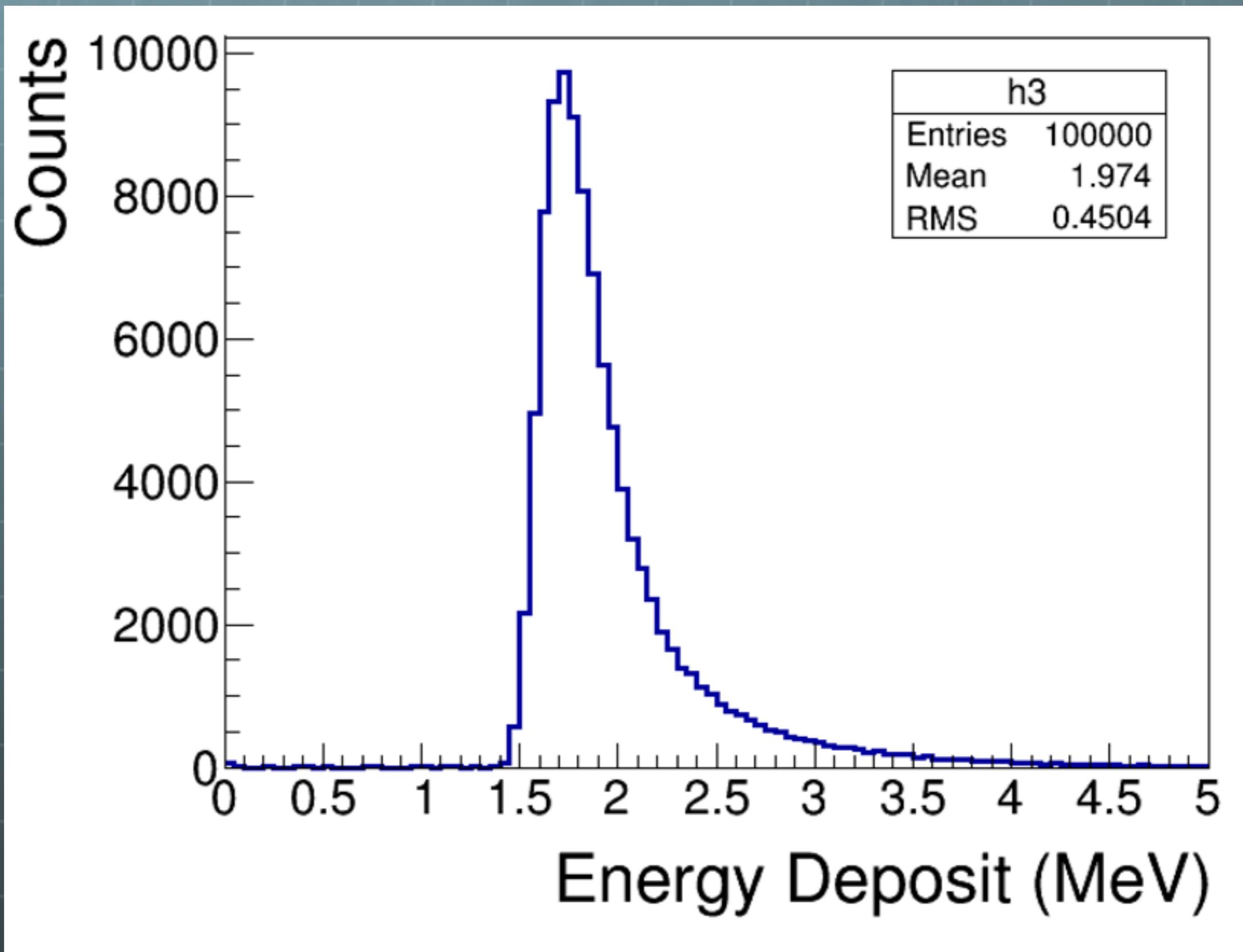
$m_e c^2 = 0.511 \text{ MeV}$

$W_{\max} = 14.109 \text{ MeV}$

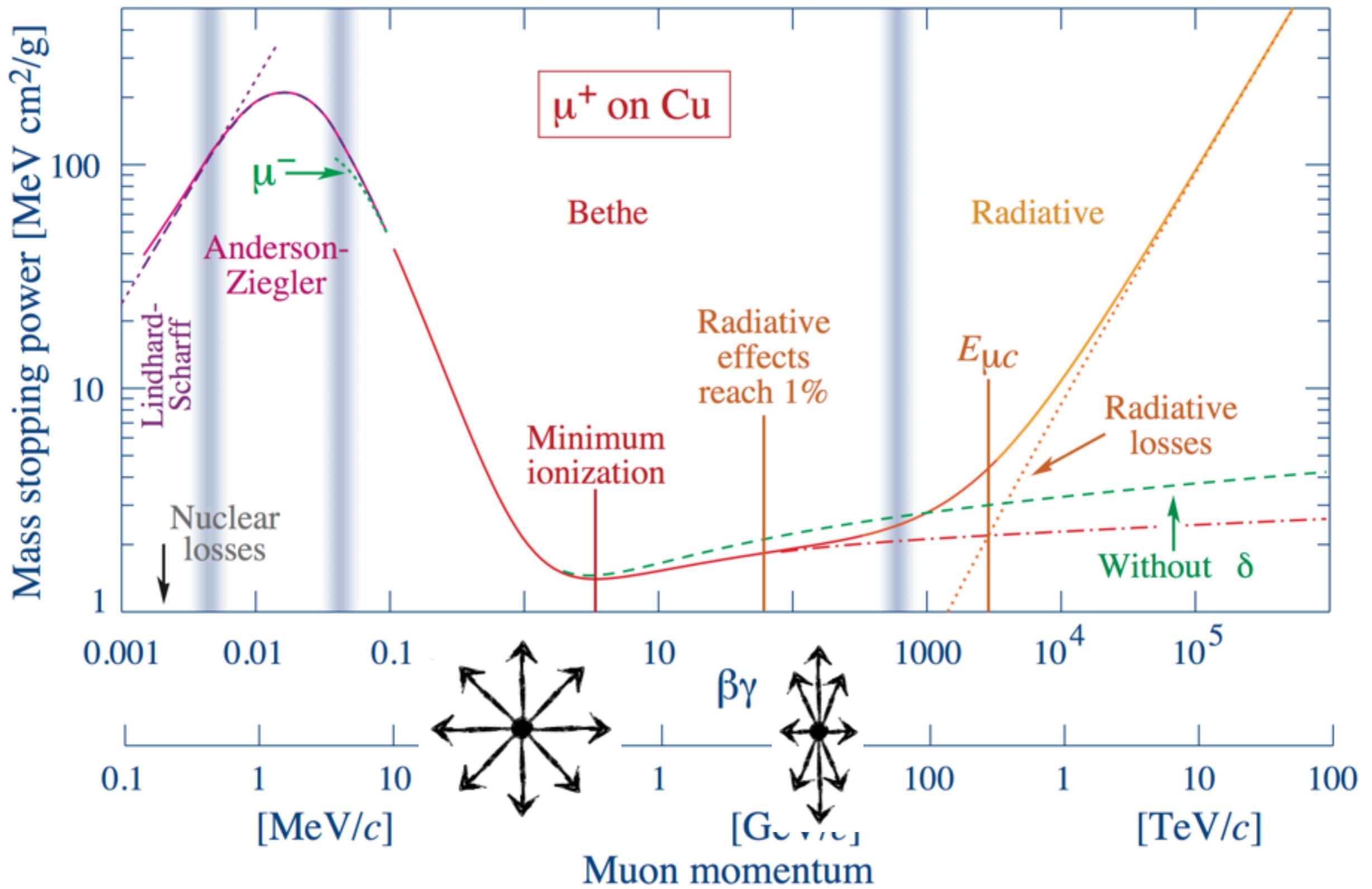
Quantity	Value	Units	Value	Units
$\langle Z/A \rangle$	0.54141			
Specific gravity	1.032	g cm^{-3}		
Mean excitation energy	64.7	eV		
Minimum ionization	1.956	$\text{MeV g}^{-1}\text{cm}^2$	2.019	MeV cm^{-1}
Nuclear collision length	57.3	g cm^{-2}	55.56	cm
Nuclear interaction length	81.3	g cm^{-2}	78.80	cm
Pion collision length	84.8	g cm^{-2}	82.19	cm
Pion interaction length	113.3	g cm^{-2}	109.8	cm
Radiation length	43.90	g cm^{-2}	42.54	cm
Critical energy	94.11	MeV (for e^-)	91.62	MeV (for e^+)
Molière radius	9.89	g cm^{-2}	9.586	cm
Plasma energy $\hbar\omega_p$	21.54	eV		
Muon critical energy	1195.	GeV		
Index of refraction (Na D)	1.580			

Cosmic ray momentum spectrum @ KEK

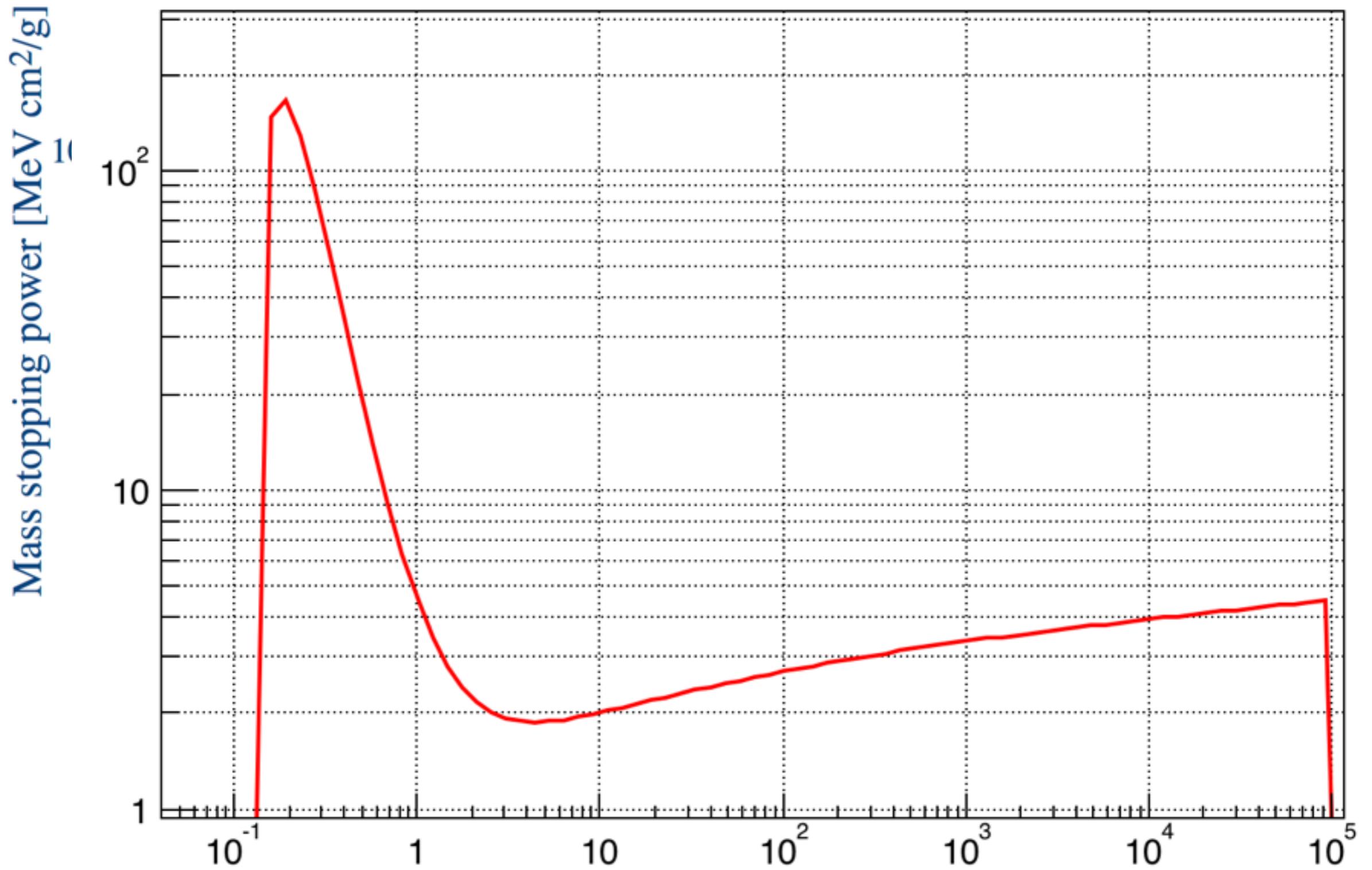




$$\left\langle -\frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$



$$\left\langle -\frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$



Particle Energy Deposit

$\beta\gamma > 3.5$:

$$\left\langle \frac{dE}{dx} \right\rangle \approx \frac{dE}{dx} \Big|_{\min}$$

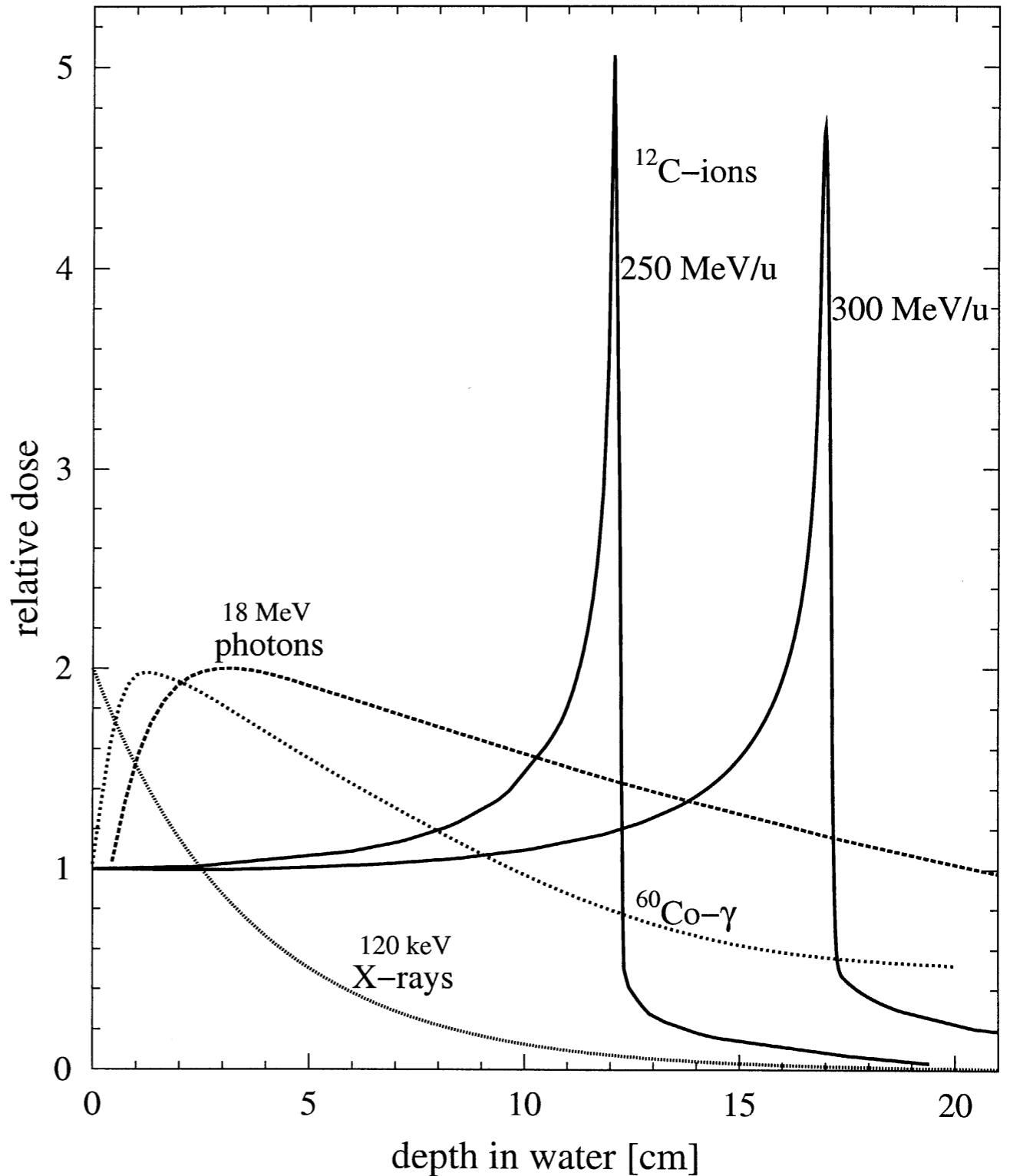
$\beta\gamma < 3.5$:

$$\left\langle \frac{dE}{dx} \right\rangle \gg \frac{dE}{dx} \Big|_{\min}$$

Applications:

Tumor therapy

Possibility to precisely deposit dose
at well defined depth by E_{beam} variation
[see Journal Club]



Straggling -Landau Distribution-

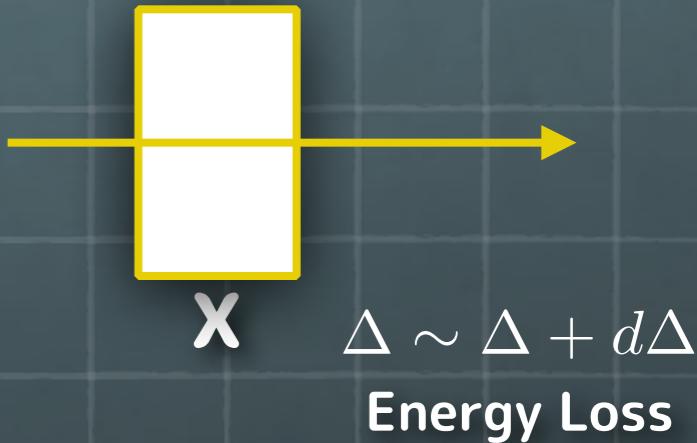


For a small energy loss, prob. of fluctuations



Let unknown function $f(x, \Delta)$

$$\frac{\partial f}{\partial x} = \int_0^\infty \omega(\varepsilon) [f(x, \Delta - \varepsilon) - f(x, \Delta)] d\varepsilon$$



$$f(x, \Delta) = \frac{1}{2\pi i} \int_{i\infty+\sigma}^{+i\infty+\sigma} e^{p\Delta} \phi(p, x) dp$$

두께 x 의 매질을 통과하는 하전 입자가 $\Delta \sim \Delta + d\Delta$ 사이의 에너지를 남길 확률을 $f(x, D)$ 라고 하자. 이 unknown 함수가 어떤 분포를 갖는가를 계산한 사람이 Landau인데, Laplace transform을 포함하는 수학적 계산 과정은 그냥 믿기로 한다면, 함수 f 를 다음과 같이 얻게되고, 확률함수 ϕ 를 제한된 과정에서 앞에서 얻은 에너지 loss를 주고, 이 값을 무한대까지 적분함으로써 얻을 수 있다는 것을 보여줍니다.

단위계를 거리에거 에너지 x_i 로 변환하면, probability function을 변수 lambda의 변수를 갖는 함수로 표현할 수 있고, 이 확률 함수가 최대가 되는 조건은 이와 같이 되고, 이 값을 most probable value라고 함. 이 함수는 이와 같이 주어지고, 이러한 분포를 Landau distribution이라고 함.

With a variable

$$f(x, \Delta) = \frac{1}{\xi} \phi(\lambda)$$

$$\phi(\lambda) = \frac{1}{2\pi i} \int_{-i\infty+\sigma}^{+i\infty+\sigma} e^{u \ln u + \lambda u} du$$

$$\lambda = \frac{\Delta - \xi(\ln \frac{\xi}{\varepsilon'} + 1 - C)}{\xi}$$

Most probable value of energy loss :

$$f(x, \Delta) d\Delta = \phi\left(\frac{\Delta - \Delta_0}{\xi}\right) d\left(\frac{\Delta}{\xi}\right)$$

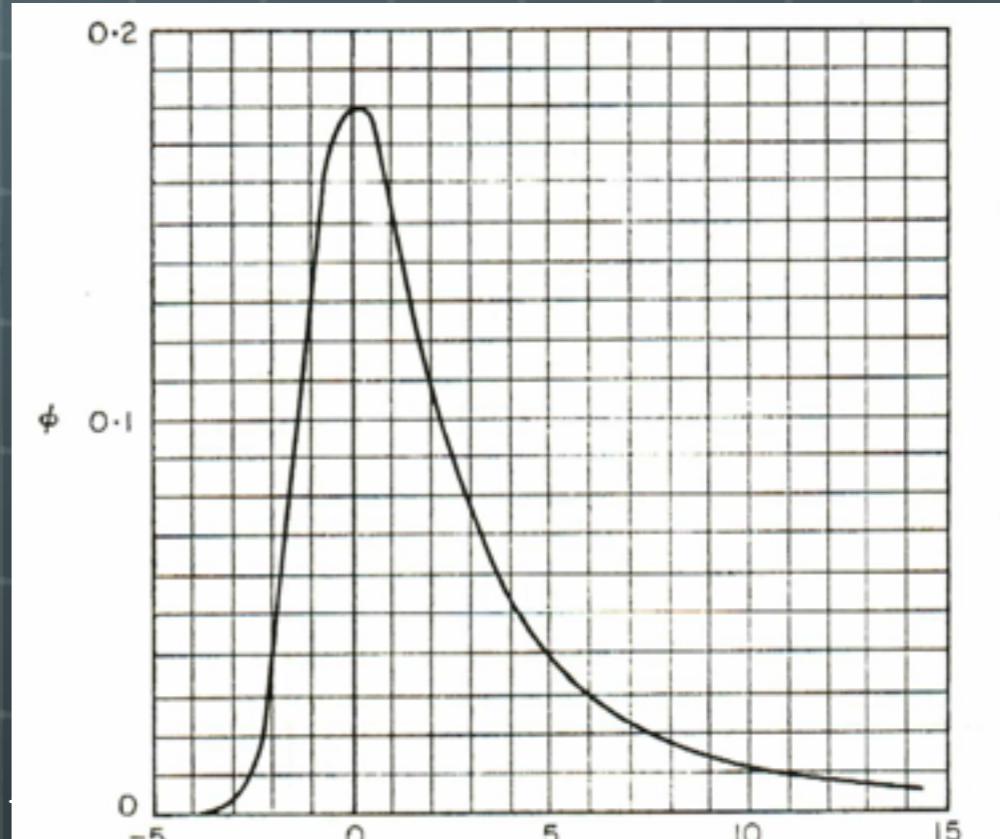


FIG. 1.

단위계를 거리에거 에너지 x_i 로 변환하면, probability function을 변수 lambda의 변수를 갖는 함수로 표현할 수 있고, 이 확률 함수가 최대가 되는 조건은 이와 같이 되고, 이 값을 most probable value라고 함. 이 함수는 이와 같이 주어지고, 이러한 분포를 Landau distribution이라고 함.

With a variable $\xi = x \frac{2\pi N e^\zeta \rho \sum Z}{mv^2 \sum A}$

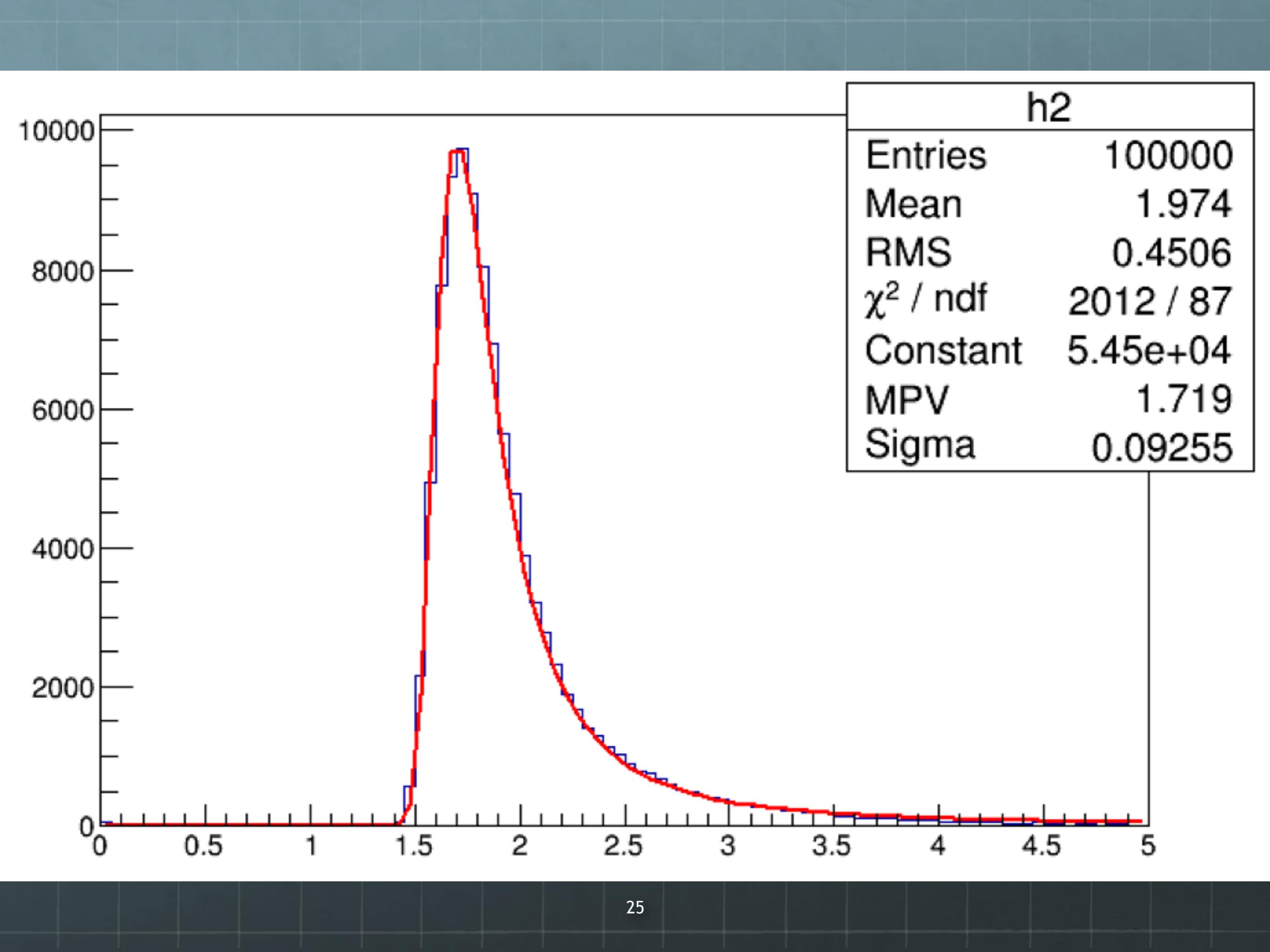
$$f(x, \Delta) = \frac{1}{\xi} \phi(\lambda)$$

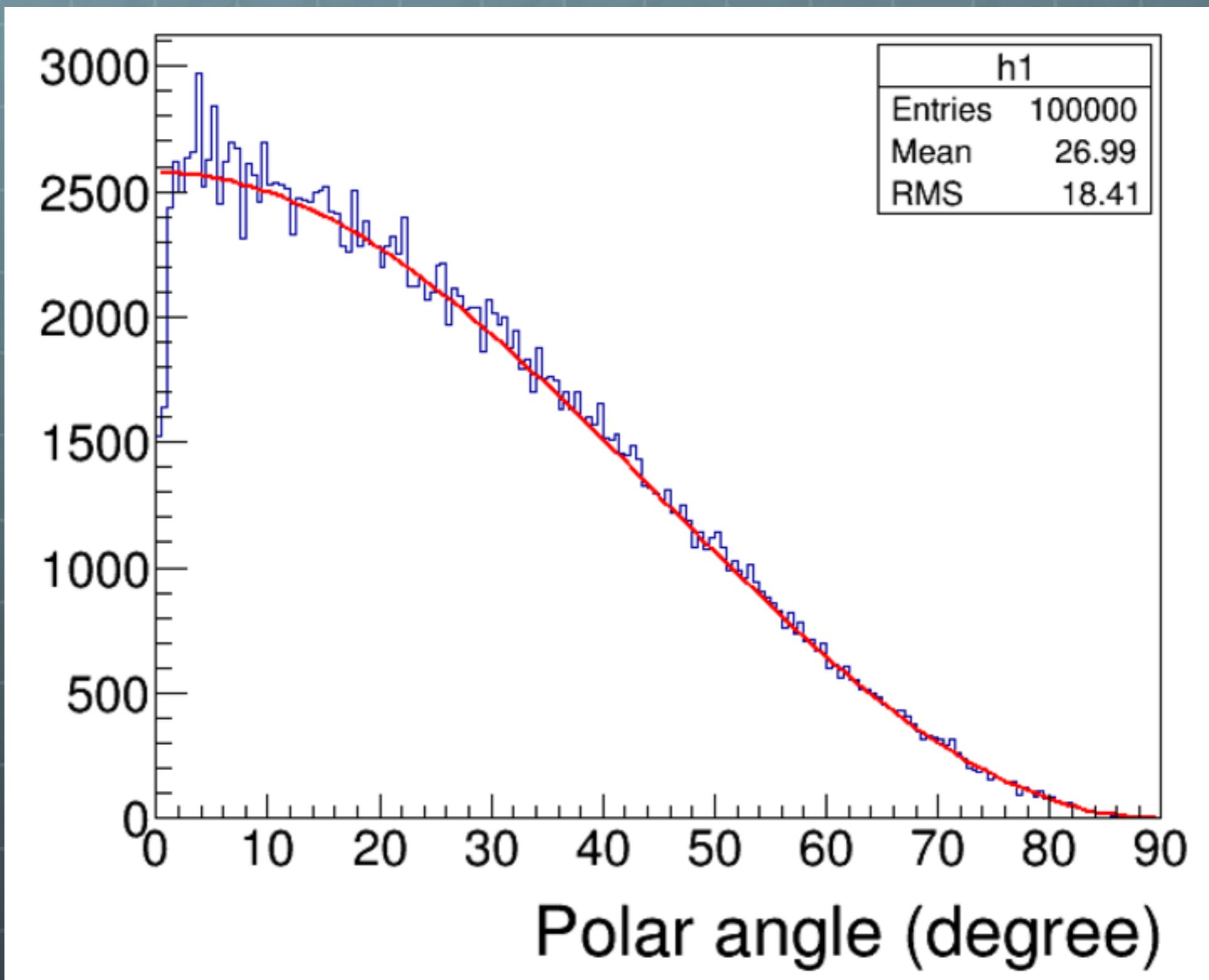
$$\phi(\lambda) = \frac{1}{2\pi i} \int_{-i\infty + \sigma}^{+i\infty + \sigma} e^{u \ln u + \lambda u} du$$
Maximum at $\lambda = -0.05$

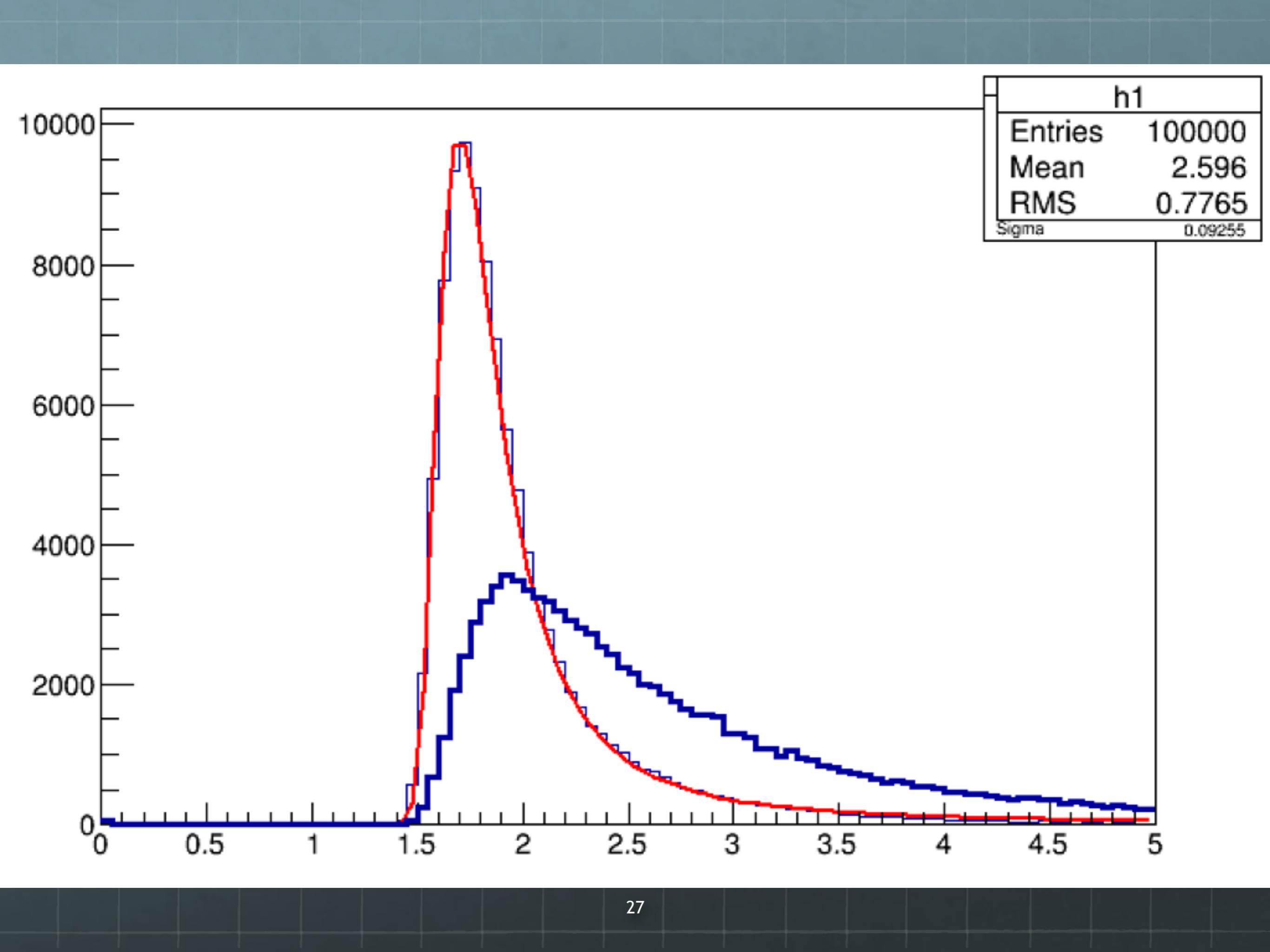
$$\lambda = \frac{\Delta - \xi \left(\ln \frac{\xi}{\varepsilon'} + 1 - C \right)}{\xi}$$

Most probable value of energy loss : $\Delta_0 = \xi \left(\ln \frac{\xi}{\varepsilon'} + 0.37 \right)$

$$f(x, \Delta) d\Delta = \phi\left(\frac{\Delta - \Delta_0}{\xi}\right) d\left(\frac{\Delta - \Delta_0}{\xi}\right)$$





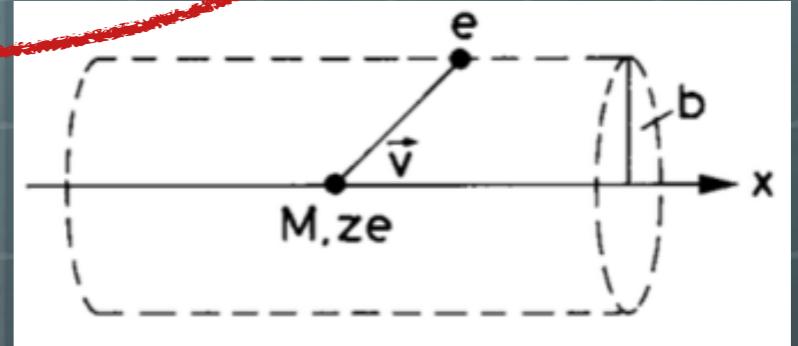


Energy LOSS of e^\pm

Energy loss of e^\pm

$$\left\langle -\frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[\frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$

$$W_{\max} = \frac{2m_e c^2 \beta^2 \gamma^2}{1 + 2\gamma m_e/M + (m_e/M)^2}$$



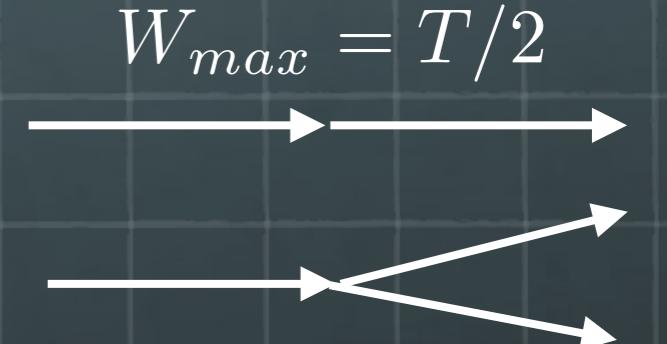
$$\frac{1}{2} \ln \left(\frac{\tau^2 (\tau + 2)}{2(I/m_e c^2)^2} \right) + F(\tau)$$

for e^-

$$F(\tau) = 1 - \beta^2 + \frac{(\tau^2/8) - (2r+1)\ln 2}{(\tau+1)^2}$$

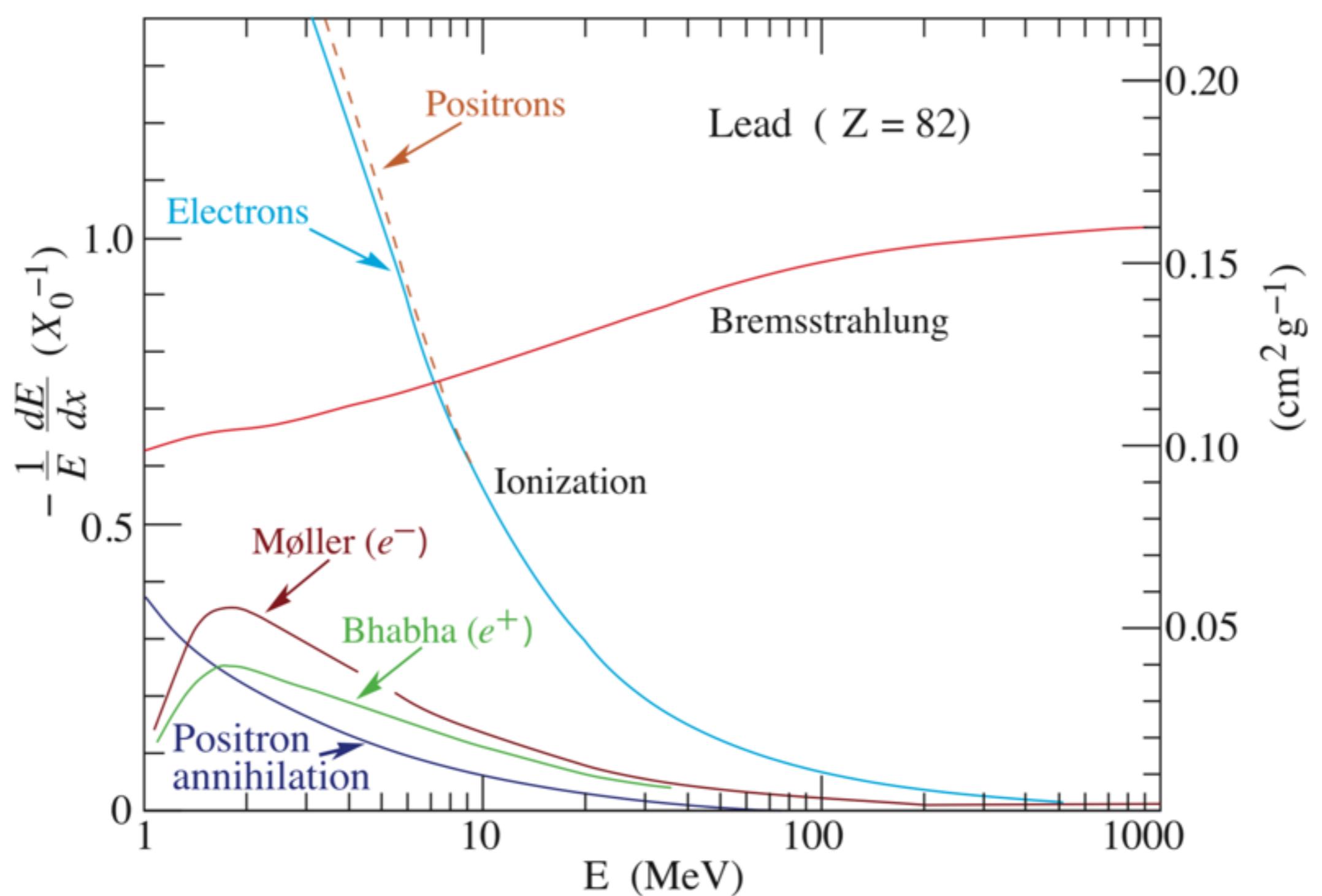
for e^+

$$F(\tau) = 2\ln 2 - \frac{\beta^2}{12} \left(23 + \frac{14}{\tau+2} + \frac{10}{(\tau+2)^2} + \frac{4}{(\tau+2)^3} \right)$$



$$W_{\max} = T$$

Energy loss of e^\pm



Bremsstrahlung

$$\frac{dE}{dx} = 4\alpha N_A z^2 \frac{Z^2}{A} \left(\frac{e^2}{mc^2} \right)^2 E \ln \frac{183}{Z^{1/3}}$$

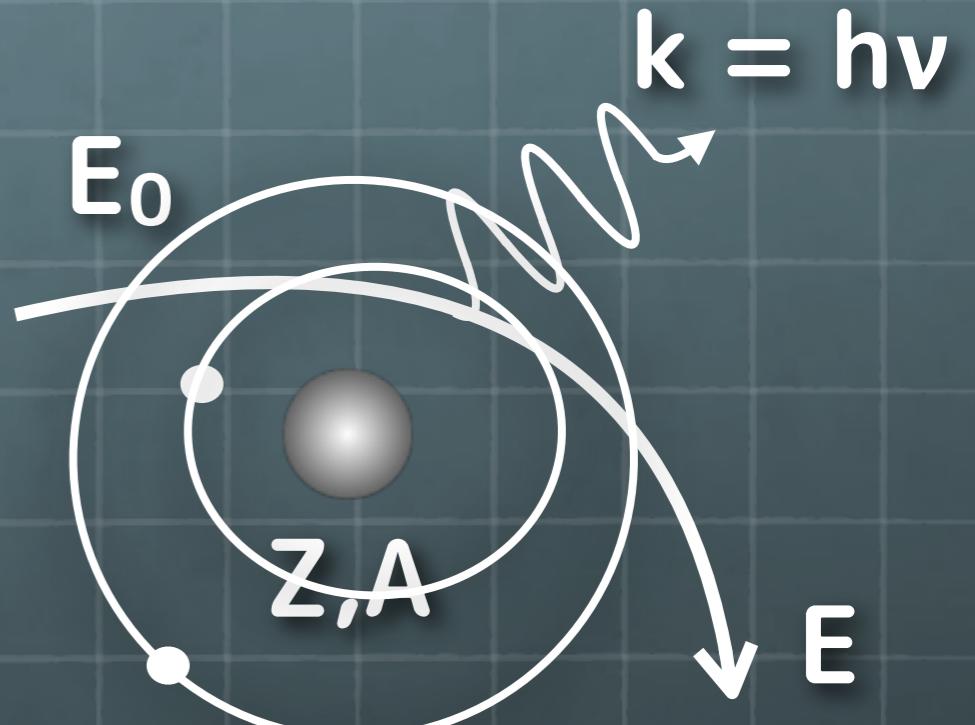
$$= K z^2 \frac{Z}{A} \left(\frac{Z}{m_e c^2} \right) \frac{\alpha}{\pi} E \cdot \ln \frac{183}{Z^{1/3}}$$

$$= X_0 E$$

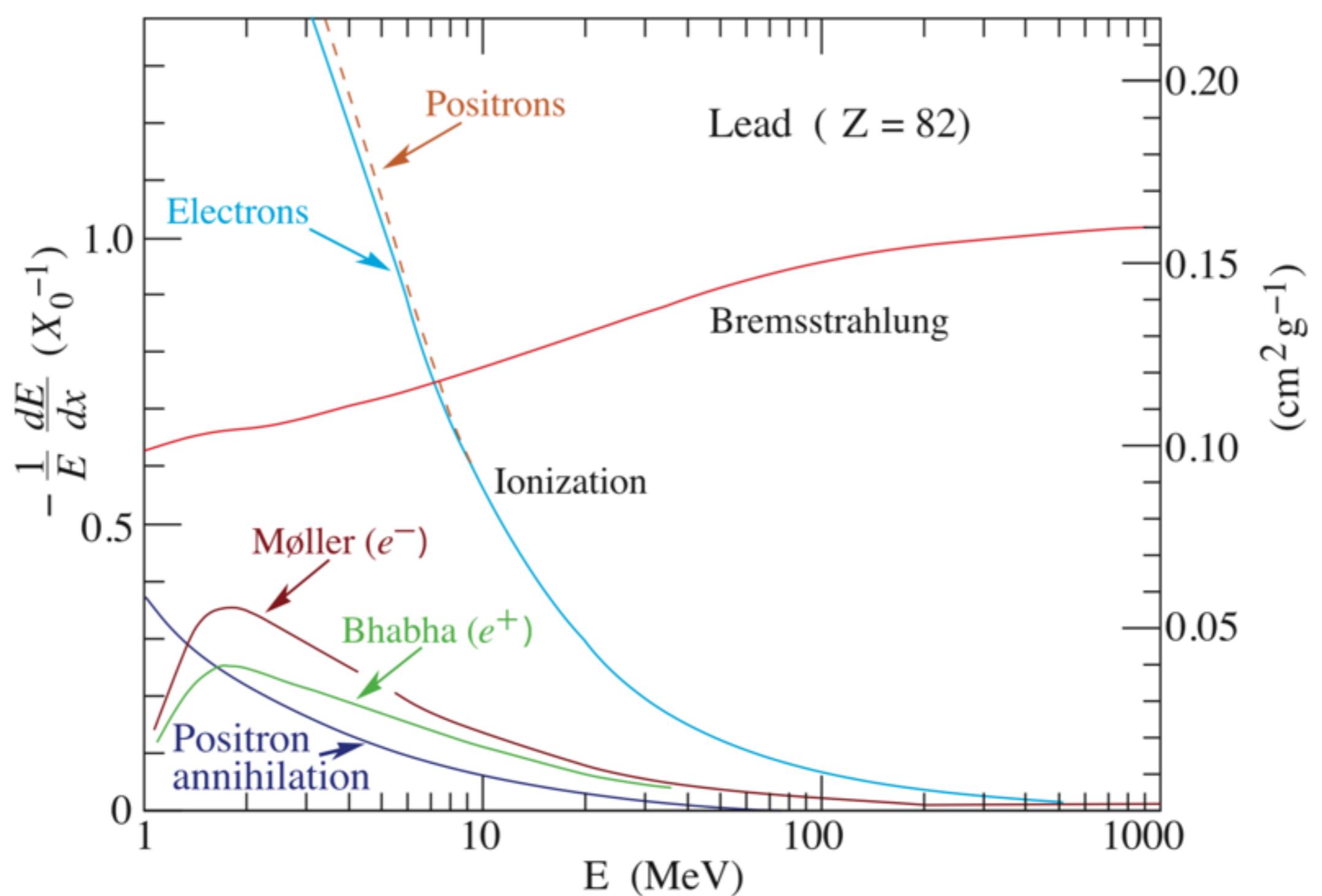
$$E(x) = E_0 e^{-x/X_0}$$

Radiation Length X_0

$$X_0 = \frac{A}{4\alpha N_A Z^2 r_e^2 \ln \frac{183}{Z^{1/3}}}$$



Energy loss of e^\pm



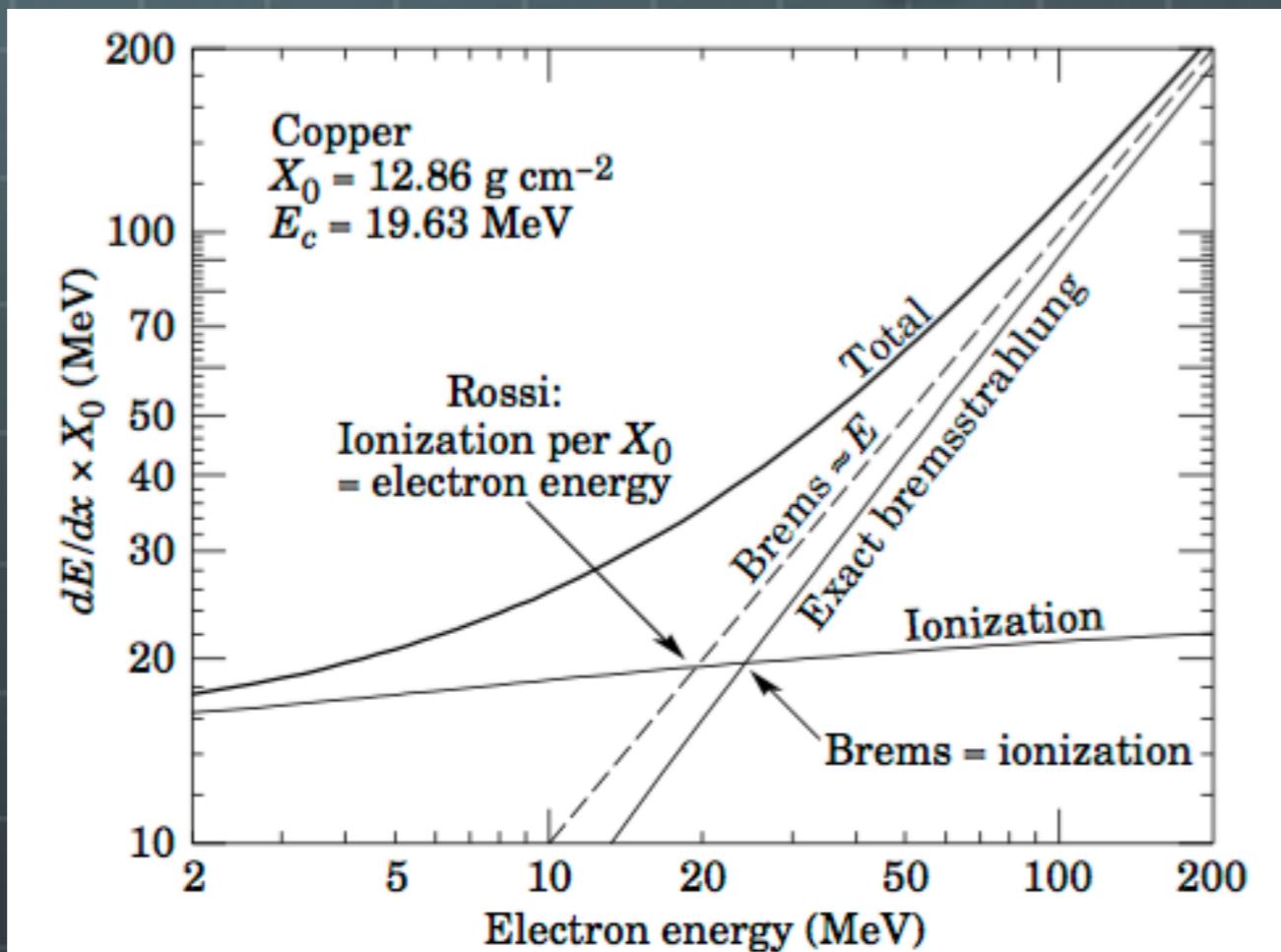
Critical energy (E_c)



Energy at which a electron losses its energy as same amount by bremsstrahlung and ionization.



Energy at which the ionization loss per X_0 is equal to the electron energy.



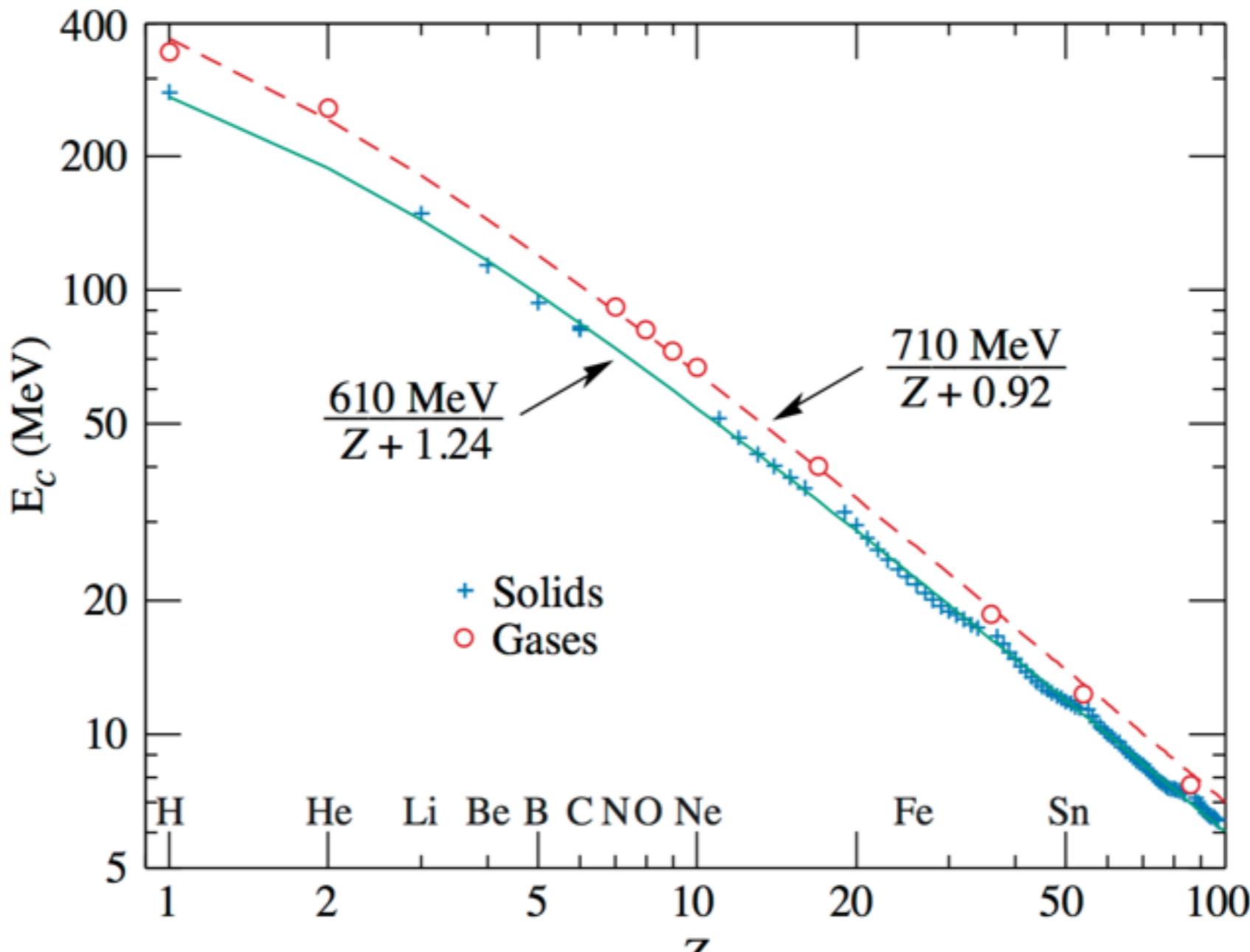


Figure 33.14: Electron critical energy for Z the chemical elements, using Rossi's definition [2]. The fits shown are for solids and liquids (solid line) and gases (dashed line). The rms deviation is 2.2% for the solids and 4.0% for the gases. (Computed with code supplied by A. Fassó.)

Cherenkov radiation



In 1934, P.A. Cherenkov observe new type of luminescence irradiating gamma rays into uranyl salt.



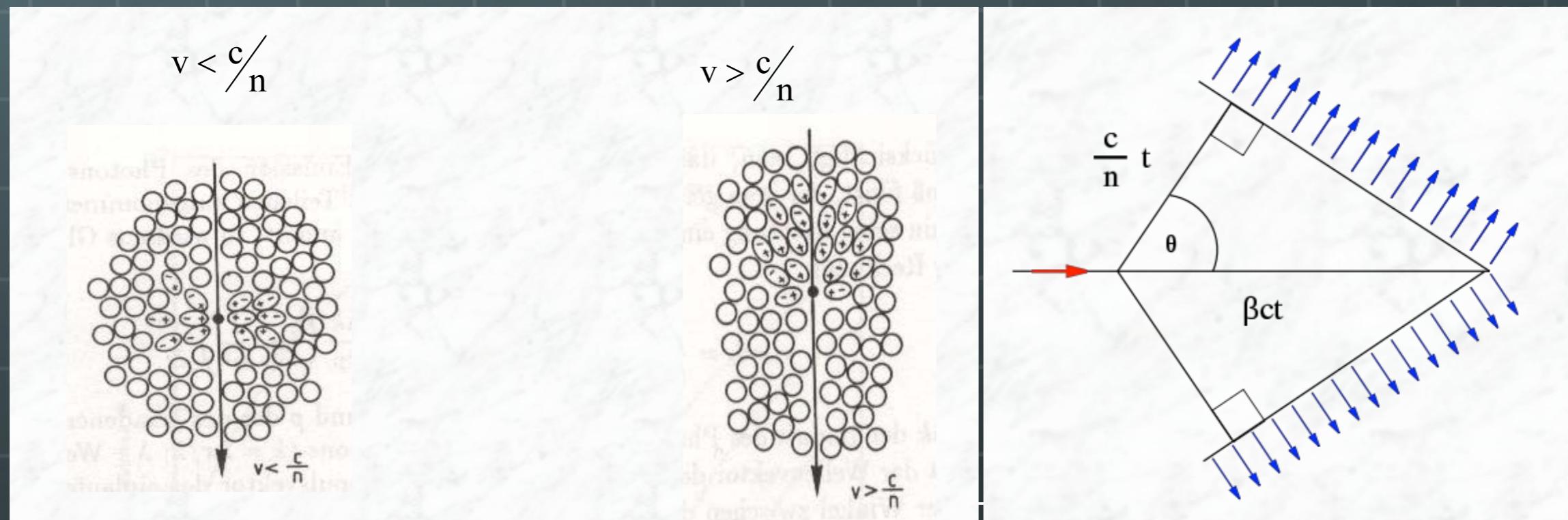
Originated by charged particle



Not to be radiative origin



Observed at a certain angle along particle direction



Refractive indices, Cherenkov threshold values

Material	$n - 1$	β -Schwelle	γ -Schwelle
festes Natrium	3.22	0.24	1.029
Bleisulfit	2.91	0.26	1.034
Diamant	1.42	0.41	1.10
Zinksulfid ($ZnS(Ag)$)	1.37	0.42	1.10
Silberchlorid	1.07	0.48	1.14
Flintglas (SFS1)	0.92	0.52	1.17
Bleifluorid	0.80	0.55	1.20
Clerici-Lösung	0.69	0.59	1.24
Bleiglas	0.67	0.60	1.25
Thalliumformiat-Lösung	0.59	0.63	1.29
Szintillator	0.58	0.63	1.29
Plexiglas	0.48	0.66	1.33
Borsilikatglas	0.47	0.68	1.36
Wasser	0.33	0.75	1.52
Aerogel	0.025 - 0.075	0.93 - 0.976	4.5 - 2.7
Pentan (STP)	$1.7 \cdot 10^{-3}$	0.9983	17.2
CO_2 (STP)	$4.3 \cdot 10^{-4}$	0.9996	34.1
Luft (STP)	$2.93 \cdot 10^{-4}$	0.9997	41.2
H_2 (STP)	$1.4 \cdot 10^{-4}$	0.99986	59.8
He (STP)	$3.3 \cdot 10^{-5}$	0.99997	123

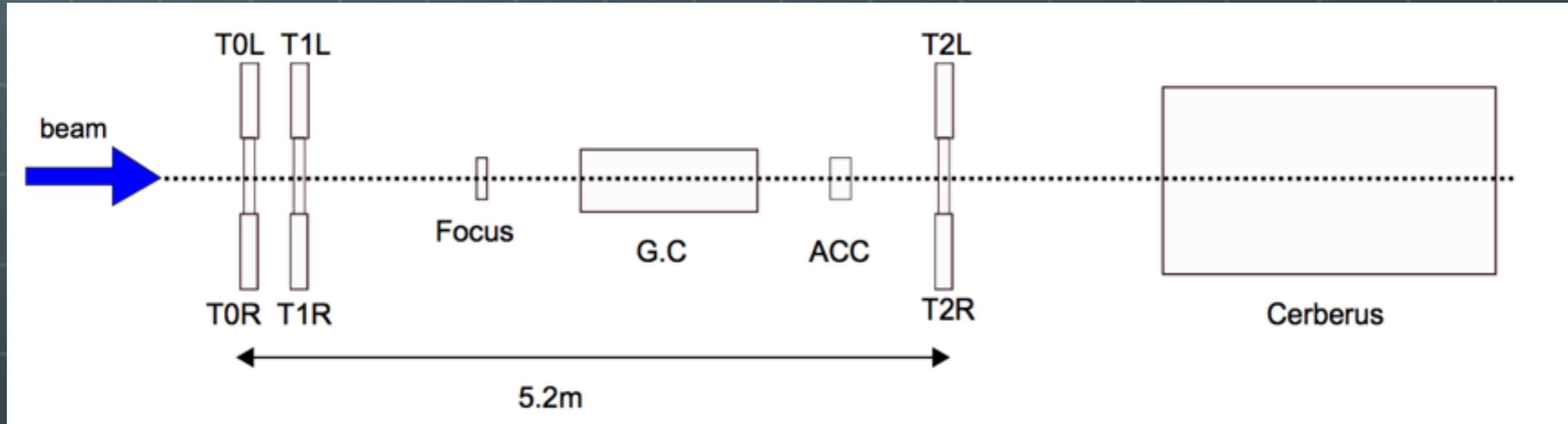
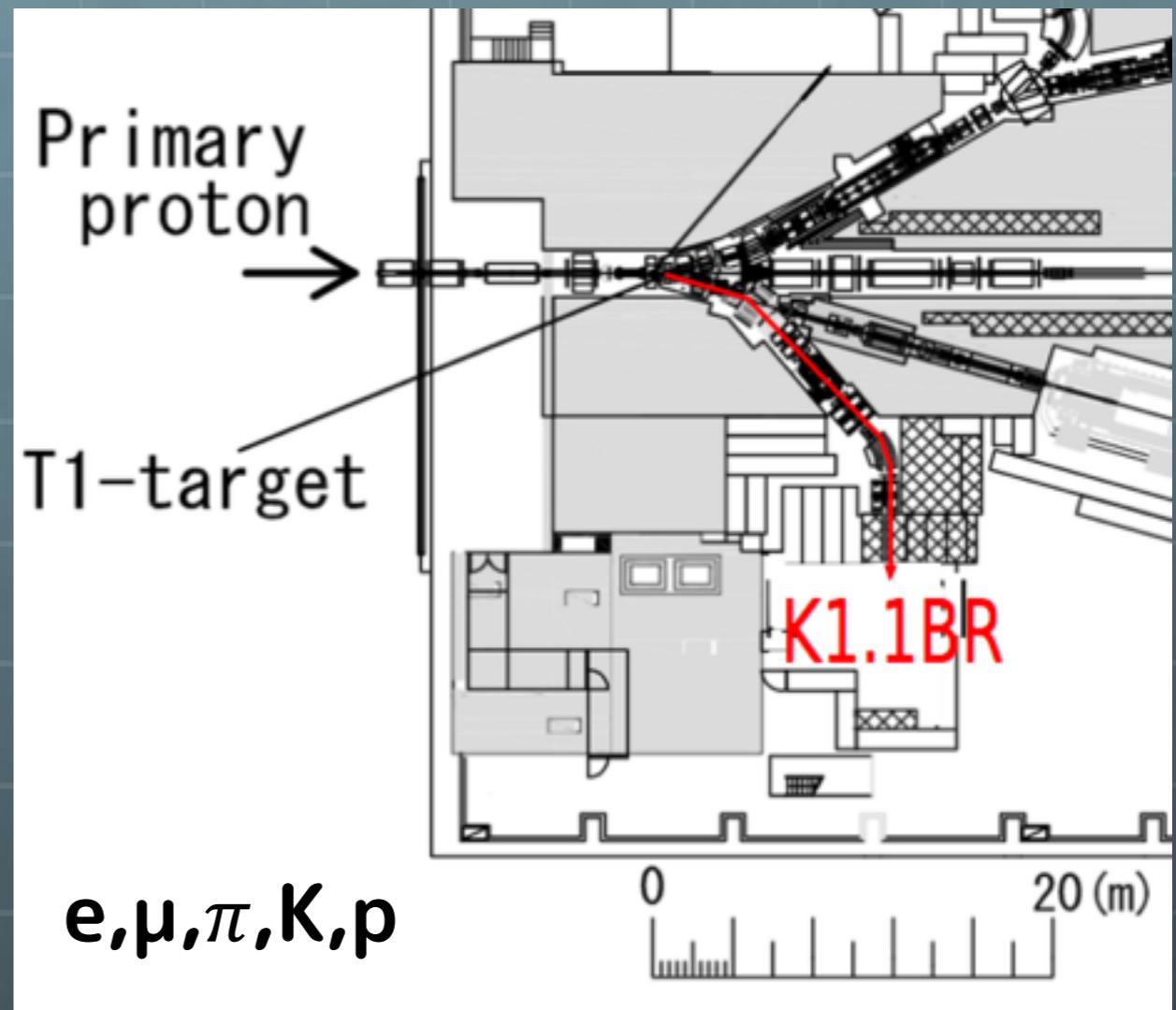
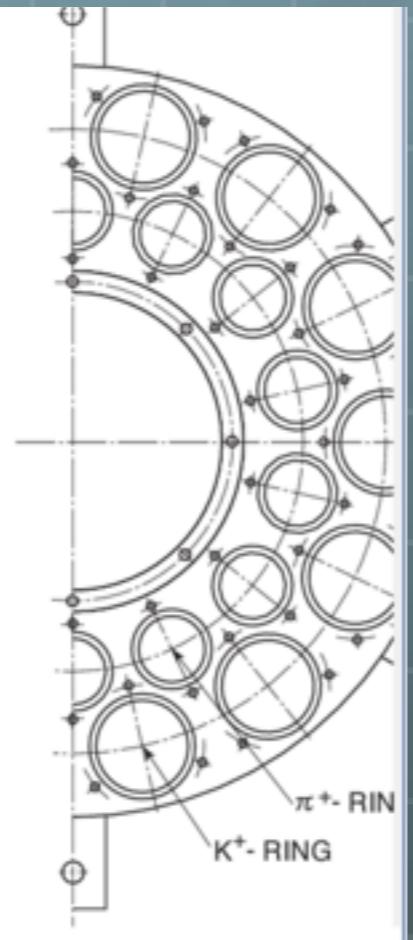
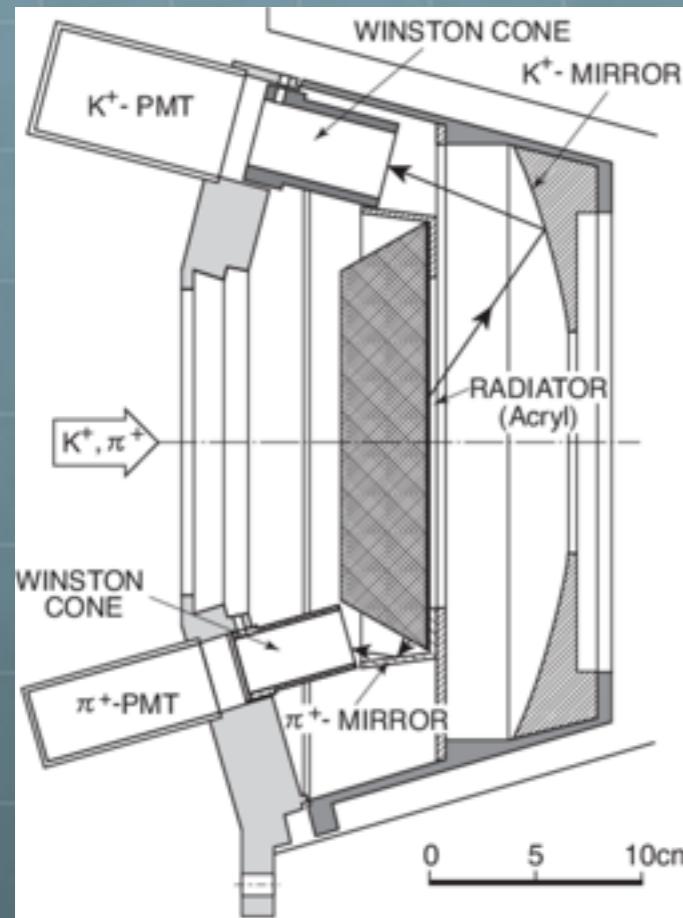
problematic: region between liquids and gases

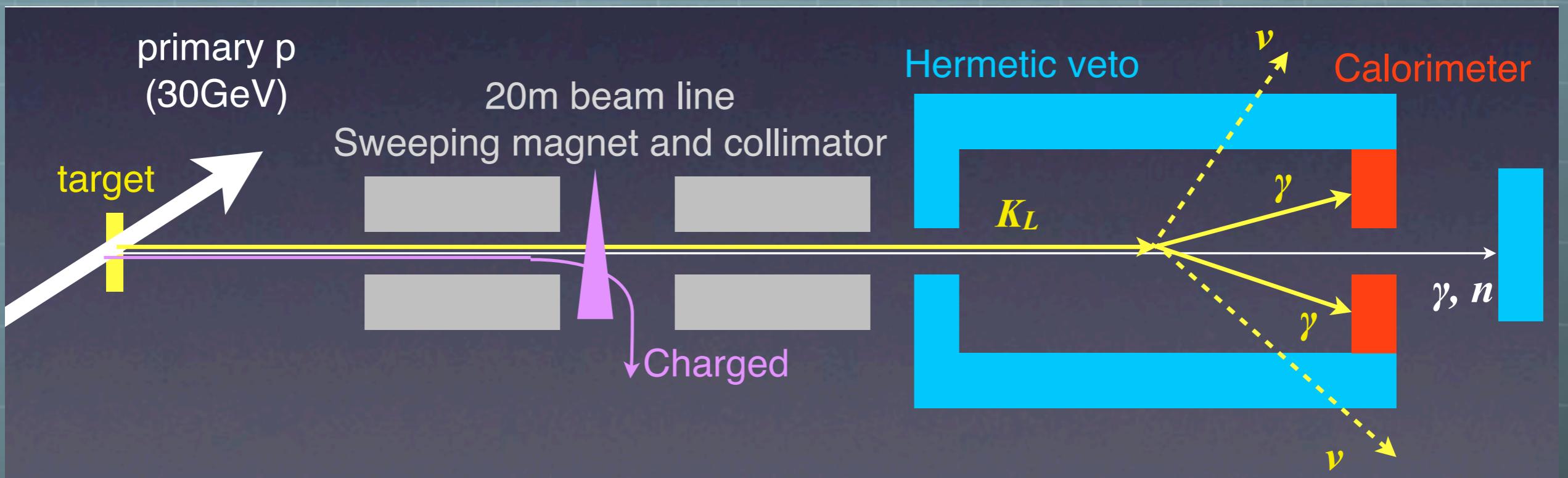
Aerogel: mixture of $m(SiO_2) + 2m(H_2O)$

light structure with inclusions of air, bubbles with diameter $< \lambda_{Licht}$

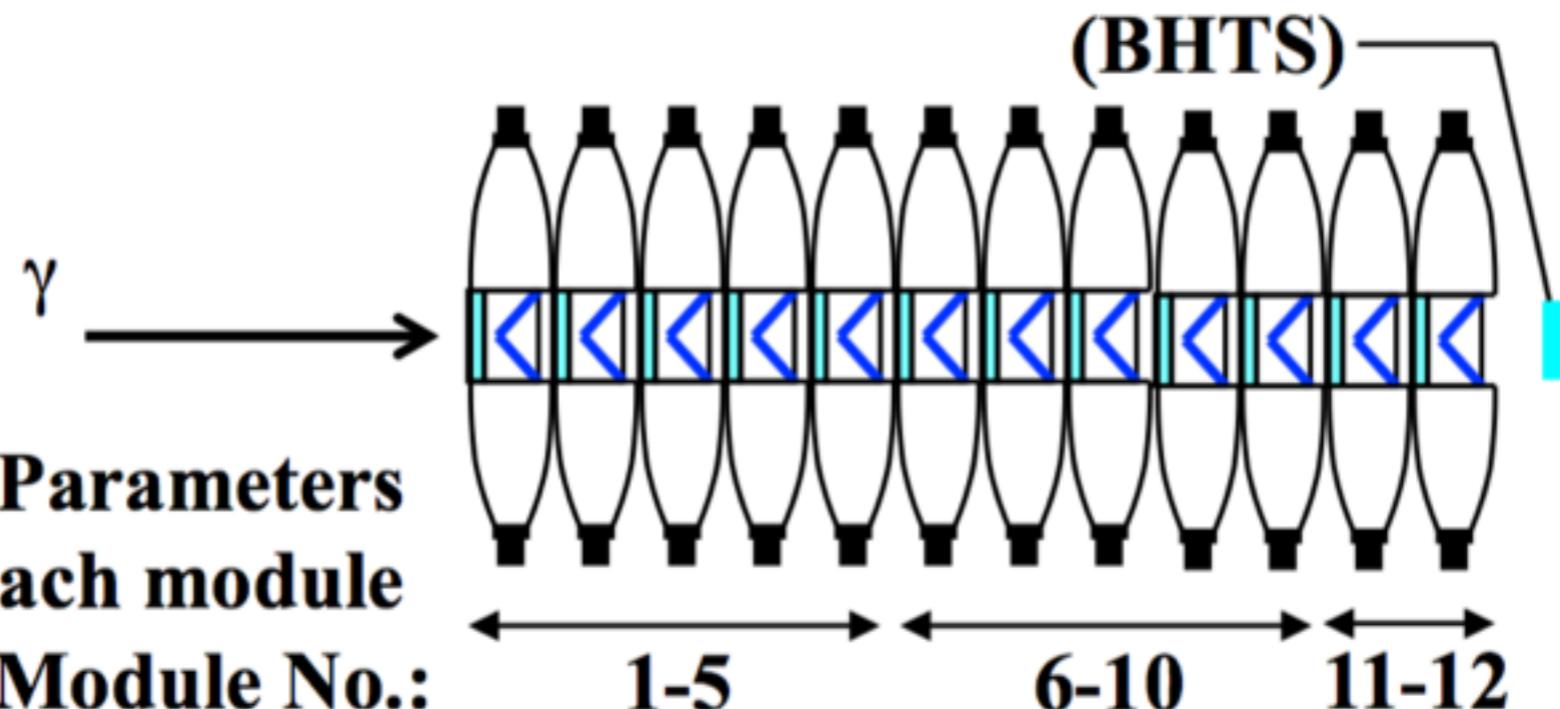
→ n : average from n_{air} , n_{SiO_2} , n_{H_2O}

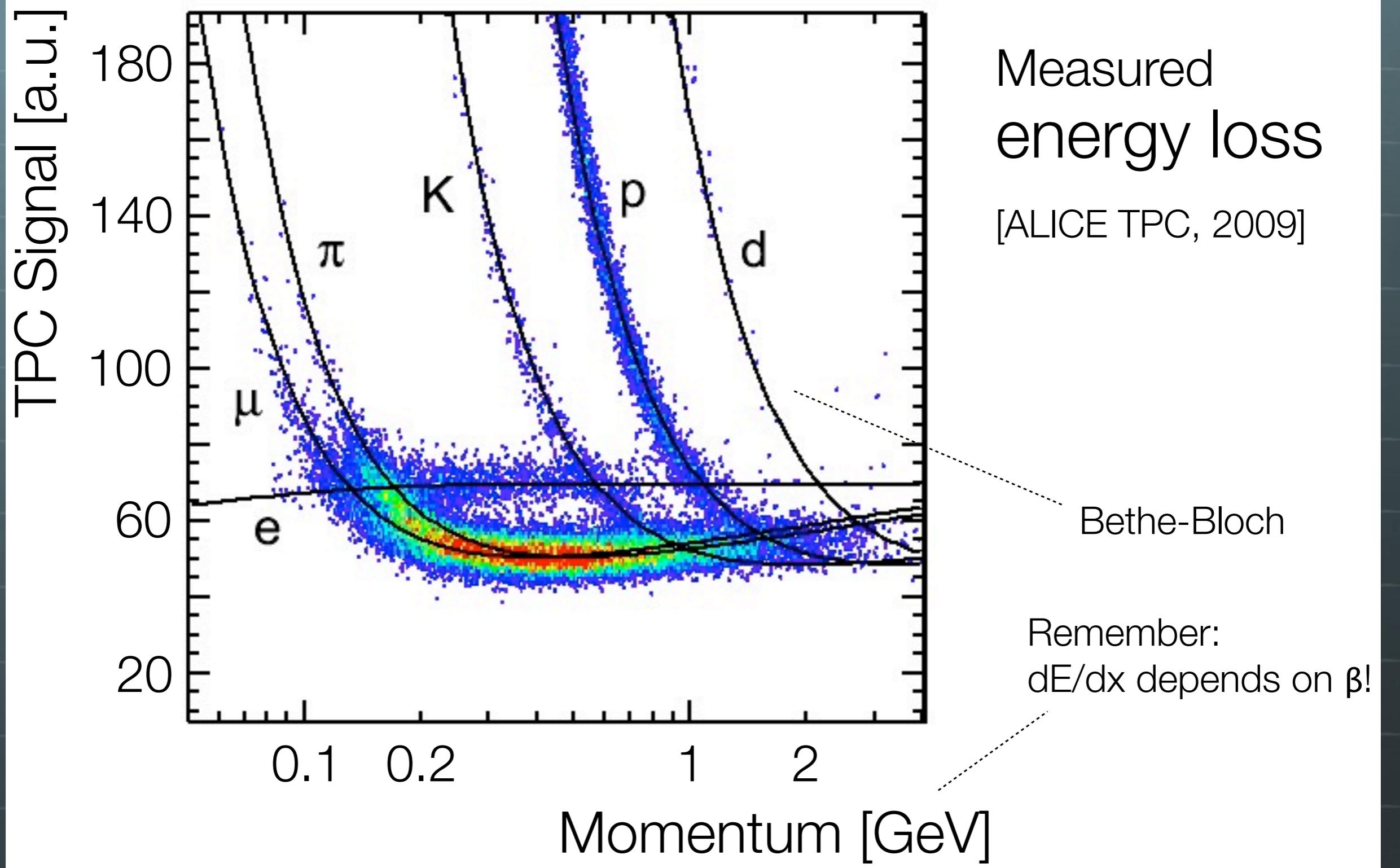
Tabelle 6.2: Cherenkov-Radiatoren [94, 32, 313]. Der Brechungsindex für Gase bezieht sich auf $0^\circ C$ und 1 atm (STP). Festes Natrium ist für Wellenlängen unterhalb von 2000 \AA transparent [373, 209].





Trigger scintillator for calibration





Transition radiation (TRD)

Charged particle passes through materials with different dielectric properties

- particle forms dipole with the mirror charge
- dipole changes with time
- radiation

- radiated energy W proportional to the energy of particle!

$$W = \frac{1}{3} \alpha \hbar \omega_p \gamma$$

- with ω_p Plasma frequency

$$\omega_p = \sqrt{\frac{N_e e^2}{\epsilon_0 m_e}} \quad \hbar \omega_p = 20 \text{ eV}$$

- only important for highly relativistic particles

- energy: keV (x-rays)

- $\theta \propto 1/\gamma$: emission in very forward direction

- probability for photon emission very small → many transitions (foils with gaps)

- # photons $\langle N \rangle \sim W / h\nu \sim O(\alpha) = 1/137$ α fine structure constant

- energy loss due to TRD negligible for single transition

- important for particle ID at high energies, other effects used for PID $\propto \beta$ ($\beta \approx 1$)

Review article: B. Dolgoshein; NIM A 326 (1993) 434

Air (Vacuum)



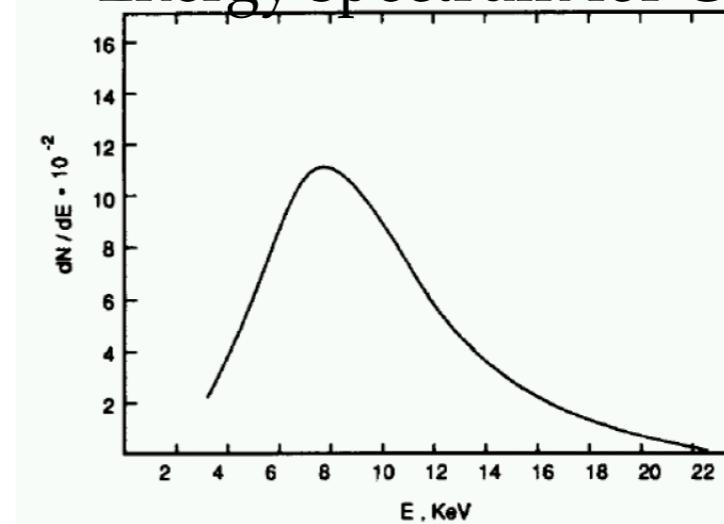
charged part.

Dielectric medium



mirror charge

Energy spectrum for CH₂ foil

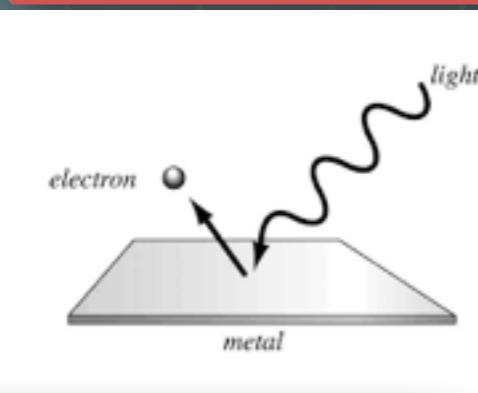


keV!

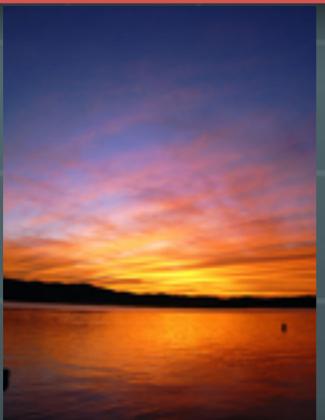
Interaction of photons

Interaction of photons

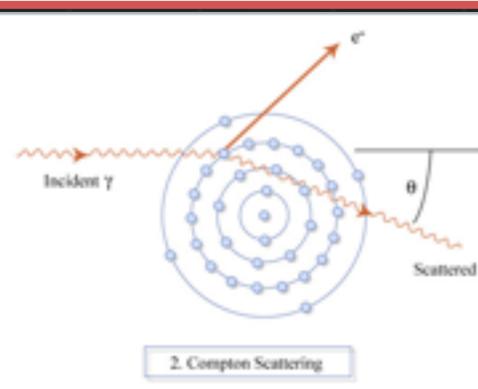
Photoelectric effect



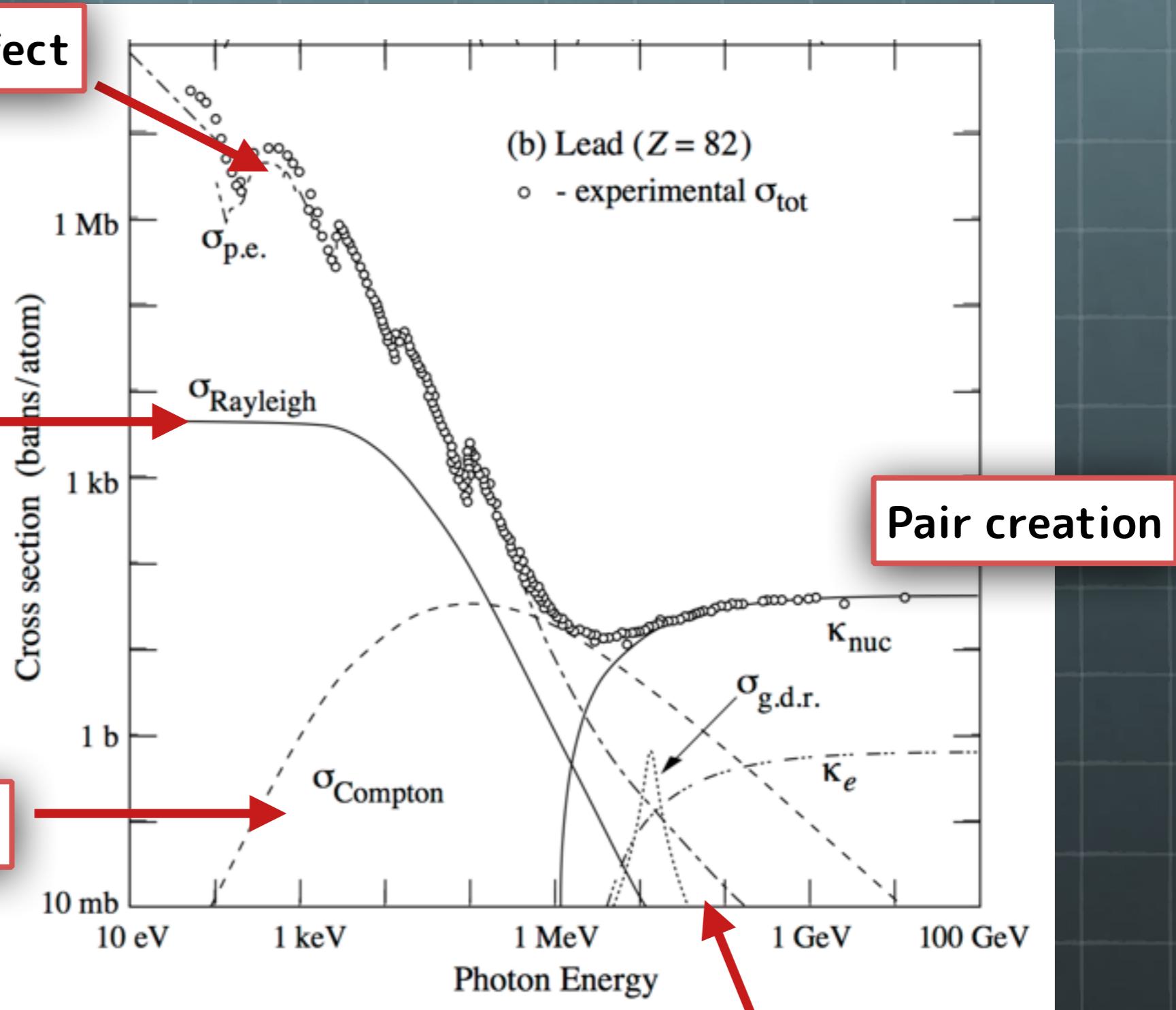
Rayleigh scattering



Compton scattering



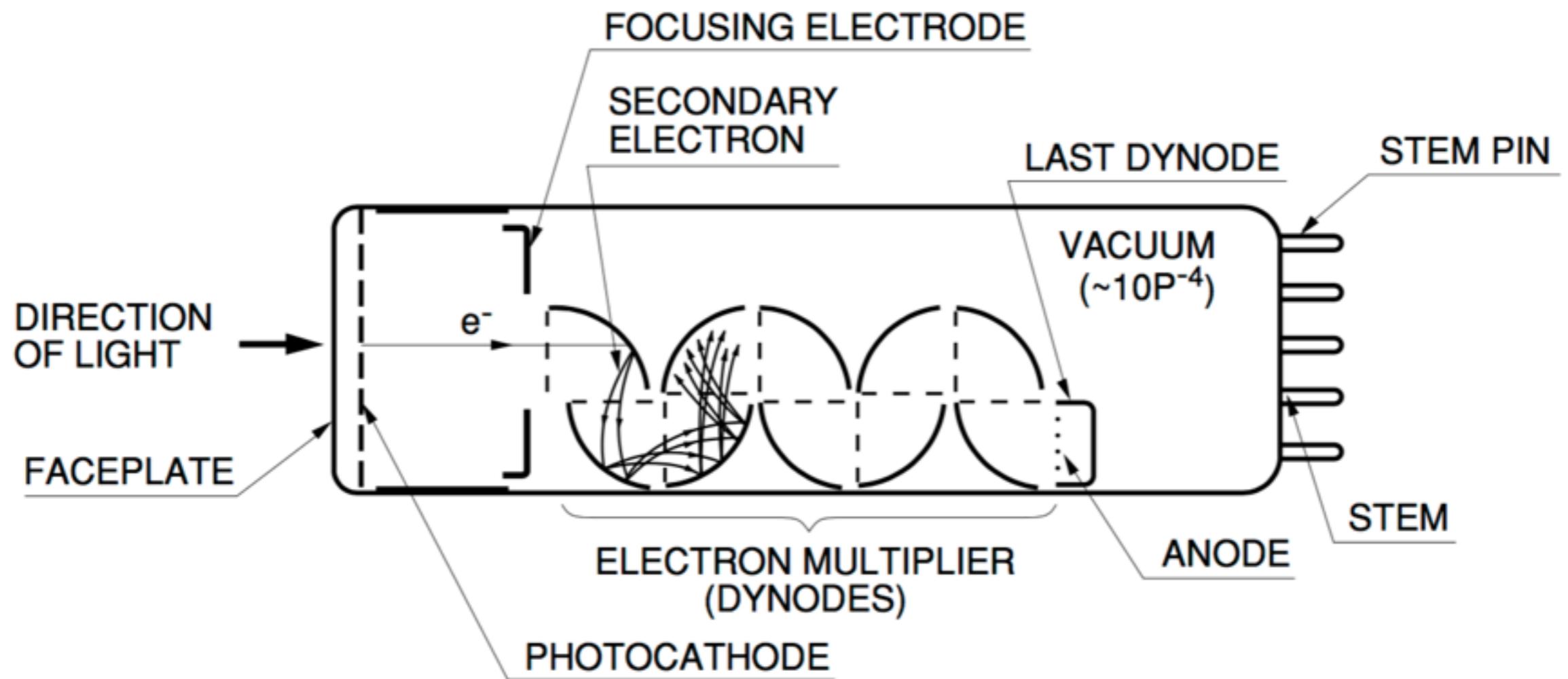
(b) Lead ($Z = 82$)
○ - experimental σ_{tot}



Pair creation

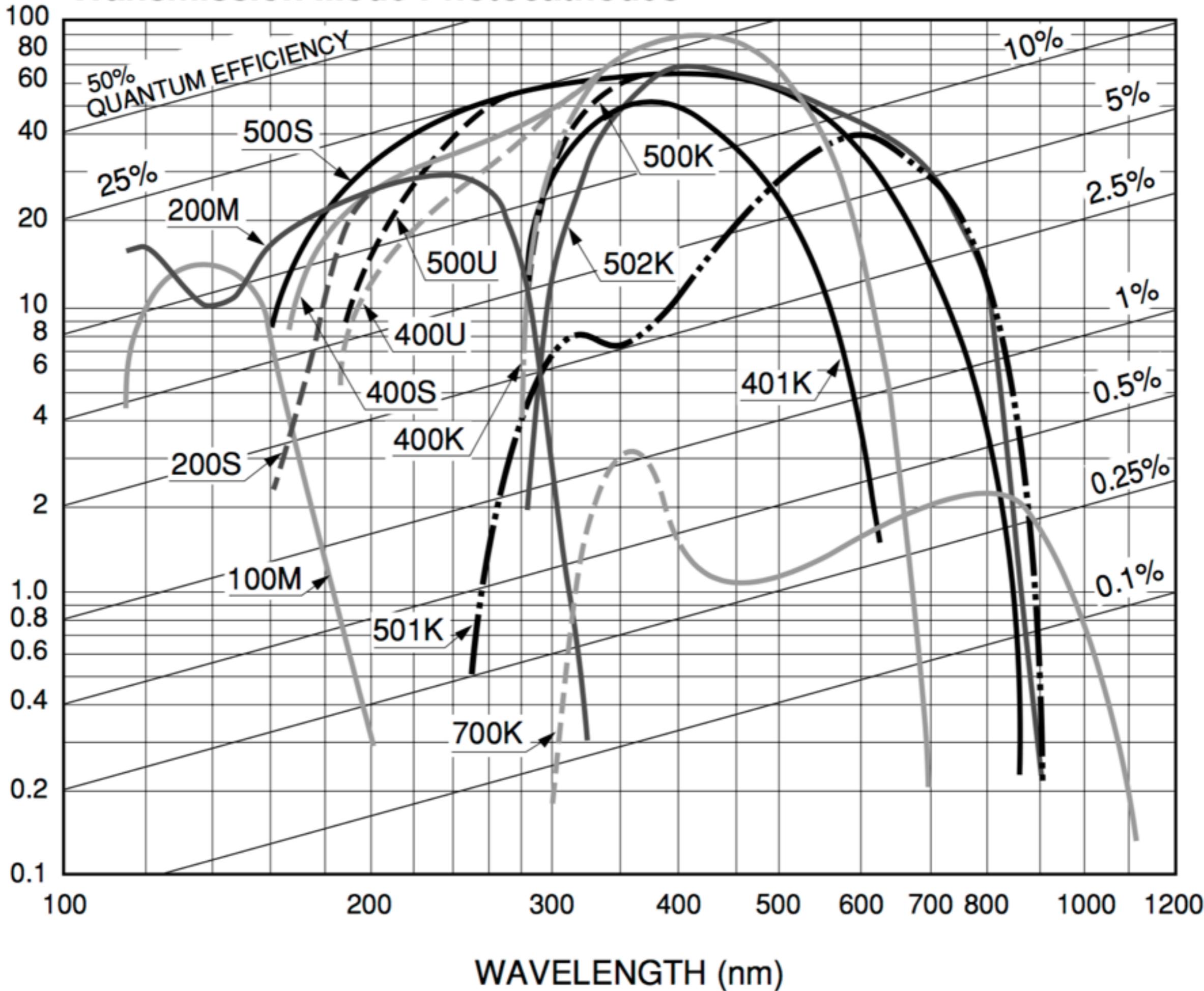
Photonuclear interaction

PMT



Transmission Mode Photocathodes

PHOTOCATHODE RADIANT SENSITIVITY (mA/W)



Transmission mode photocathodes

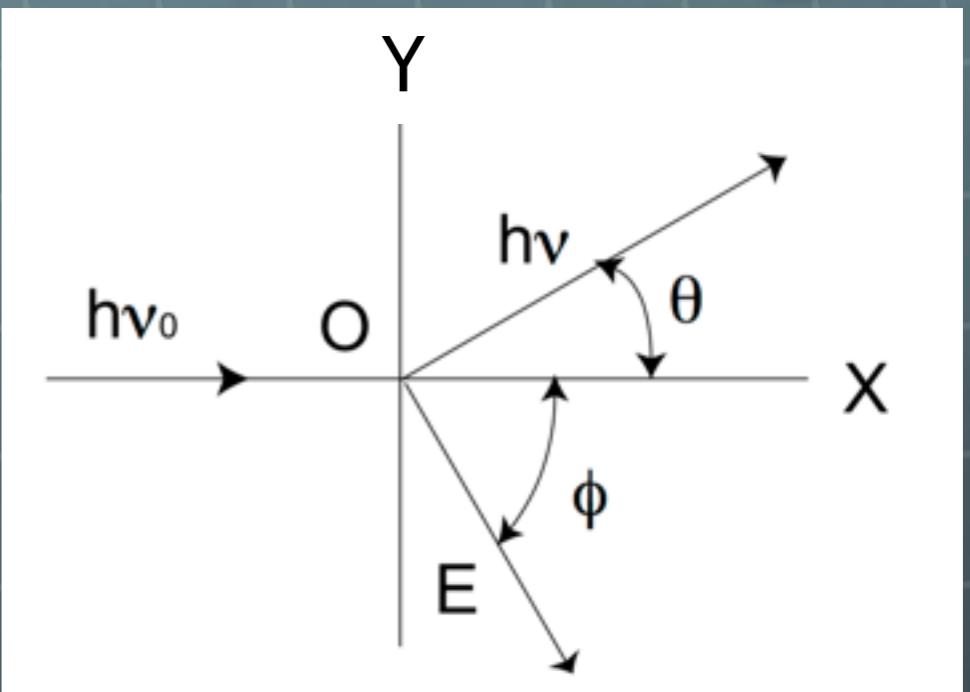
Curve Code (S number)	Photocathode Material	Window Material	Luminous Sensitivity (Typ.) (μ A/lm)	Spectral Response				
				Spectral Range (nm)	Peak Wavelength			
					Radiant Sensitivity (mA/W)		(nm)	(%)
100M	Cs-I	MgF ₂	—	115 to 200	14	140	13	130
200S	Cs-Te	Quartz	—	160 to 320	29	240	14	210
200M	Cs-Te	MgF ₂	—	115 to 320	29	240	14	200
400K	Bialkali	Borosilicate	95	300 to 650	88	420	27	390
400U	Bialkali	UV	95	185 to 650	88	420	27	390
400S	Bialkali	Quartz	95	160 to 650	88	420	27	390
401K	High temp. bialkali	Borosilicate	40	300 to 650	51	375	17	375
500K (S-20)	Multialkali	Borosilicate	150	300 to 850	64	420	20	375
500U	Multialkali	UV	150	185 to 850	64	420	25	280
500S	Multialkali	Quartz	150	160 to 850	64	420	25	280
501K (S-25)	Multialkali	Borosilicate	200	300 to 900	40	600	8	580
502K	Multialkali	Borosilicate (prism)	230	300 to 900	69	420	20	390
700K (S-1)	Ag-O-Cs	Borosilicate	20	400 to 1200	2.2	800	0.36	740
—	GaAsP(Cs)	—	—	300 to 720	180	580	40	540
—	GaAs(Cs)	—	—	380 to 890	85	800	14	760
—	InP/InGaAsP(Cs)	—	—	950 to 1400	21	1300	2.0	1000 to 1300
—	InP/InGaAs(Cs)	—	—	950 to 1700	24	1500	2.0	1000 to 1550

Compton Scattering

$$h\nu = \frac{h\nu_0}{1 + \left(\frac{h\nu_0}{m_e c^2}\right)(1 - \cos \theta)},$$

$$E = h\nu_0 - h\nu = m_e c^2 \frac{2(h\nu_0)^2 \cos^2 \phi}{(h\nu_0 + m_e c^2)^2 - (h\nu_0)^2 \cos^2 \phi},$$

$$\tan \phi = \frac{1}{1 + \left(\frac{h\nu_0}{m_e c^2}\right)} \cot \frac{\theta}{2},$$



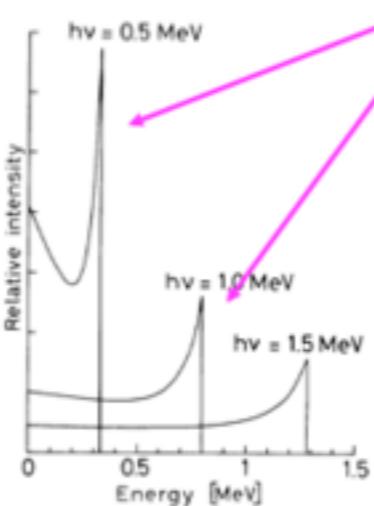
We can also calculate the recoil kinetic energy (T) spectrum of the electron:

$$\frac{d\sigma}{dT} = \frac{\pi r_e^2}{m_e c^2 \gamma^2} \left(2 + \frac{s^2}{\gamma^2 (1-s)^2} + \frac{s}{(1-s)} \left(s - \frac{2}{\gamma} \right) \right) \quad \text{with } s = T / E_{\gamma,in}$$

This cross section is strongly peaked around T_{\max} :

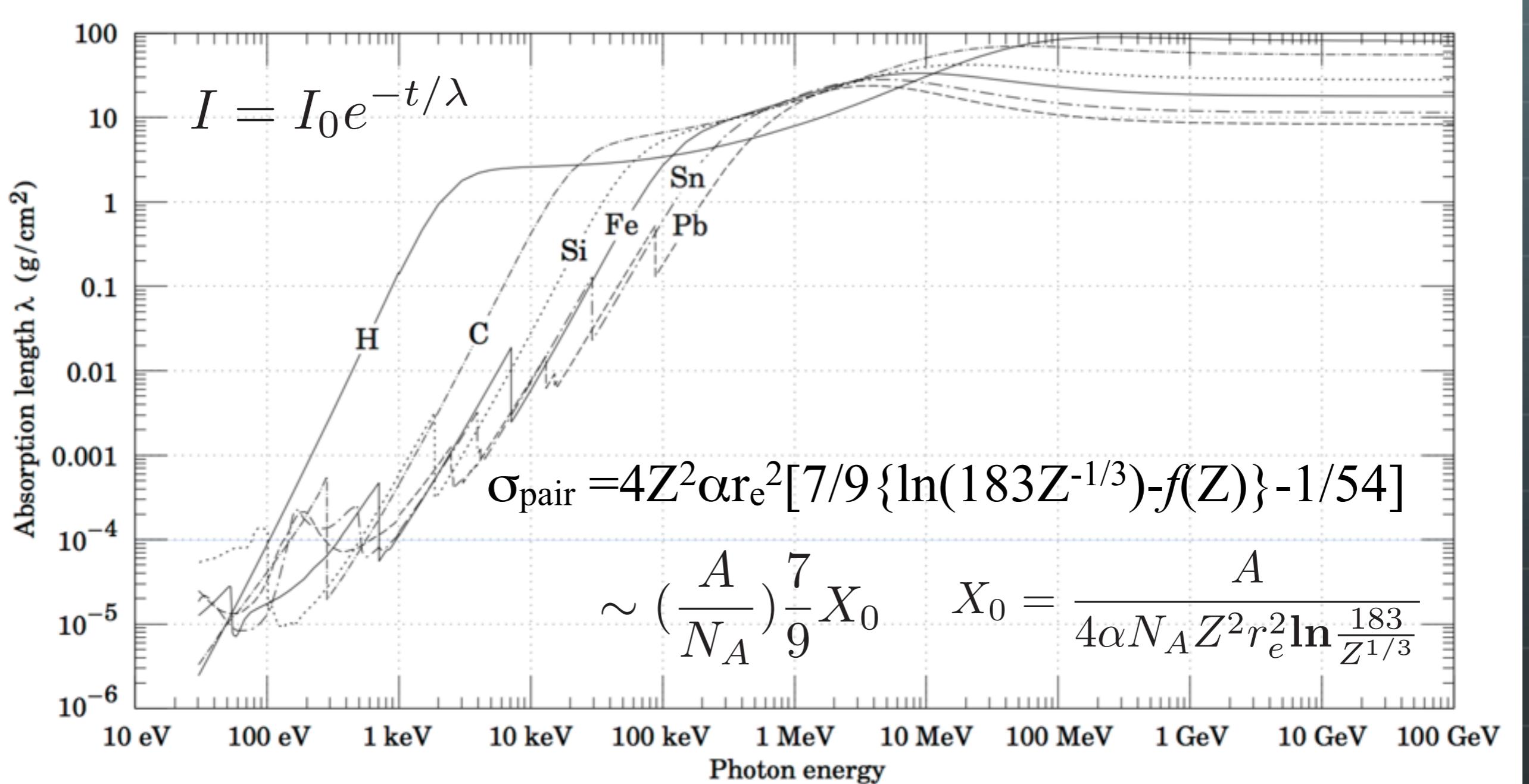
$$T_{\max} = E_{\gamma,in} \frac{2\gamma}{1+2\gamma}$$

T_{\max} is known as the Compton Edge



Kinetic energy distribution of Compton recoil electrons

Pair Production



RESOLUTION OF THE Σ^- -MASS ANOMALY

Walter H. Barkas, John N. Dyer,* and Harry H. Heckman
Lawrence Radiation Laboratory, University of California, Berkeley, California
(Received 29 May 1963)



Mass of hyperons are measured by range in the emulsion

Table I. Observed ranges.

Process	Measured	Mean range	Number of emulsion stacks used
(1) $\Sigma^+ \rightarrow p + \pi^0$	144 p	$1677 \pm 2 \mu$	5
(2) $\Sigma^+ \rightarrow n + \pi^+$	48 π^+	$92.74 \pm 0.34 \text{ mm}$	2
(3) $K + p \rightarrow \Sigma^+ + \pi^-$	40 Σ^+ 24 π^-	$818.8 \pm 1.7 \mu$ $88.58 \pm 0.51 \text{ mm}$	6
(4) $K + p \rightarrow \Sigma^- + \pi^+$	94 Σ^- 63 π^+	$708.9 \pm 1.5 \mu$ $78.45 \pm 0.25 \text{ mm}$	5

expected as 684 (5)

It is well known that stopping theory based on the first Born approximation fails when the particle velocity becomes comparable to the velocities of many of the electrons in the stopping material. It is perhaps not surprising that there