

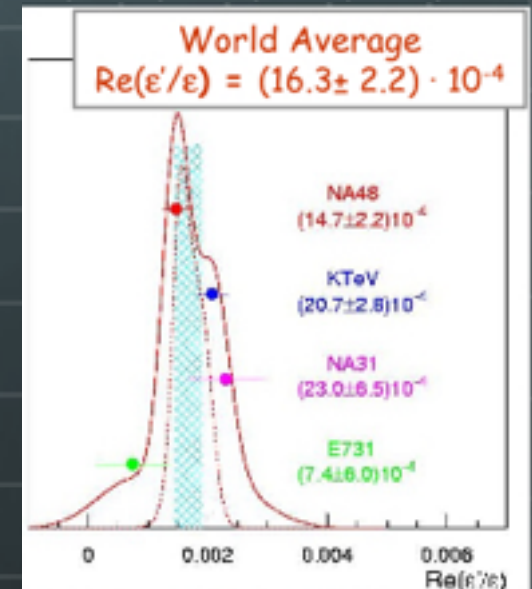
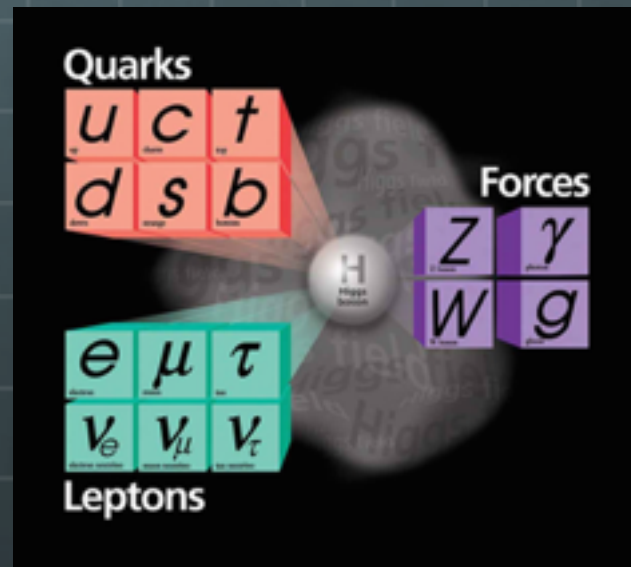
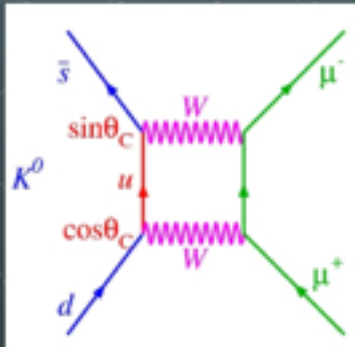
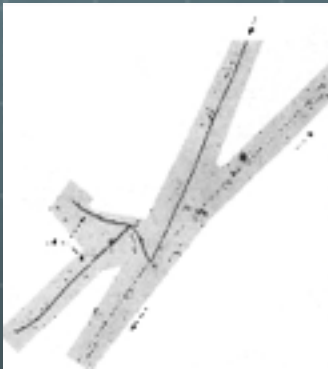
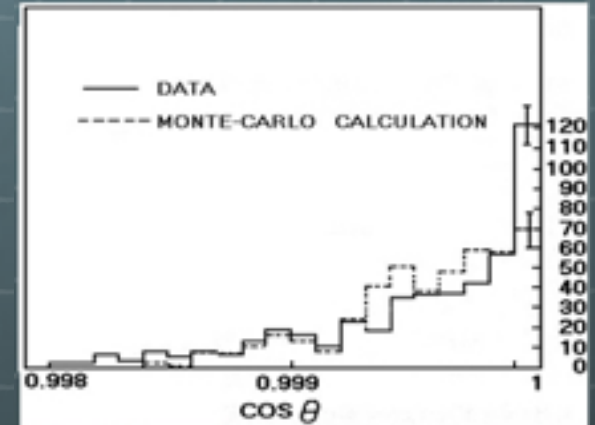
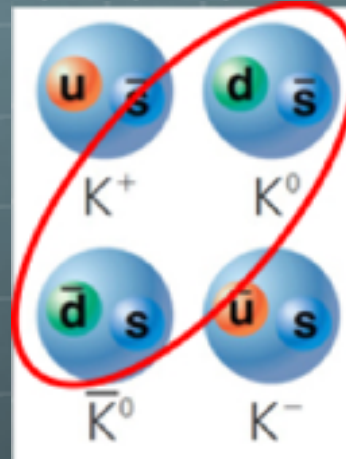
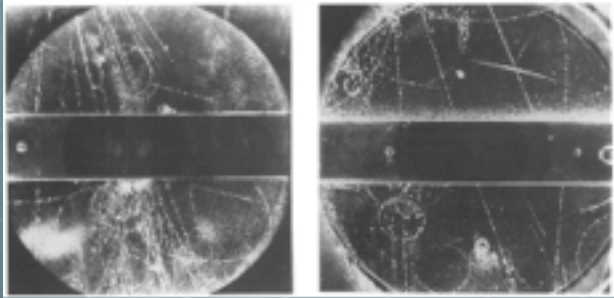
$K_L \rightarrow \pi^0 \nu \nu$ Experiment

G.Y.Lim
IPNS/KEK







2015

Studies on kaon decays

- Cornerstone of the SM -



Search for new physics

-  s-quark
-  Flavor structure
-  $\sim 500 \text{ MeV}/c^2$, $\sim 10\text{m}$ in $c\tau$
-  High intensity secondary kaon beam
-  Stopped and in-flight decays of kaons
-  Small number of decay products

(parts of) K_L^0 DECAY MODES

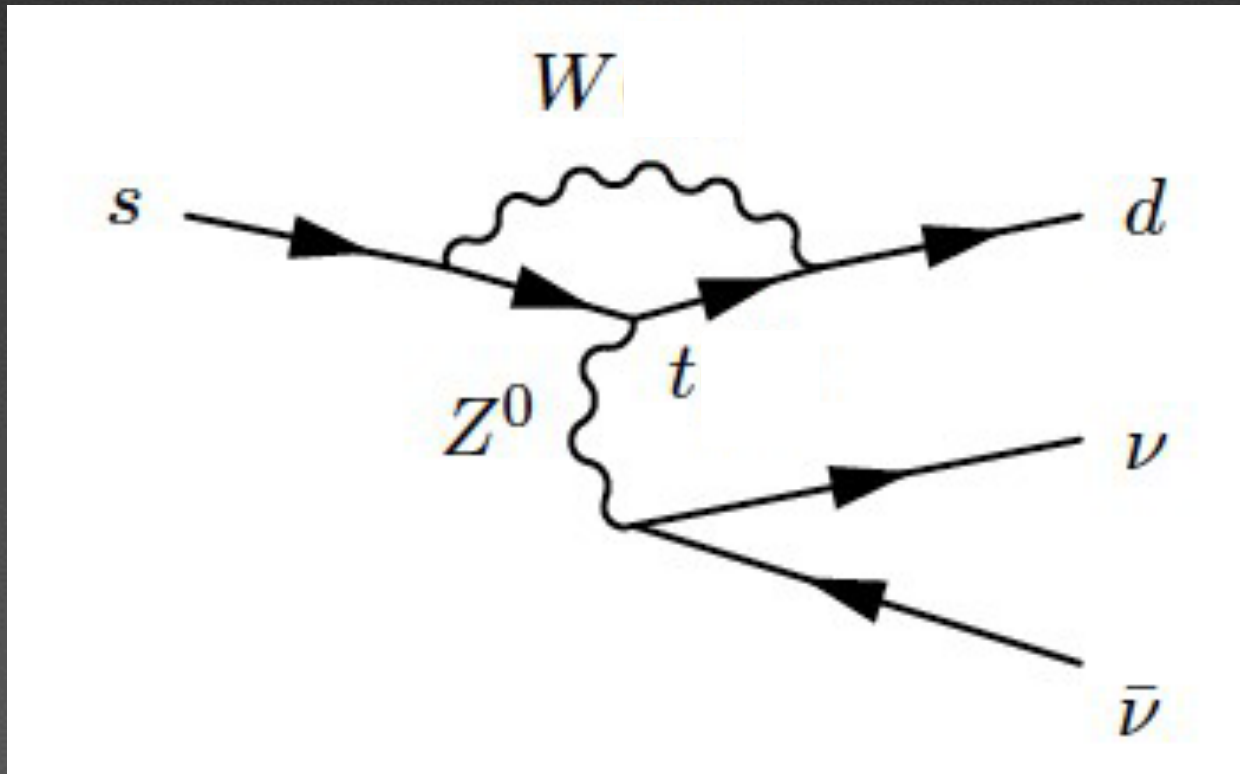
Charge conjugation \times Parity (CP) or Lepton Family number (LF) violating modes, or $\Delta S = 1$ weak neutral current ($S1$) modes

$\mu^+ \mu^-$	$S1$	$(6.84 \pm 0.11) \times 10^{-9}$
$e^+ e^-$	$S1$	$(9 \pm 4) \times 10^{-12}$
$\pi^+ \pi^- e^+ e^-$	$S1$ [r]	$(3.11 \pm 0.19) \times 10^{-7}$
$\pi^0 \pi^0 e^+ e^-$	$S1$	$< 6.6 \times 10^{-9}$
$\pi^0 \pi^0 \mu^+ \mu^-$	$S1$	$< 9.2 \times 10^{-11}$
$\mu^+ \mu^- e^+ e^-$	$S1$	$(2.69 \pm 0.27) \times 10^{-9}$
$e^+ e^- e^+ e^-$	$S1$	$(3.56 \pm 0.21) \times 10^{-8}$
$\pi^0 \mu^+ \mu^-$	$CP, S1$ [s]	$< 3.8 \times 10^{-10}$
$\pi^0 e^+ e^-$	$CP, S1$ [s]	$< 2.8 \times 10^{-10}$
$\pi^0 \nu \bar{\nu}$	$CP, S1$ [t]	$< 2.6 \times 10^{-8}$
$\pi^0 \pi^0 \nu \bar{\nu}$	$S1$	$< 8.1 \times 10^{-7}$
$e^\pm \mu^\mp$	LF [o]	$< 4.7 \times 10^{-12}$
$e^\pm e^\pm \mu^\mp \mu^\mp$	LF [o]	$< 4.12 \times 10^{-11}$
$\pi^0 \mu^\pm e^\mp$	LF [o]	$< 7.6 \times 10^{-11}$
$\pi^0 \pi^0 \mu^\pm e^\mp$	LF	$< 1.7 \times 10^{-10}$

<http://pdg.lbl.gov/2015/tables/rpp2015-tab-mesons-strange.pdf>

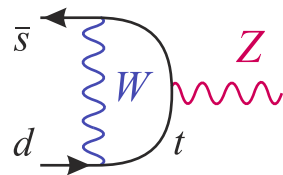
Rare precess : $\Delta t \cdot \Delta E \sim \hbar$

$$K_L \rightarrow \pi^0 \nu \bar{\nu} \text{ \& \> } K^+ \rightarrow \pi^+ \nu \bar{\nu}$$

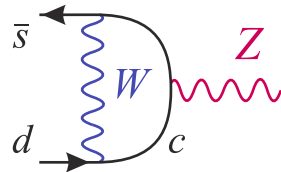


Flavor Changing Neutral Current

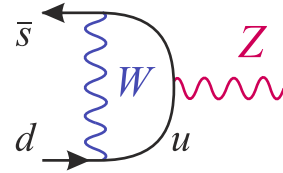
Highly suppressed process (in SM)



$$\frac{m_t^2}{M_W^2} (\text{Re } V_{ts}^\dagger V_{td} \sim \lambda^5)$$



$$\frac{m_c^2}{M_W^2} (\text{Re } V_{cs}^\dagger V_{cd} \sim \lambda)$$



$$\frac{m_u^2}{M_W^2} (\text{Re } V_{us}^\dagger V_{ud} \sim \lambda)$$

$$K^+ \rightarrow \pi^+ \nu \bar{\nu}$$

$$K_2 \rightarrow \pi^0 \nu \bar{\nu}$$

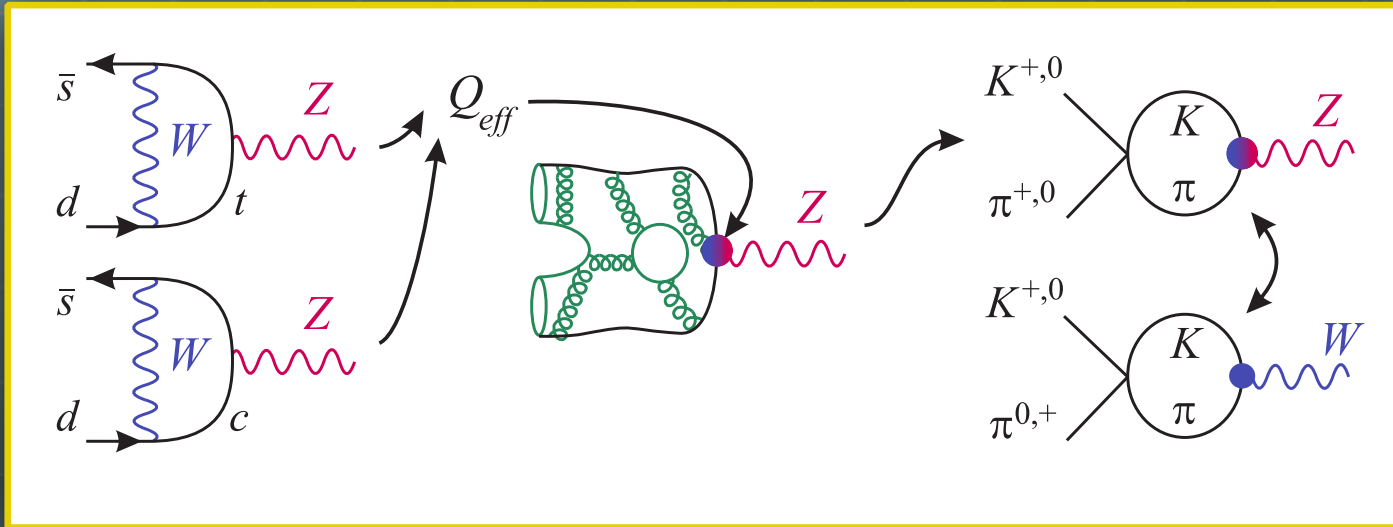
$$\frac{m_t^2}{M_W^2} (\text{Im } V_{ts}^\dagger V_{td} \sim \lambda^5)$$

$$\frac{m_c^2}{M_W^2} (\text{Im } V_{cs}^\dagger V_{cd} \sim \lambda^5)$$

$$\frac{m_u^2}{M_W^2} (\text{Im } V_{us}^\dagger V_{ud} = 0)$$

C. Smith, arXiv:1409.6162

Hadronic matrix elements



$$K \rightarrow \pi \nu \bar{\nu}$$

$$K \rightarrow \pi l \bar{l}$$

C. Smith, arXiv:1409.6162

The cleanest mode in theory

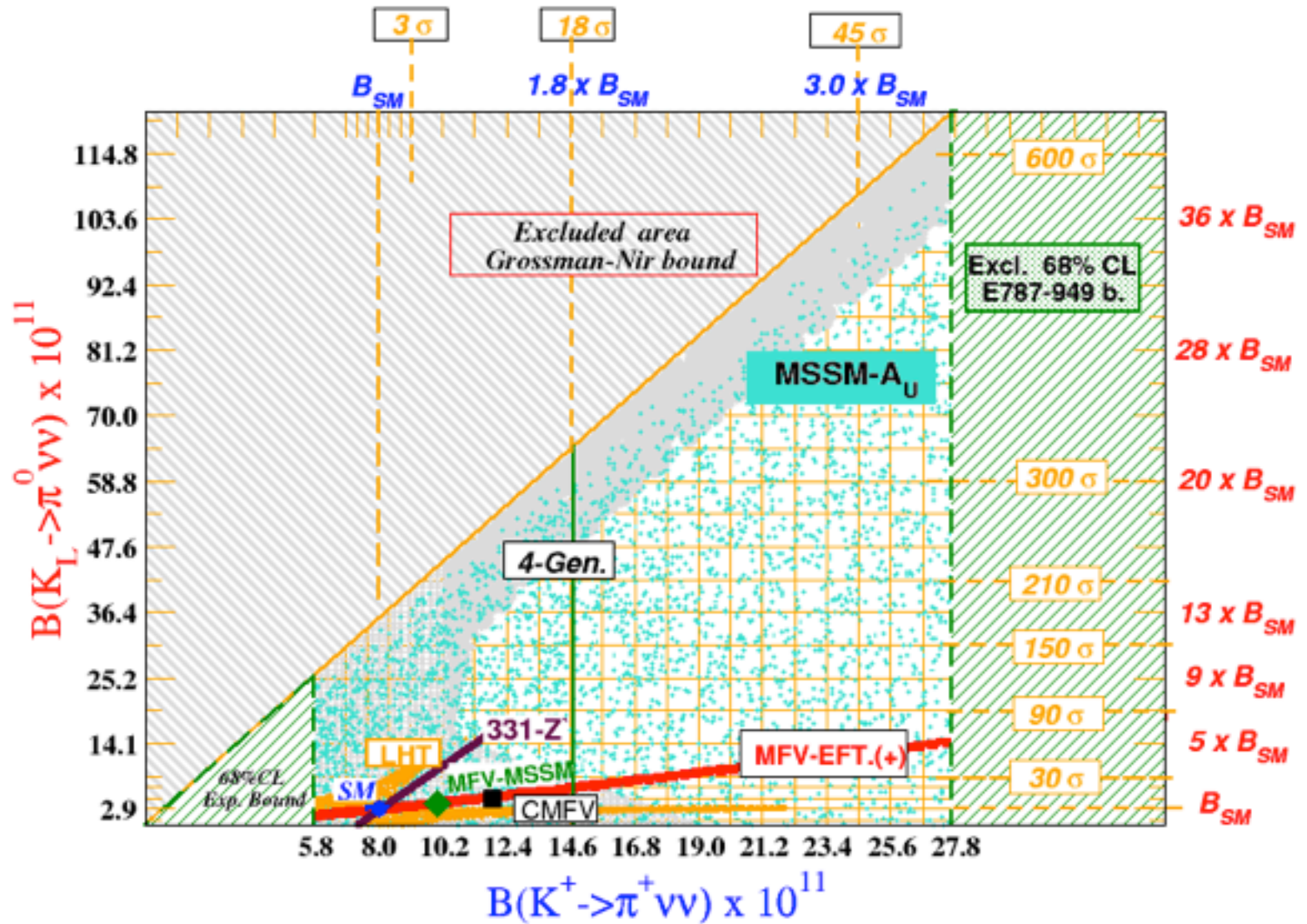
$$\mathcal{B}(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = \kappa_+ (1 + \Delta_{\text{EM}}) \cdot \left[\left(\frac{\text{Im} \lambda_t}{\lambda^5} X(x_t) \right)^2 + \left(\frac{\text{Re} \lambda_c}{\lambda} P_c(X) + \frac{\text{Re} \lambda_t}{\lambda^5} X(x_t) \right)^2 \right]$$

$$\longrightarrow (9.11 \pm 0.72) \times 10^{-11}$$

$$\mathcal{B}(K_L \rightarrow \pi^0 \nu \bar{\nu}) = \kappa_L \cdot \left(\frac{\text{Im} \lambda_t}{\lambda^5} X(x_t) \right)^2$$

$$\longrightarrow (3.00 \pm 0.31) \times 10^{-11}$$

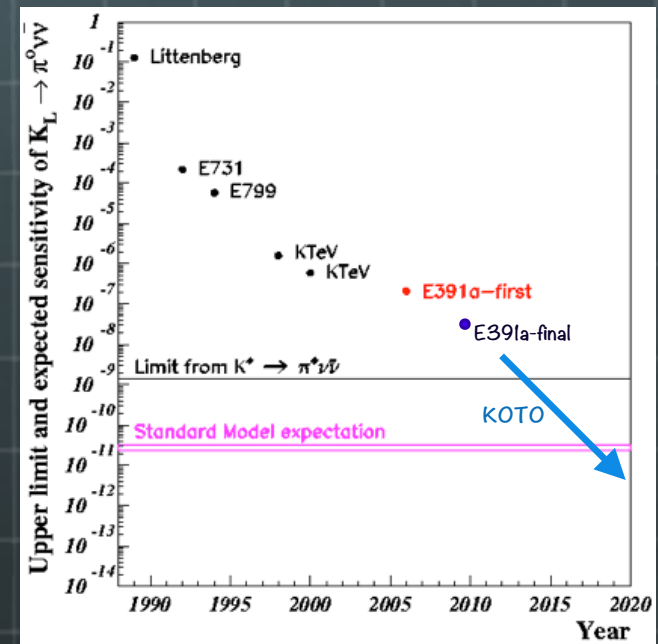
A. Buras, arXiv:1503.02693



Signal of the $K_L \rightarrow \pi^0 \nu \bar{\nu}$ Decay

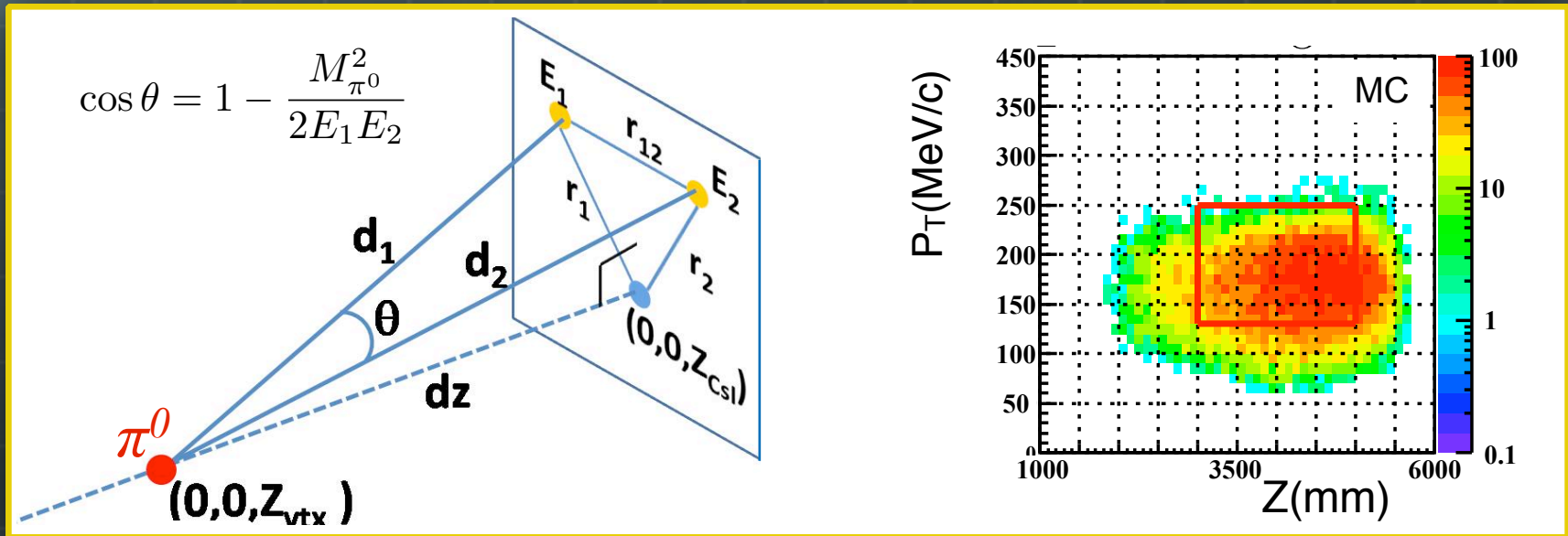
- No information of incident K_L
- Only decaying particle is K_L in the neutral beam.
- Momentum distribution can be obtained by using monitoring modes such as $K_L \rightarrow \pi^0 \pi^0 \pi^0$

- One clear π^0 and only one
 - Properly reconstruct π^0
 - $\pi^0 \rightarrow e^+ e^- \gamma$, $\pi^0 \rightarrow \gamma \gamma$
 - No any other decay products
 - Hermetic veto detector

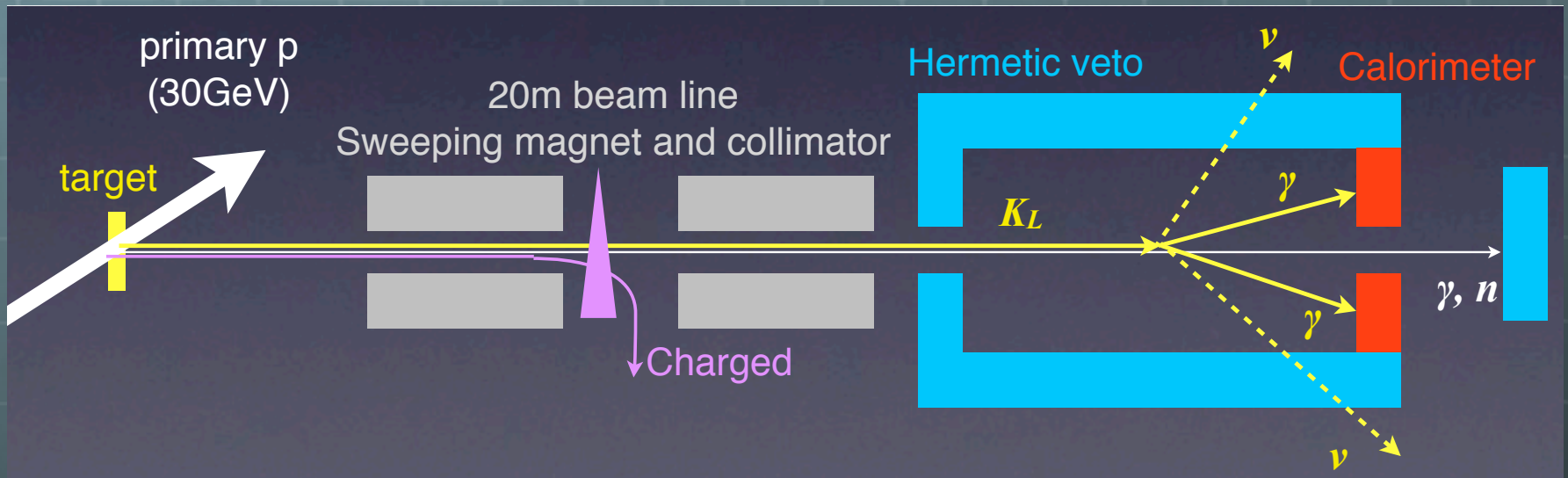


π^0 reconstruction

- EM calorimeter provides energies and incident positions of two photons
- With an assumption that π^0 decays at beam center and π^0 rest mass, we can obtain a distance between calorimeter and decaying vertex.
- Pair of neutrinos take away transverse momentum (P_T)

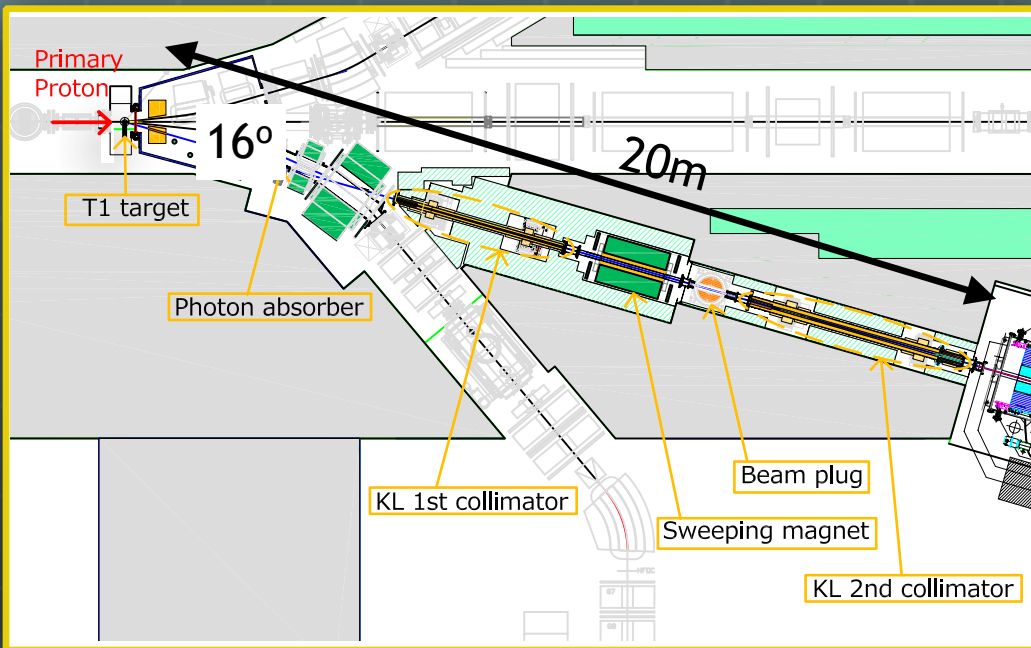


Experimental Method

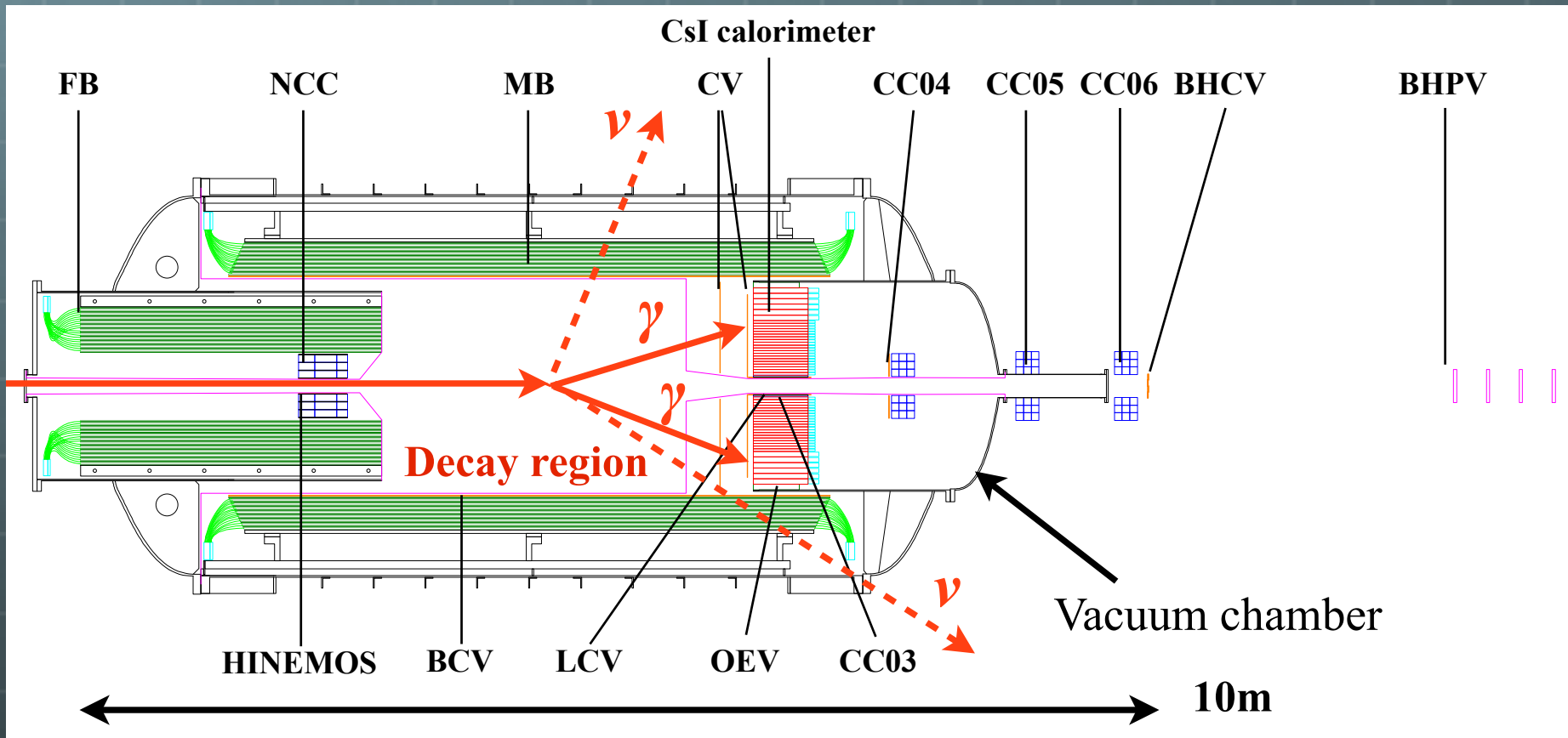


$2\gamma + \text{Nothing}$

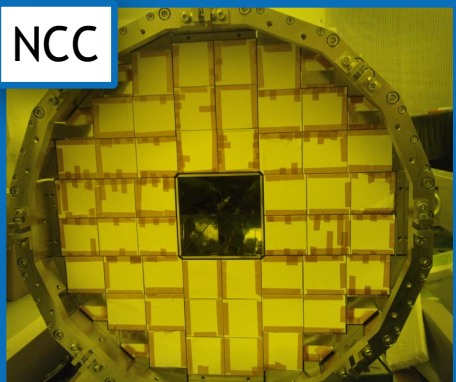
KL Beam Line



Detector



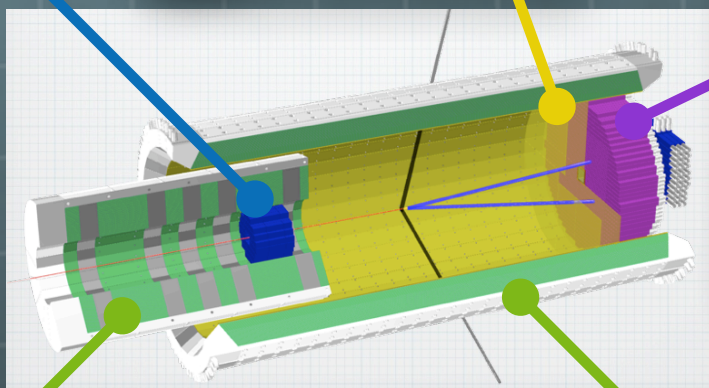
NCC



CV



CsI



FB



MB



Data taking in 2013

- January : Engineering run in vacuum
 - 11 kW X 1.5 days + 15 kW X 6.5 days
- March-April: Tuning for physics run
 - 15 kW X 5 days
- May: Physics run
 - 20 kW X 1 day + 24kW X 4 days (~ 100 hours)

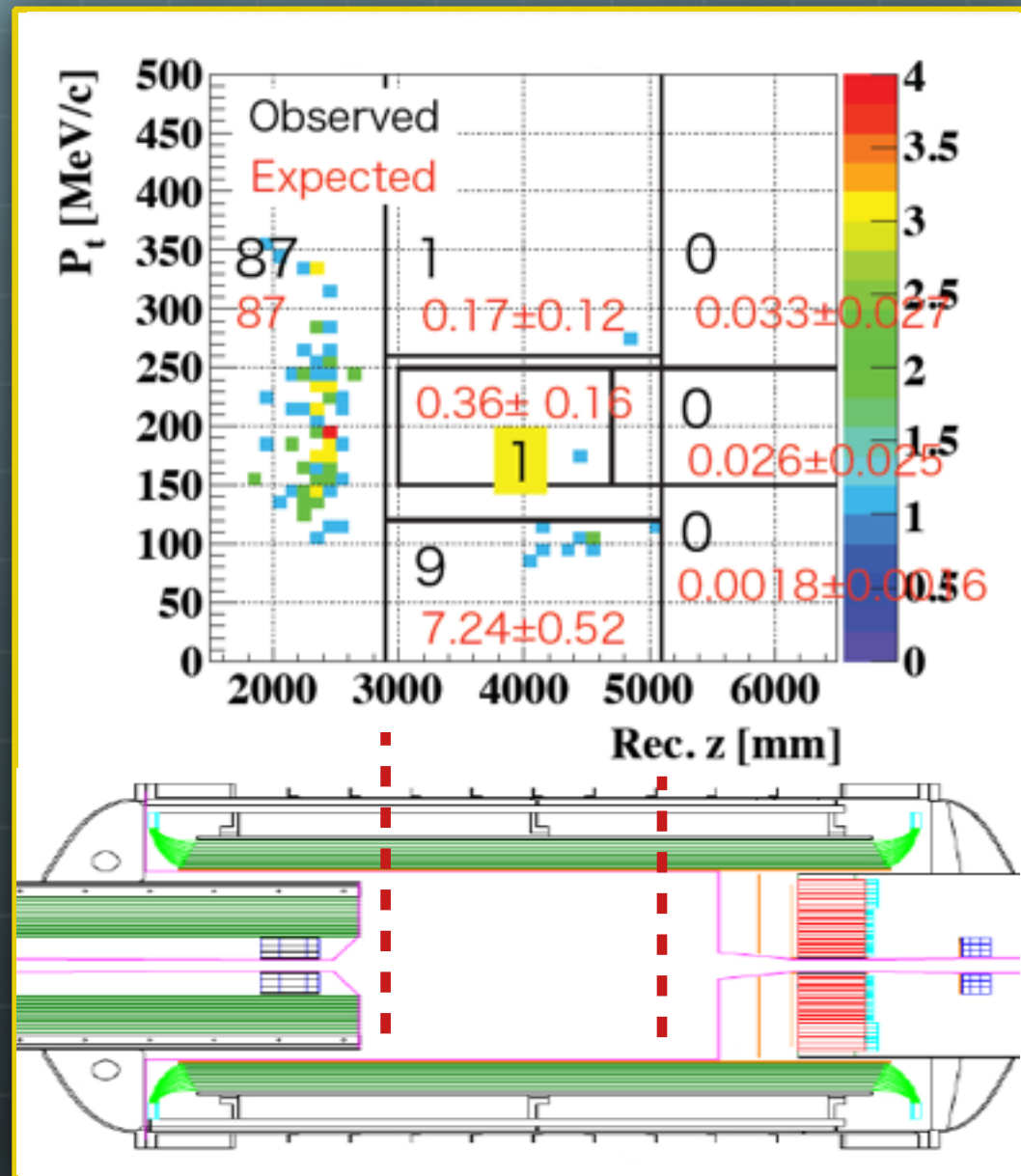
Results of May 2013

Removed B.G. events learned from the E391a. (π^0 production at the detectors)

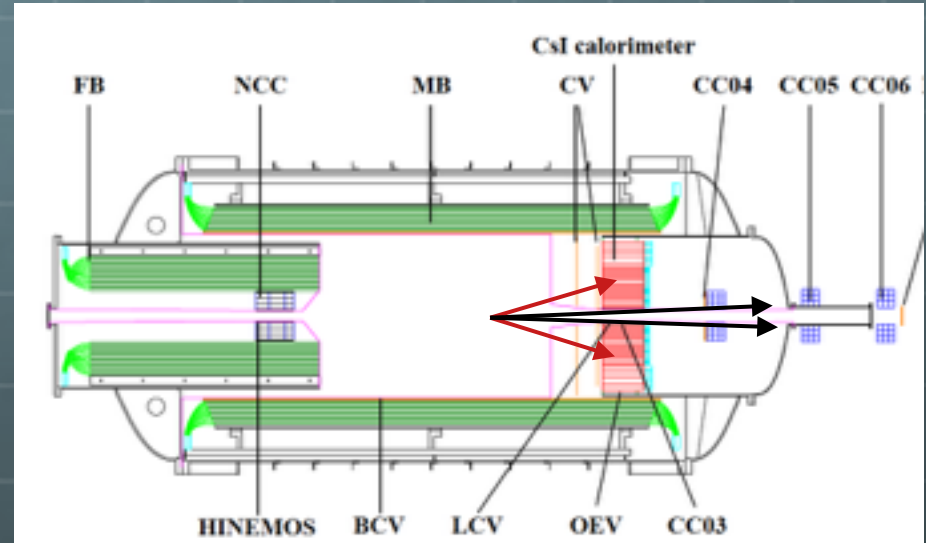
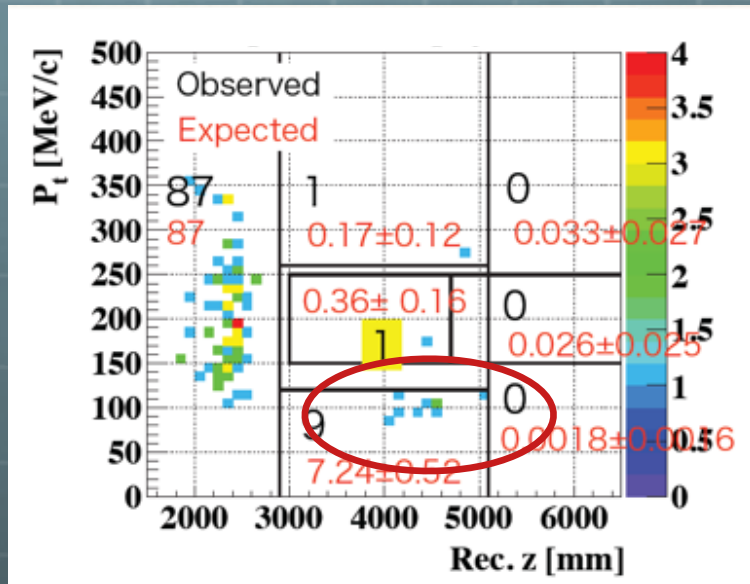
We found two new sources of the B.G.

Upgraded detector for run 2015

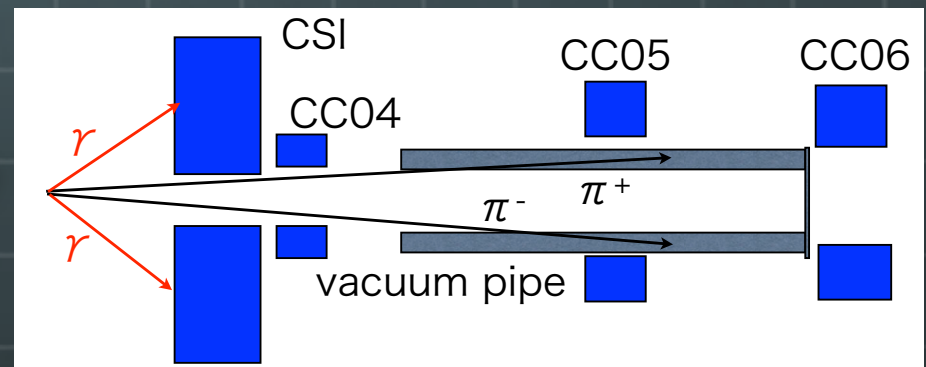
$$\text{S.E.S} = 1.29 \times 10^{-8}$$



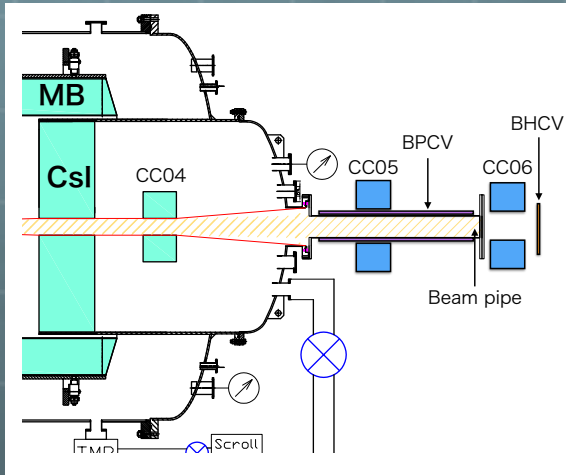
$$K_L \rightarrow \pi^+ \pi^- \pi^0$$



π^+ and/or π^- were not detected due to interaction inside vacuum pipe



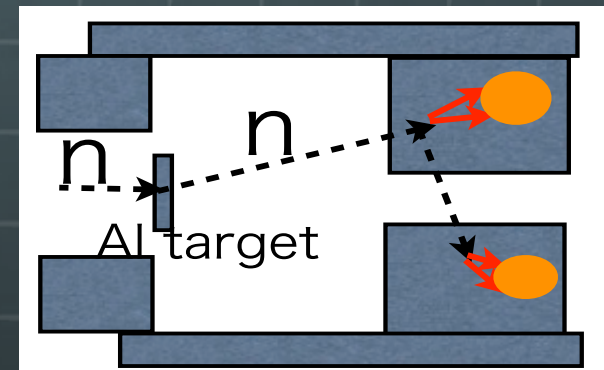
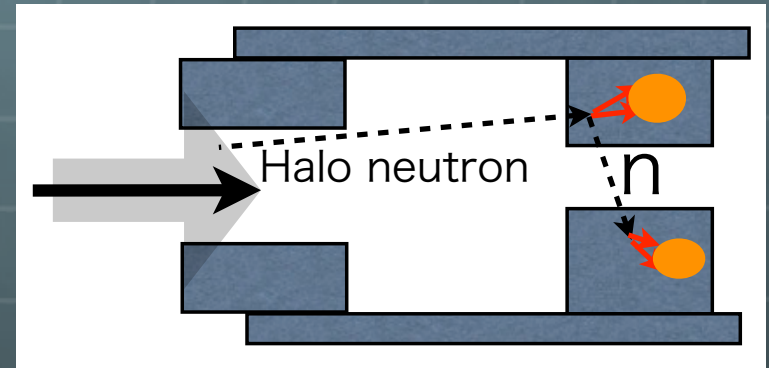
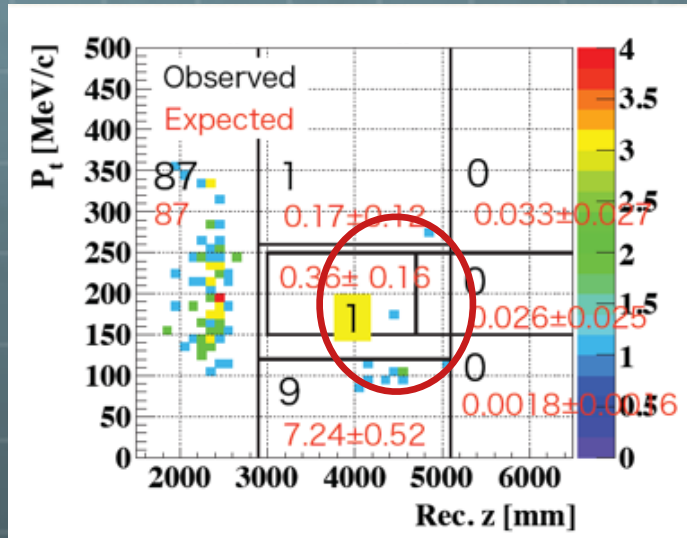
$$K_L \rightarrow \pi^+ \pi^- \pi^0$$



- Beam pipe with lighter material
- Stainless Steel -> Aluminum
- Beam Pipe Charged Veto
- 5mm-thick Plastic Scintillator
- Wavelength shifting fiber readout

1/60 additional reduction

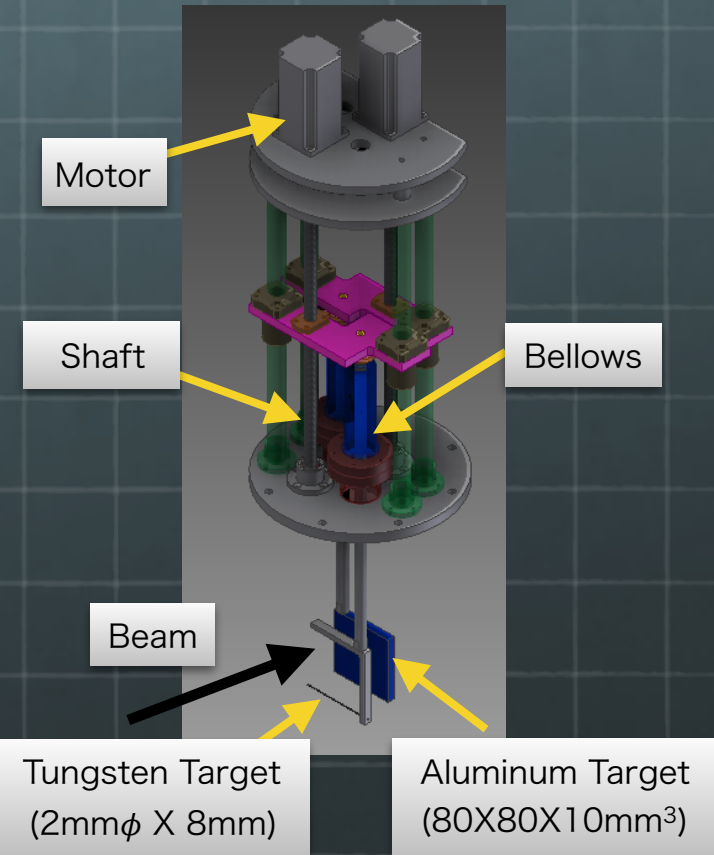
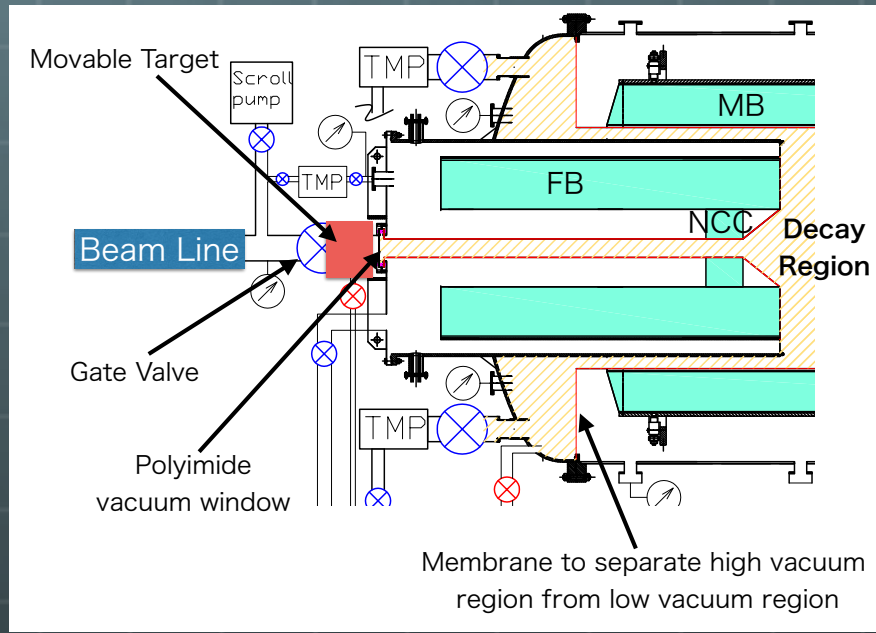
Halo neutron events



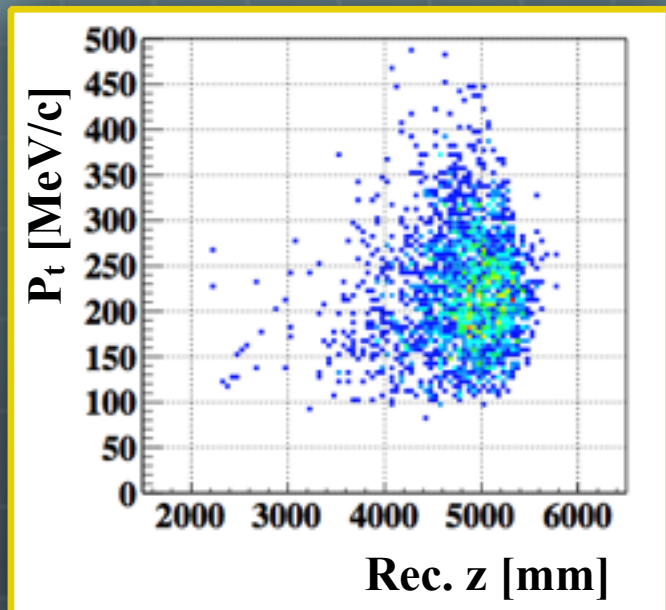
- Single neutron produce two clusters
- Newly founded background source
- Studied by using aluminum target data

Halo neutron events

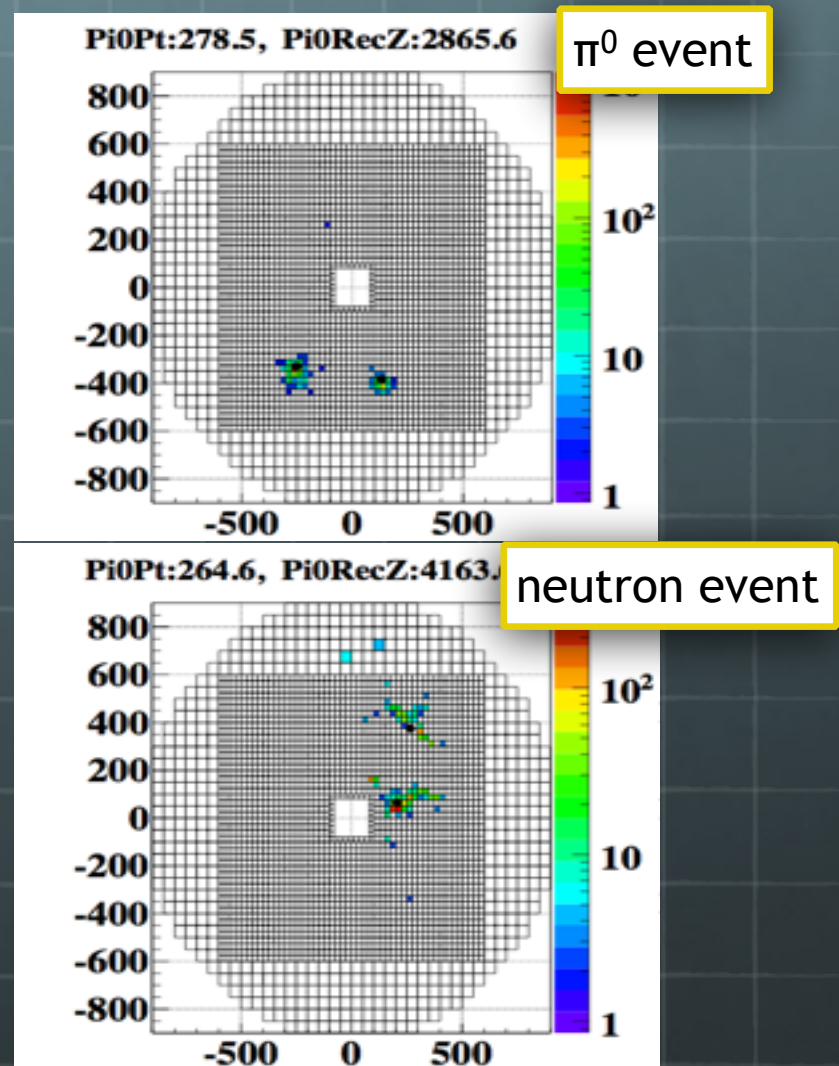
- 🌐 To reduce scattering source
- 🌐 To take data for enhanced neutron events



Enhanced neutron events

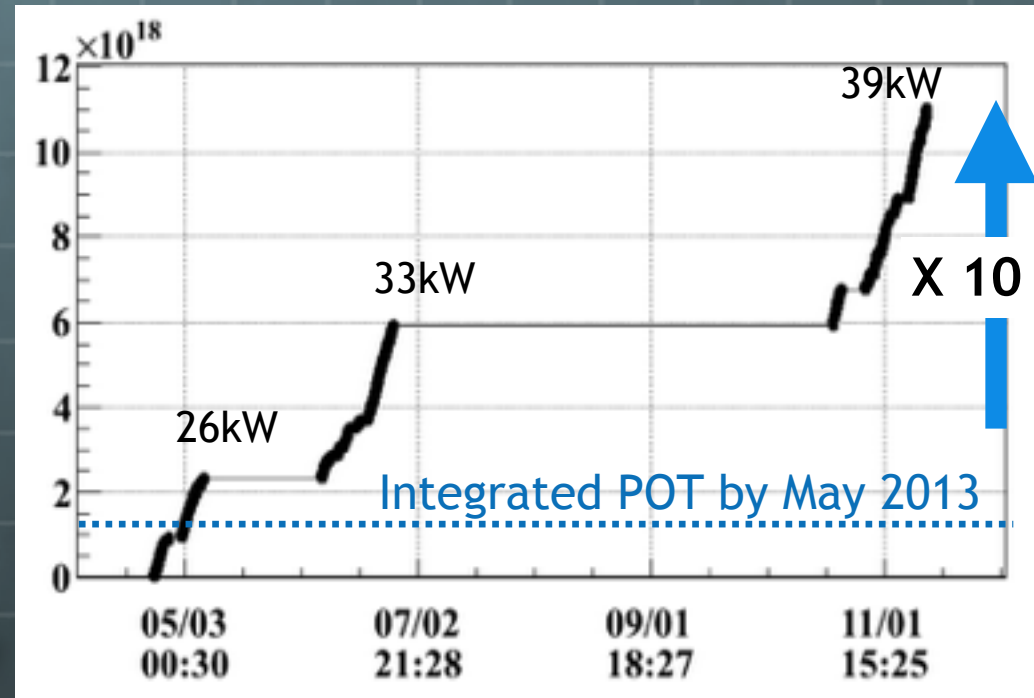


- 70-hour data taking with Al-target (>15 times more than May 2013)
- To study cluster and pulse shape in the calorimeter
- To develop a method to discriminate neutron induced events from the π^0 events

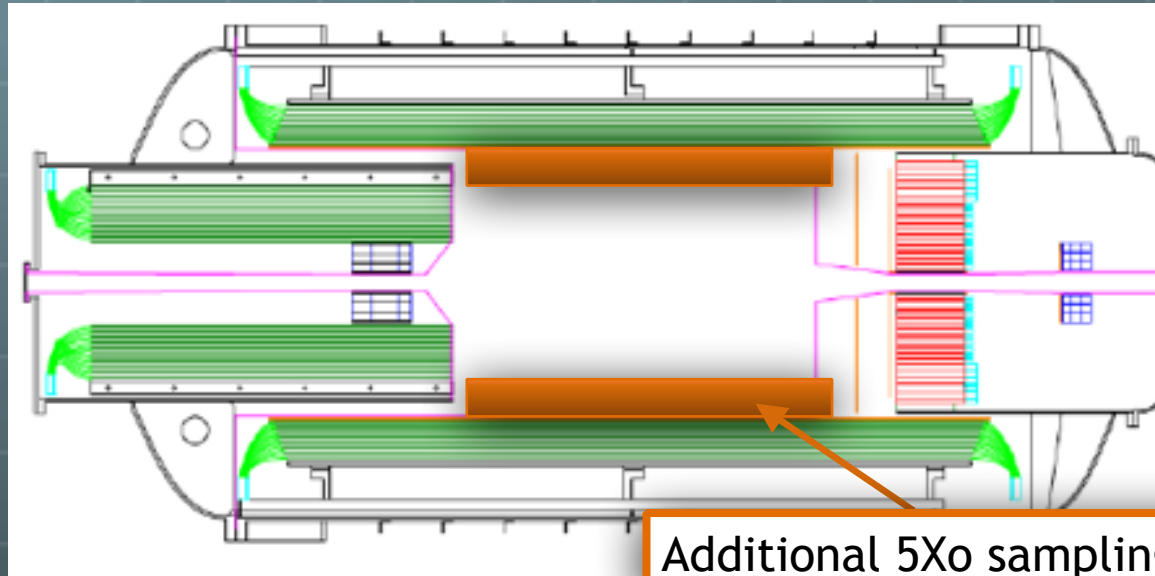


Run in 2015

- 10 times larger POT than run in May 2013
- Newly installed detectors work properly
- Detailed performance check in undergoing
- With the data collecting in coming beam time (Oct.~ Dec.), we aim at across the Grossman-Nir limit.

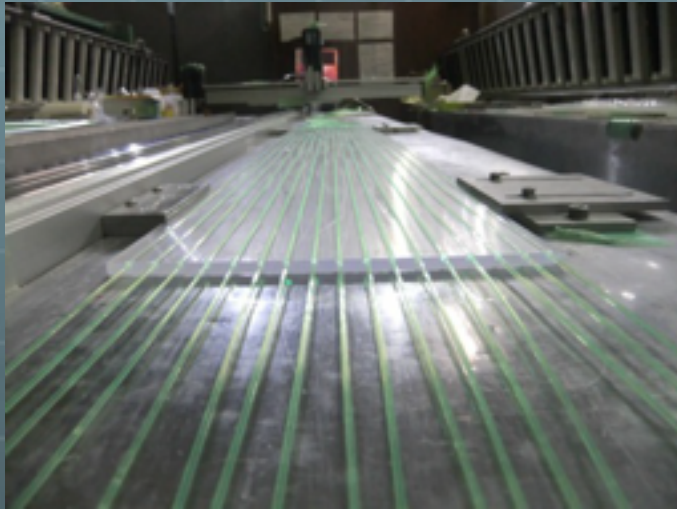


Inner Barrel



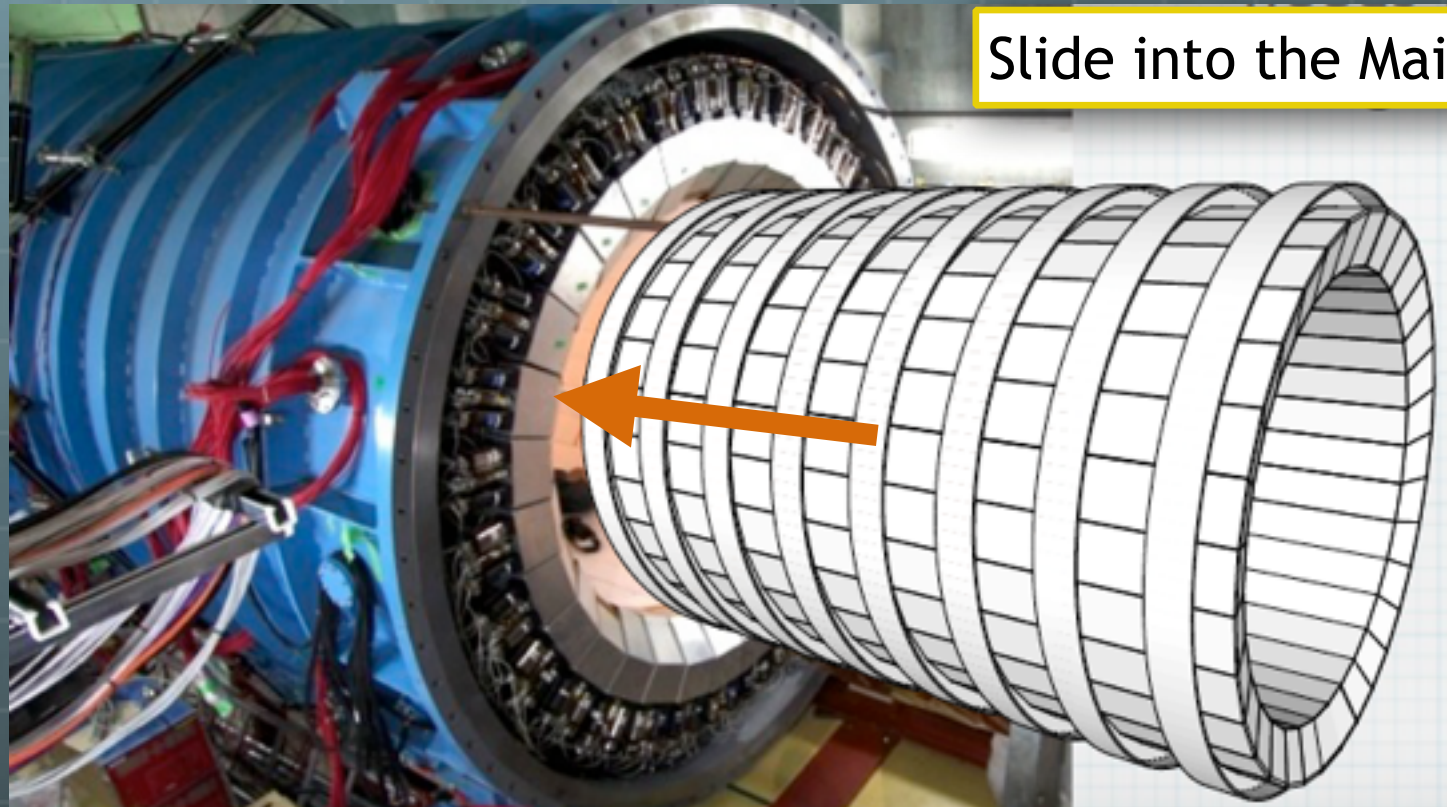
- Additional reduction for detection inefficiency due to punch-through as a factor of 50
- Increase better visible ration portion
- Reduction of background as a factor of three

Inner Barrel



- 🌐 Alternation lead sheet (1mm) and plastic scintillator(5mm)
- 🌐 Wave length shifting fiber read-out (BCF-92, $\phi 1.5\text{mm}$)

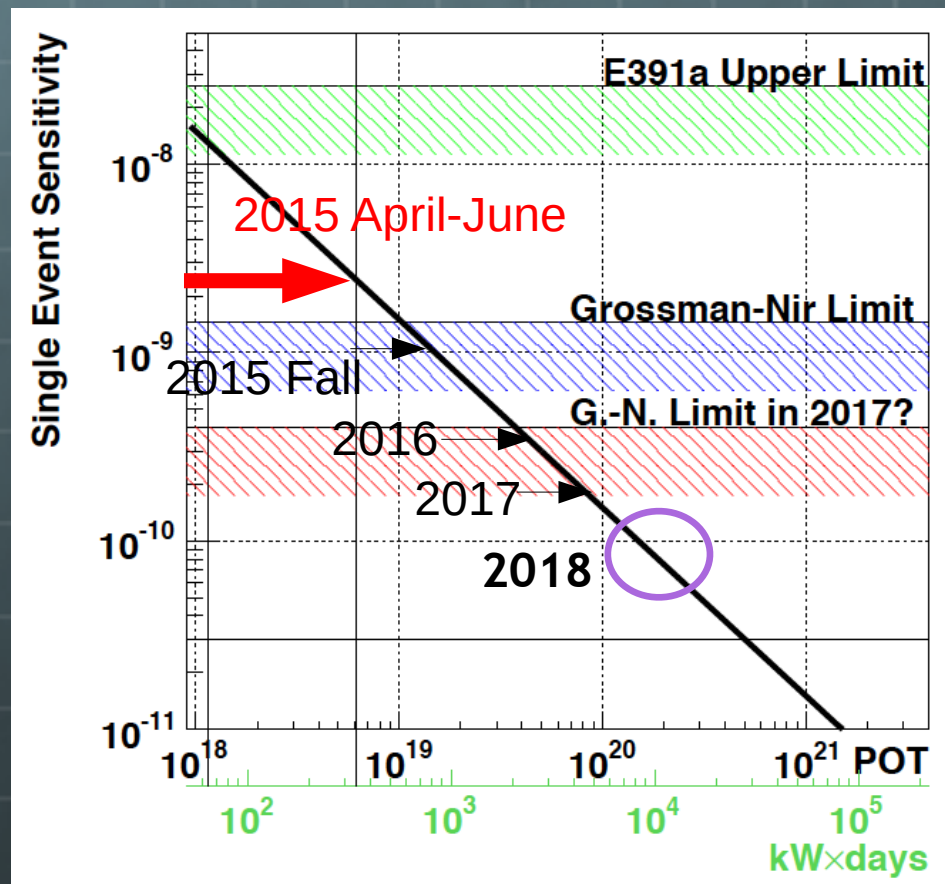
Inner Barrel









- 🌐 Fabrication of the 32 modules is undergoing
- 🌐 Assembling of the modules is preparing to start this month
- 🌐 Installation will be done in middle of next January

3 years from now

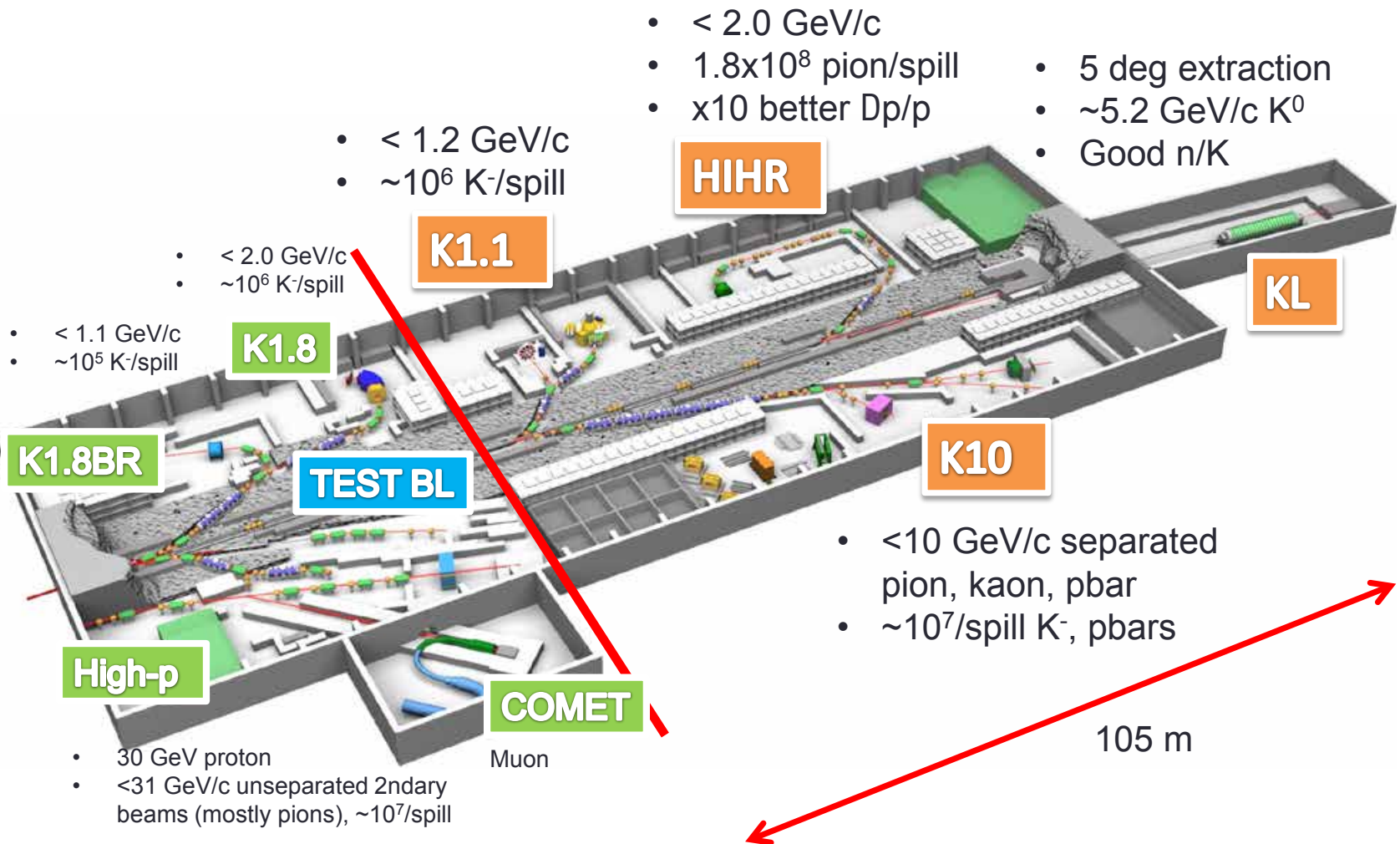
- Up-graded detector System
- Higher beam intensity of J-PARC



In near future

-  To perform precise measurement of the branching ratio (KOTO II)
-  It is worth to design an experiment to determine branching ratio with comparable uncertainty to that of the theoretical calculation
-  For the experiment
 -  Correct understanding all background sources
 -  Higher beam intensity
 -  Larger detector acceptance

Hadron Hall Extension



One solution

Study in 2014

5 deg extraction

36m beamline

8usr solid angle

4m-diamter calorimeter

20m decay volume

~100kW x 3e7 sec

KOTO

16deg

20m

7.8usr

2m

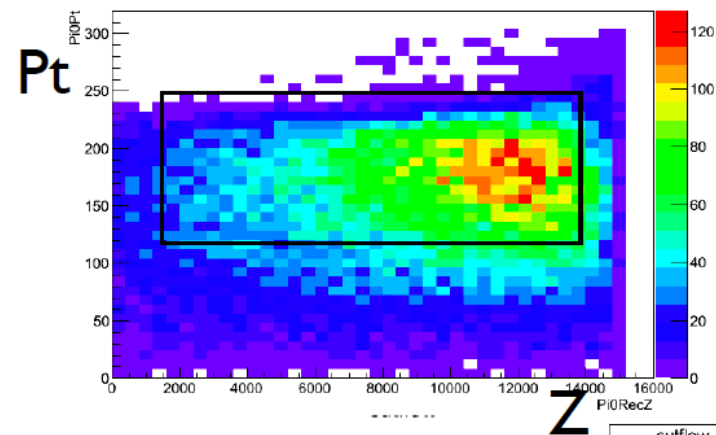
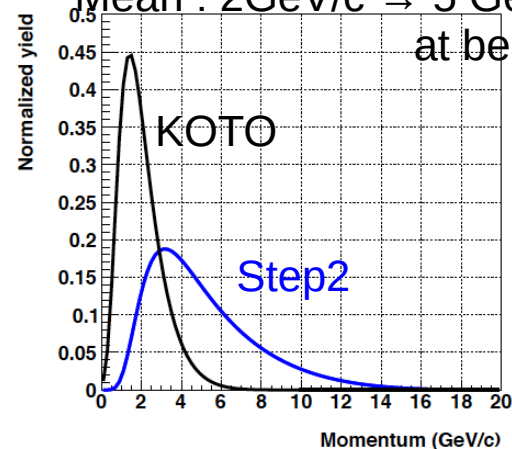
3.3m

~350 signal with S/N ~3 (K decay)
 $\Delta\text{BR}/\text{BR} \sim 6\%$

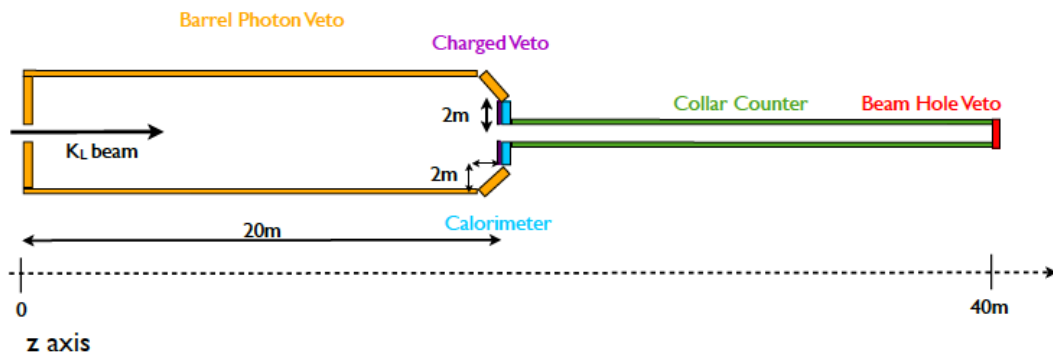
Peak : 1.4 GeV/c \rightarrow 3 GeV/c

Mean : 2 GeV/c \rightarrow 5 GeV/c

at beam exit



beam power	109kW
proton energy	30GeV
proton intensity	$7.5 \times 10^{13}/\text{spill}$
Spill length / Beam repetition	0.7s/3.3s
Target	Ni(1 λ)
Absorber	Pb(7cm)
Extraction angle	5°



100 events/ 10^7 s ?

- # of events = Branching ratio * # of K-decay * Acceptance
- For $Br = 3 \times 10^{-11}$, Acceptance = 0.1
 - Needed K-decay for 100 events : 3.3×10^{14} (33MHz)
 - Event definition time as 10ns
 - Accidental coincidence = 0.33
 - Increasing of instantaneous rate as 3 : Always two decay inside detector
 - Accidental hits have longer decay region (factor of 5 ?)
- Totally new approach is needed

Summary

- Measurement of the $K_L \rightarrow \pi^0 \nu \nu$ decay will enable us to find new physics effect.
- KOTO is aiming for the first event.
- High statistics experiment is needed even though lots of difficulties.
 - large number of KAONS - High intensity beam
 - large acceptance and enough background suppression.
 - Fully understand accidental activities.
- We started hadron hall extension
 - KOTO-II will be performed



ELSEVIER

10 April 1997

PHYSICS LETTERS B

Physics Letters B 398 (1997) 163–168

$K_L \rightarrow \pi^0 \nu \bar{\nu}$ beyond the Standard Model \star

Yuval Grossman^a, Yosef Nir^b

^a *Stanford Linear Accelerator Center, Stanford University, Stanford, CA 94309, USA*

^b *Department of Particle Physics, Weizmann Institute of Science, Rehovot 76100, Israel*

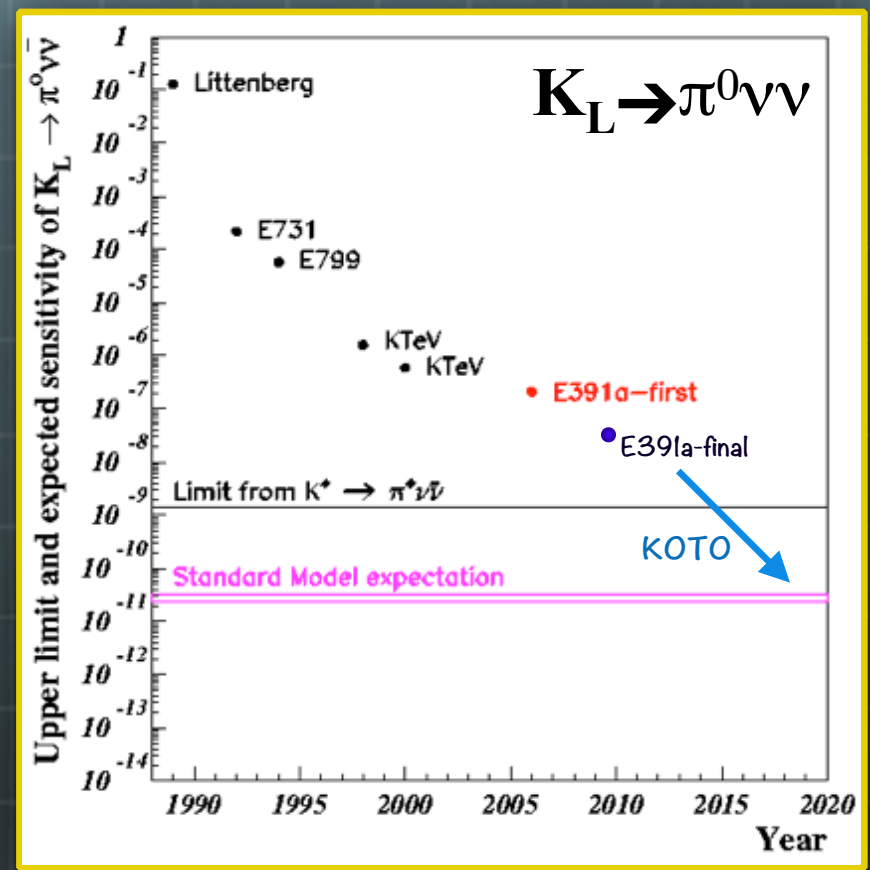
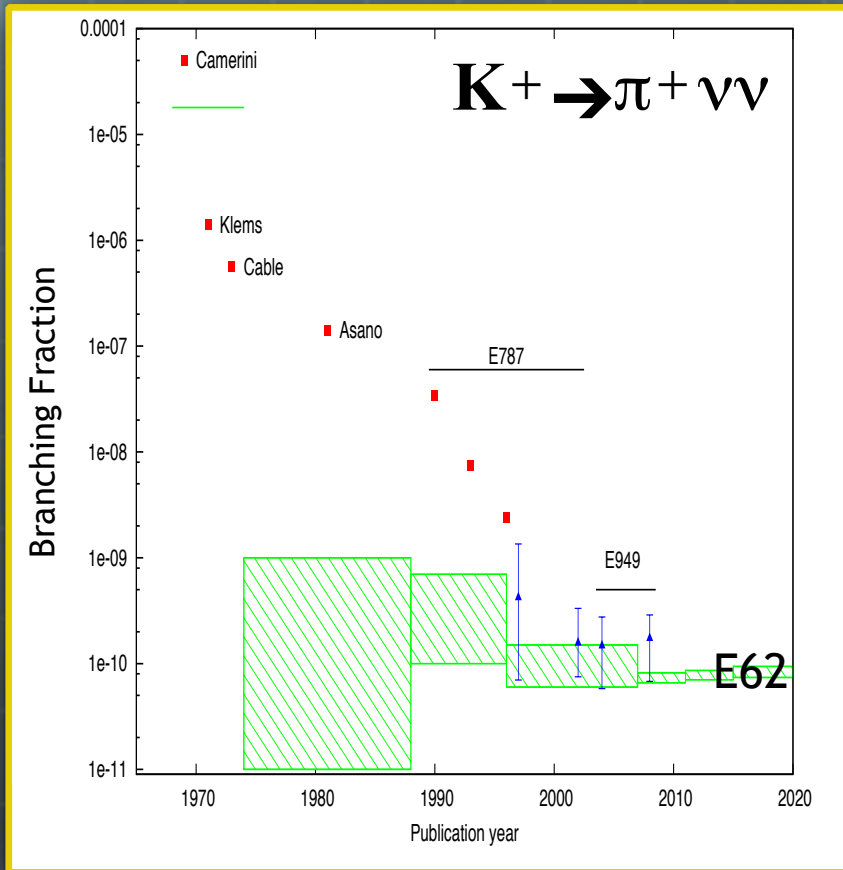
Received 29 January 1997

Editor: M. Dine

Abstract

We analyze the decay $K_L \rightarrow \pi^0 \nu \bar{\nu}$ in a model independent way. If lepton flavor is conserved the final state is (to a good approximation) purely CP even. In that case this decay mode goes mainly through CP violating interference between mixing and decay. Consequently, a theoretically clean relation between the measured rate and electroweak parameters holds in any given model. Specifically, $\Gamma(K_L \rightarrow \pi^0 \nu \bar{\nu}) / \Gamma(K^+ \rightarrow \pi^+ \nu \bar{\nu}) = \sin^2 \theta$ (up to known isospin corrections), where θ is the relative CP violating phase between the $K - \bar{K}$ mixing amplitude and the $s \rightarrow d \nu \bar{\nu}$ decay amplitude. The experimental bound on $\text{BR}(K^+ \rightarrow \pi^+ \nu \bar{\nu})$ provides a model independent upper bound: $\text{BR}(K_L \rightarrow \pi^0 \nu \bar{\nu}) < 1.1 \times 10^{-8}$. In models with lepton flavor violation, the final state is not necessarily a CP eigenstate. Then CP conserving contributions can dominate the decay rate. © 1997 Published by Elsevier Science B.V.

Experimental Status



Still far from the goal ...