Measurements of the elastic, inelastic and total *pp* cross sections with the ATLAS, CMS and TOTEM detectors

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Outline:

- Physics motivation.
- Inelastic cross section at $\sqrt{s} = 13$ TeV with the ATLAS and CMS detectors.
- Elastic, inelastic and total cross sections at $\sqrt{s} =$ 8 TeV with the ATLAS and TOTEM detectors.

- The elastic (σ_{el}), inelastic (σ_{inel}) and total (σ_{tot}) *pp* cross sections are fundamental quantities which cannot be calculated with perturbative QCD.
- Regge theory provides a description but data is needed to constrain models.
- σ_{tot} gives the upper bound on any *pp* process and is seen to rise with collision energy.
- A substantial fraction of σ_{inel} is diffractive processes. Measurement of σ_{inel} will constrain models also for cosmic-ray shower in the atmosphere.



Inelastic cross section measurement at $\sqrt{s} = 13$ TeV with ATLAS and CMS

arXiv:1606.02625 CMS-PAS-FSQ-15-005

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Elastic, inelastic and total pp cross sections

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- Strategy: Measure the inelastic cross section in a fiducial region and extrapolate to full phase-space with input from theoretical models.
 - The better detector coverage, the smaller extrapolation uncertainty.
- The fiducial σ_{inel} is the number of observed events corrected for background, pile-up, efficiencies and luminosity.

ATLAS:

 MBTS plastic scintillators at z = ±3.6 m covering 2.07 < |η| < 3.86.

WLS fibers

CMS:

- HF calorimeter of iron absorbers and quartz fibers covering $3.0 < |\eta| < 5.2$.
- CASTOR calorimeter of tungsten and quartz covering $-6.6 < \eta < -5.2$.



Inelastic cross section - Tuning models (ATLAS)

- The inelastic cross section is the sum of the non-diffractive and the diffractive cross section.
- The ratio $f_D = (\sigma_{SD} + \sigma_{DD})/\sigma_{inel}$ is poorly known and differs between models.
- The fraction of single-sided events, *R*_{SS}, is related to *f*_D and used to tune *f*_D in the models.

- Using the *f*_D-tuned models, the hit multiplicity in the MBTS for the models are compared to data:
 - The DL (Donnachie-Landshoff) pomeron flux model is best.
 - The EPOS LHC and QGSJET-II models (developed for cosmic-ray showering) are worst.
- CMS checks that models predict correct ratio of cross sections between the HF-only and the HF+CASTOR phase-space regions.



CMS:

- The average of the model extrapolation factors is used to go from fiducial to full phase-space cross section.
- Maximum difference between models is used as systematic uncertainty.

ATLAS:

• A precise (and independent) measurement of σ_{inel} at $\sqrt{s}=$ 7 TeV is used:

$$\sigma_{\text{inel}} = \sigma_{\text{inel}}^{\text{fid}} + (\sigma_{\text{inel}, \ 7 \ \text{TeV}}^{\text{ALFA}} - \sigma_{\text{inel}, \ 7 \ \text{TeV}}^{\text{fid}}) \times \frac{\sigma_{\text{inel}}^{\text{MC}}(\xi < 10^{-6})}{\sigma_{\text{inel}, \ 7 \ \text{TeV}}^{\text{MC}}(\xi < 5 \times 10^{-6})}$$

• The difference between Pythia8 DL and Pythia8 MBR is used as systematic uncertainty.

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Inelastic cross section - Results

CMS:
$$\sigma_{inel}^{fid,HF} = 65.8 \pm 0.8(exp.) \pm 1.8(lum.) \text{ mb}$$

$$\sigma_{inel}^{fid,HF+CASTOR} = 66.9 \pm 0.4(exp.) \pm 2.0(lum.) \text{ mb}$$

$$\sigma_{inel} = 71.3 \pm 0.5(exp.) \pm 2.1(lum.) \pm 2.7(ext.) \text{ mb}$$

ATLAS:
$$\sigma_{inel}^{fid} = 68.1 \pm 0.6(exp.) \pm 1.3(lum.) \text{ mb}$$

$$\sigma_{inel} = 79.3 \pm 0.6(exp.) \pm 1.3(lum.) \pm 2.5(ext.) \text{ mb}$$

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ξ > 10⁻⁶

 σ_{inel}

 $\xi_{\rm X} > 10^{-7} \text{ or } \xi_{\rm Y} > 10^{-6}$

Elastic, inelastic and total cross sections measurement at $\sqrt{s} = 8$ TeV with ATLAS and TOTEM

Phys. Lett. B (2016) 158 Phys. Rev. Lett. 111, 012001 (2013) Nucl. Phys. B 899 (2015) 527-546 CERN-PH-EP-2015-325 • From the optical theorem we get:

$$\begin{split} \sigma_{\text{tot}}^2 &= \frac{16\pi(\hbar c)^2}{1+\rho^2} \frac{1}{L} \frac{dN_{\text{el}}}{dt} \Big|_{t=0} \quad \text{with} \quad \rho = \frac{\text{Re}\left[F_{\text{el}}(t)\right]}{\text{Im}\left[F_{\text{el}}(t)\right]} \quad \text{(ATLAS)} \\ \text{equiv.} \qquad \sigma_{\text{tot}} &= \frac{16\pi(\hbar c)^2}{1+\rho^2} \frac{1}{N_{\text{el}}+N_{\text{inel}}} \frac{dN_{\text{el}}}{dt} \Big|_{t=0} \quad \text{(TOTEM)}. \end{split}$$

• The four-momentum transfer t is calculated as

$$t=-(p\times\theta^*)^2.$$

where the scattering angle θ^* is calculated from the proton trajectories and *p* is the beam momentum.

- Data are taken in runs with low pile-up ($\mu \leq 0.1$) and high β^* collision optics since $t_{\min} \propto \frac{p^2}{\beta^*}$:
 - $\beta^* = 90$ m results from ATLAS (one dataset) and TOTEM (two datasets).
 - $\beta^* = 1$ km results from TOTEM (one dataset).

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- ATLAS and TOTEM use tracking detectors in Roman Pots at z ~ 230 m to approach outgoing beams in vertical direction.
 - ATLAS uses 10 \times 2 orthogonal layers of scintillating fibers giving \approx 30 μ m tracking resolution.
 - TOTEM uses a stack of 10 silicon strip detectors giving \approx 11 μ m tracking resolution.
- In addition, TOTEM has two tracking telescopes:
 - T1 is a cathode strip chamber at $z = \pm 9$ m covering $3.1 \le |\eta| \le 4.7$.
 - T2 is based on gas electron multiplier chambers at $z = \pm 13.5$ m covering $5.3 \le |\eta| \le 6.5$.



Elastic analysis - event selection

- Elastic events are selected when all four detectors in an arm have a track.
- The tracks are required to fulfill certain correlations between inner-outer stations and between left and right side.
- ATLAS $\beta^* = 90$ m: 3.8 M elastic events.
- TOTEM $\beta^* = 90$ m: 0.65 M elastic events.
- TOTEM $\beta^* = 1$ km: 0.35 M elastic events.





Elastic analysis - Experimental effects

- **Background** comes from beam halo, single diffractive and central diffractive protons.
- Fraction is \leq 0.12 % estimated from antigolden topology.



- Detector acceptance is highly dependent on detector distance to the beam and beam divergence.
- Found from simulation tuned to data.
- *t*-resolution is influenced by detector resolution and beam divergence.
- Relative *t*-resolution is better than 10 % and corrected for by unfolding.
- Track reconstruction inefficiency is data driven.



Elastic analysis - Theoretical prediction

 The differential elastic cross section is a superposition of the strong interaction amplitude F_N and the Coulomb amplitude F_C added in quadrature giving

$$\frac{\mathrm{d}\sigma_{\mathrm{el}}}{\mathrm{d}t} \propto \frac{G^4(t)}{|t|^2} + \sigma_{\mathrm{tot}}^2(1+\rho^2) \cdot \exp(-\mathbf{B}|t|) - \frac{\sigma_{\mathrm{tot}}G^2(t)}{|t|} \left[\sin(\phi(t)) + \rho\cos(\phi(t))\right] \cdot \exp\left(\frac{-\mathbf{B}|t|}{2}\right)$$

The corrected differential cross section is fitted with σ_{tot} and B as free parameters.

 $0.01 \le |t| \le 0.2 \text{ GeV}^2$ (1. dataset) $0.02 \le |t| \le 0.2 \text{ GeV}^2$ (2. dataset)

- TOTEM neglects the Coulomb and interference terms.
- TOTEM $\beta^* = 1$ km: $6 \cdot 10^{-4} \le |t| \le 0.2 \text{ GeV}^2$.



- TOTEM measures simultaneously the elastic and inelastic rate and is hence independent of luminosity.
- The inelastic rate is determined with the T1 (3.1 \leq $|\eta| \leq$ 4.7) and T2 (5.3 \leq $|\eta| \leq$ 6.5) telescopes, detecting about 95% of the inelastic rate.
- The strategy is similar to the CMS and ATLAS inelastic cross section measurement.
- Events are triggered by the T2 telescope and corrected for experimental effects.
- Uncertainty from extrapolation of fiducial region is dominating ("Low mass diffraction").

Source	Correction	Uncertainty	Effect on
Beam gas	0.45%	0.45%	all rates
Trigger efficiency	1.2%	0.6%	all rates
Pileup	2.8%	0.6%	all rates
T2 reconstruction	0.35%	0.2%	$N_{\text{inel}}, N_{ \eta < 6.5}$
"T1 only"	1.2%	0.4%	$N_{\rm inel}, N_{ \eta < 6.5}$
Internal gap covering T2	0.4%	0.2%	$N_{\rm inel}, N_{ \eta < 6.5}$
Central diffraction	0%	0.35%	$N_{\text{inel}}, N_{ \eta < 6.5}$
Low mass diffraction	4.8%	2.4%	$N_{\rm inel}$

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Elastic analysis - Results

σ_{el} is the integral of the nuclear part:

$$\sigma_{\rm el} = \frac{\sigma_{\rm tot}^2}{B} \frac{1+\rho^2}{16\pi(\hbar c)^2}.$$

•
$$\sigma_{\text{inel}} = \sigma_{\text{tot}} - \sigma_{\text{el}}$$
.

	ATLAS ($\beta^* = 90 \text{ m}$)	TOTEM ($\beta^* = 90 \text{ m}$)	TOTEM ($\beta^* = 1 \text{ km}$)
σ_{tot}	96.07 ± 0.92	101.7 ± 2.9	102.9 ± 2.3
σ_{el}	24.33 ± 0.39	27.1 ± 1.4	-
σ_{inel}	71.73 ± 0.71	74.7 ± 1.7	-

- The total cross section is still rising with energy.
- The difference between ATLAS and ۲ TOTEM $\beta^* = 90$ m corresponds to 1.9 σ .
- Using $\beta^* = 1$ km data, TOTEM also measured

$$\rho = 0.12 \pm 0.03$$

Model extrapolation from lower energies:

$$\rho=0.140\pm0.007$$



(fit not updated with latest ATLAS result)

Elastic, inelastic and total pp cross sections

• Hints of a slight deviation from an exponential fall-off of the elastic nuclear amplitude was reported at ISR and SppS, but not at the Tevatron.

• TOTEM used
$$\frac{d\sigma}{dt}(t) = \frac{d\sigma}{dt}\Big|_{t=0} \exp\left(\sum_{i=0}^{N_b} b_i t^i\right)$$
 to exclude a purely exponential form with more than $7\sigma!$

Nb	χ^2/ndf	σ_{tot}
1	117.5/28	(101.7 ± 2.9) mb
2	29.3/27	(101.5 ± 2.1) mb
3	25.5/26	(101.9 ± 2.1) mb

- This approach was impossible for ATLAS as the t-dependent systematic uncertainties are too large.
 - Average beam energy uncertainty dominates.
 - ATLAS uses a larger uncertainty than TOTEM.
 - An official value will be available later this year.



• ATLAS has tested different B-parametrizations giving an RMS of 0.28 mb on σ_{tot} .

Summary

• The inelastic cross section at $\sqrt{s} = 13$ TeV was measured by a MC extrapolation from the fiducial region:

$$\begin{split} &\sigma_{\text{inel}} = &79.3 \pm 2.9 \text{ mb (ATLAS)} \\ &\sigma_{\text{inel}} = &71.3 \pm 3.5 \text{ mb (CMS)} \end{split}$$

The largest uncertainty contribution comes from the extrapolation and the luminosity determination.

• The elastic, inelastic and total cross sections at $\sqrt{s} = 8$ TeV have been measured by ATLAS and TOTEM exploiting the optical theorem:

ATLAS (luminosity-dependent):

TOTEM (luminosity-independent):

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- $\begin{array}{lll} \sigma_{tot} = (96.07 \pm 0.92) \text{ mb} & \sigma_{tot} = (101.7 \pm 2.9) \text{ mb} \\ \sigma_{el} = (24.33 \pm 0.39) \text{ mb} & \sigma_{el} = (27.1 \pm 1.4) \text{ mb} \\ \sigma_{inel} = (71.73 \pm 0.71) \text{ mb} & \sigma_{inel} = (74.7 \pm 1.7) \text{ mb} \end{array}$
- TOTEM has excluded a single-exponential shape of $d\sigma_{el}/dt$ with more than 7σ .
- TOTEM also measured σ_{tot} = (102.9 \pm 2.3) mb and ρ = 0.12 \pm 0.03 with the β^* = 1 km dataset.

Backup slides

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Backup - ATLAS test of nuclear slope parametrization

- ATLAS has tried different parametrizations for the nuclear slope.
- The upper limit of the fit range was increased to $|t| = 0.3 \text{ GeV}^2$ in order to increase the sensitivity of additional parameters.
- The quality of the fit is increased due to the higher number of free parameters.

	$\sigma_{ m tot}[{ m mb}]$	Model
Nominal	96.07 ± 0.86	$f_{\rm N}(t) = (\rho + i) \frac{\sigma_{\rm tot}}{\hbar c} e^{-Bt/2}$
Ct^2	96.16 ± 0.80	$f_{\rm N}(t) = (\rho + i) \frac{\sigma_{\rm tot}}{\hbar c} e^{-Bt/2 - Ct^2/2} $
$c\sqrt{-t}$	96.40 ± 0.80	$f_{\rm N}(t) = (\rho + i) \frac{\sigma_{\rm tot}}{\hbar c} e^{-Bt/2 - c/2(\sqrt{4\mu^2 - t - 2\mu})} , \mu = m_{\pi}$
SVM	96.16 ± 0.80	$f_{\rm N}(t) = \rho \frac{\sigma_{\rm tot}}{\hbar c} e^{-B_R t/2} + i \frac{\sigma_{\rm tot}}{\hbar c} e^{-B_I t/2}$
BP	96.81 ± 0.95	$f_{\rm el}(t) = i \left[G^2(t) \sqrt{A} e^{-Bt/2} + e^{i\phi} \sqrt{C} e^{-Dt/2} \right]$
BSW	96.67 ± 0.99	$\operatorname{Re} f_{\mathrm{el}}(t) = c_1(t_1+t)e^{-b_1t/2}$, $\operatorname{Im} f_{\mathrm{el}}(t) = c_2(t_2+t)e^{-b_1t/2}$

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Backup - Results for the nuclear B-slope

- ATLAS measurement: $B = 19.73 \pm 0.24 \text{ GeV}^{-2}$
- TOTEM measurement: $B = 19.9 \pm 0.3 \text{ GeV}^{-2}$
- Pre-LHC expectations was a linear evolution of the B-slope with ln(s)
- LHC measurements of the B-slope favours a second ln²(s) term.



Backup - *t*-reconstruction methods

• The scattering angle is calculated from the proton transverse positions far from IP:

$$\begin{pmatrix} u \\ \theta_u \end{pmatrix} = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} u^* \\ \theta^*_u \end{pmatrix}, \quad u = (x, y).$$

Subtraction method:

$$\theta_u^* = \frac{u_A - u_C}{M_{12,A} + M_{12,C}}, \quad u = x, y$$

• Local angle method method:

$$\theta_x^* = rac{ heta_{x,A} - heta_{x,C}}{M_{22,A} + M_{22,C}}, \quad \theta_y^* ext{ as for subtraction}$$

• Local subtraction method:

$$\theta_{x,S}^* = \frac{M_{11,S}^{241} \cdot x_{237,S} - M_{11,S}^{237} \cdot x_{241,S}}{M_{11,S}^{241} \cdot M_{12,S}^{237} - M_{11,S}^{237} \cdot M_{12,S}^{241}}, \quad S = A, C, \quad \theta_y^* \text{ as for subtraction}$$

• Lattice method:

$$\theta_x^* = M_{12}^{-1} \cdot x + M_{22}^{-1} \cdot \theta_x$$
, θ_y^* as for subtraction

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- The beam optics has direct influence on the *t*-reconstruction through the transport matrix.
- Different *t*-reconstructions gives different results
 ⇒ the initial **design** optics needs modifications.
- Both TOTEM and ATLAS use elastic data to constrain an optics fit including magnet strengths whereby an effective optics in obtained.



$$\frac{y_L}{y_R} = \frac{M_{12}^L}{M_{12}^R} \, .$$



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