Measurements of the elastic, inelastic and total *pp* cross sections with the ATLAS, CMS and TOTEM detectors

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Outline:

- Physics motivation.
- Inelastic cross section at \sqrt{s} = 13 TeV with the ATLAS and CMS detectors.
- Elastic, inelastic and total cross sections at $\sqrt{s} = 8$ TeV with the ATLAS and TOTEM detectors.

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Physics motivation

- **The elastic (σ_{el}), inelastic (σ_{inel}) and total (σ_{tot})** *pp* **cross sections are fundamental** quantities which cannot be calculated with perturbative QCD.
- Regge theory provides a description but data is needed to constrain models. \bullet
- σ_{tot} gives the upper bound on any *pp* process and is seen to rise with collision energy. \bullet
- A substantial fraction of σ_{inel} is diffractive processes. Measurement of σ_{inel} will constrain models also for cosmic-ray shower in the atmosphere.

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Inelastic cross section measurement at $\sqrt{s} =$ 13 TeV with ATLAS and CMS

arXiv:1606.02625 CMS-PAS-FSQ-15-005

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- Strategy: Measure the inelastic cross section in a fiducial region and extrapolate to full phase-space with input from theoretical models.
	- The better detector coverage, the smaller extrapolation uncertainty.
- \bullet The fiducial σ_{inel} is the number of observed events corrected for background, pile-up, efficiencies and luminosity.

ATLAS:

MBTS plastic scintillators at $z = +3.6$ m covering $2.07 < |\eta| < 3.86$.

WLS fibers **Scintillators**

CMS:

- HF calorimeter of iron absorbers and quartz fibers covering $3.0 < |\eta| < 5.2$.
- CASTOR calorimeter of tungsten and quartz covering $-6.6 < \eta < -5.2$.

Inelastic cross section - Tuning models (ATLAS)

- **The inelastic cross section is the sum of the** non-diffractive and the diffractive cross section.
- **•** The ratio $f_D = (\sigma_{SD} + \sigma_{DD})/\sigma_{inel}$ is poorly known and differs between models.
- \bullet The fraction of single-sided events, R_{SS} , is related to f_D and used to tune f_D in the models.

- Using the *f_D*-tuned models, the hit multiplicity in the MBTS for the models are compared to data:
	- The DL (Donnachie-Landshoff) pomeron flux model is best.
	- **The FPOS LHC and OGSJFT-II models** (developed for cosmic-ray showering) are worst.
- CMS checks that models predict correct ratio of cross sections between the HF-only and the HF+CASTOR phase-space regions.

 CMS

- The average of the model extrapolation factors is used to go from fiducial to full phase-space cross section.
- Maximum difference between models is used as systematic uncertainty.

ATLAS:

A precise (and independent) measurement of σ_inel at $\sqrt{s} =$ 7 TeV is used:

$$
\sigma_{inel}=\sigma_{inel}^{fid}+(\sigma_{inel,\ 7\ TeV}^{ALFA}-\sigma_{inel,\ 7\ TeV}^{fid})\times\frac{\sigma_{inel}^{MC}(\xi<10^{-6})}{\sigma_{inel,\ 7\ TeV}^{MC}(\xi<5\times10^{-6})}
$$

The difference between Pythia8 DL and Pythia8 MBR is used as systematic uncertainty.

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Inelastic cross section - Results

CMS:

\n
$$
\sigma_{inel}^{fid,HF} = 65.8 \pm 0.8(\text{exp.}) \pm 1.8(\text{lum.}) \text{ mb}
$$
\n
$$
\sigma_{inel}^{fid,HF+CASTOR} = 66.9 \pm 0.4(\text{exp.}) \pm 2.0(\text{lum.}) \text{ mb}
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\sigma_{inel} = 71.3 \pm 0.5(\text{exp.}) \pm 2.1(\text{lum.}) \pm 2.7(\text{ext.}) \text{ mb}
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Elastic, inelastic and total cross sections measurement at $\sqrt{s} =$ 8 TeV with ATLAS and TOTEM

> Phys. Lett. B (2016) 158 Phys. Rev. Lett. 111, 012001 (2013) Nucl. Phys. B 899 (2015) 527-546 CERN-PH-EP-2015-325

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• From the optical theorem we get:

$$
\sigma_{\text{tot}}^2 = \frac{16\pi(\hbar c)^2}{1+\rho^2} \frac{1}{L} \frac{dN_{\text{el}}}{dt}\Big|_{t=0} \quad \text{with} \quad \rho = \frac{\text{Re}\left[F_{\text{el}}(t)\right]}{\text{Im}\left[F_{\text{el}}(t)\right]} \quad \text{(ATLAS)}
$$
\nequiv.
$$
\sigma_{\text{tot}} = \frac{16\pi(\hbar c)^2}{1+\rho^2} \frac{1}{N_{\text{el}} + N_{\text{inel}}} \frac{dN_{\text{el}}}{dt}\Big|_{t=0} \quad \text{(TOTEM)}.
$$

The four-momemtum transfer *t* is calculated as

$$
t=-(p\times\theta^*)^2.
$$

where the scattering angle θ^* is calculated from the proton trajectories and *p* is the beam momentum.

- Data are taken in runs with low pile-up (μ \lesssim 0.1) and high β^* collision optics since $t_{\sf min}$ $\propto \frac{\rho^2}{\beta^*}$ $\frac{\mu}{\beta^*}$:
	- $\beta^* = 90$ m results from ATLAS (one dataset) and TOTEM (two datasets).
	- $\beta^* = 1$ km results from TOTEM (one dataset).

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- ATLAS and TOTEM use tracking detectors in Roman Pots at *z* ∼ 230 m to approach outgoing beams in vertical direction.
	- ATLAS uses 10×2 orthogonal layers of scintillating fibers giving ≈ 30 *µ*m tracking resolution.
	- TOTEM uses a stack of 10 silicon strip detectors giving ≈ 11 *µ*m tracking resolution.
- In addition, TOTEM has two tracking telescopes:
	- **T1** is a cathode strip chamber at $z = \pm 9$ m covering 3.1 $\leq |\eta| \leq 4.7$.
	- T2 is based on gas electron multiplier chambers at $z = \pm 13.5$ m covering $5.3 < |\eta| < 6.5$.

Elastic analysis - event selection

- **Elastic events are selected when all four** detectors in an arm have a track.
- The tracks are required to fulfill certain correlations between inner-outer stations and between left and right side.
- ATLAS $\beta^* = 90$ m: 3.8 M elastic events.
- TOTEM $\beta^* = 90$ m: 0.65 M elastic events.
- TOTEM $\beta^* = 1$ km: 0.35 M elastic events.

Elastic analysis - Experimental effects

- **Background** comes from beam halo, single diffractive and central diffractive protons.
- Fraction is < 0.12 % estimated from antigolden topology.

- **Detector acceptance** is highly dependent on detector distance to the beam and beam divergence.
- Found from simulation tuned to data.
- *t***-resolution** is influenced by detector resolution and beam divergence.
- **e** Relative *t*-resolution is better than 10 % and corrected for by unfolding.
- **Track reconstruction inefficiency** is data driven.

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Elastic analysis - Theoretical prediction

The differential elastic cross section is a superposition of the strong interaction amplitude *F^N* and the Coulomb amplitude F_C added in quadrature giving

$$
\frac{d\sigma_{el}}{dt} \propto \frac{G^4(t)}{|t|^2} + \sigma_{tot}^2(1+\rho^2) \cdot \exp(-B|t|) - \frac{\sigma_{tot}G^2(t)}{|t|} \left[\sin(\phi(t)) + \rho \cos(\phi(t))\right] \cdot \exp\left(\frac{-B|t|}{2}\right)
$$

 \bullet The corrected differential cross section is fitted with σ_{tot} and *B* as free parameters.

\n- ATLAS fit-range:
$$
0.014 \leq |t| \leq 0.1 \text{ GeV}^2
$$
.
\n- $G(t), \phi(t), \rho$ fixed.
\n

- TOTEM $\beta^* = 90$ m fit-range: $0.01 \leq |t| \leq 0.2$ GeV² (1. dataset) $0.02 \leq |t| \leq 0.2$ GeV² (2. dataset)
	- TOTEM neglects the Coulomb and interference terms.
- TOTEM $\beta^* = 1$ km: $6\cdot10^{-4} \leq |t| \leq$ 0.2 GeV².

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- TOTEM measures simultaneously the elastic and inelastic rate and is hence independent of luminosity.
- The inelastic rate is determined with the T1 (3.1 \leq |η| \leq 4.7) and T2 (5.3 \leq |η| \leq 6.5) telescopes, detecting about 95% of the inelastic rate.
- The strategy is similar to the CMS and ATLAS inelastic cross section measurement.
- Events are triggered by the T2 telescope and corrected for experimental effects.
- Uncertainty from extrapolation of fiducial region is dominating ("Low mass diffraction").

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Elastic analysis - Results

 \bullet σ_{el} is the integral of the nuclear part: σ 2 tot 1+ρ 2 2

$$
\sigma_{\text{el}} = \tfrac{\sigma_{\text{tot}}^2}{B} \tfrac{1+\rho^2}{16\pi(\hbar c)^2}.
$$

$$
\bullet\ \sigma_{inel}=\sigma_{tot}-\sigma_{el}.
$$

- The total cross section is still rising with energy.
- **•** The difference between ATLAS and TOTEM $\beta^* = 90$ m corresponds to 1.9σ.
- Using $\beta^* = 1$ km data, TOTEM also measured

$$
\rho=0.12\pm0.03
$$

Model extrapolation from lower energies:

$$
\rho=0.140\pm0.007
$$

Hints of a slight deviation from an exponential fall-off of the elastic nuclear amplitude was reported at ISR and SppS, but not at the Tevatron.

• TOTEM used
$$
\frac{d\sigma}{dt}(t) = \frac{d\sigma}{dt}\Big|_{t=0} \exp\left(\sum_{i=0}^{N_b} b_i t^i\right)
$$

to exclude a purely exponential form with
more than 7 σ !

- This approach was impossible for ATLAS as the *t*-dependent systematic uncertainties are too large.
	- **Average beam energy** uncertainty dominates.
	- **ATLAS** uses a larger uncertainty than TOTEM.
	- An official value will be available later this year.

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 \bullet ATLAS has tested different *B*-parametrizations giving an RMS of 0.28 mb on σ_{tot} .

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Summary

The inelastic cross section at \sqrt{s} = 13 TeV was measured by a MC extrapolation from the fiducial region:

> σ_{inel} =79.3 \pm 2.9 mb (ATLAS) $\sigma_{\text{inel}} = 71.3 \pm 3.5$ mb (CMS)

The largest uncertainty contribution comes from the extrapolation and the luminosity determination.

The elastic, inelastic and total cross sections at $\sqrt{s} =$ 8 TeV have been measured by ATLAS and TOTEM exploiting the optical theorem:

ATLAS (luminosity-dependent):

TOTEM (luminosity-independent):

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 $\sigma_{\text{tot}} = (96.07 \pm 0.92)$ mb $\sigma_{\text{el}} = (24.33 \pm 0.39)$ mb $\sigma_{inel} = (71.73 \pm 0.71)$ mb $\sigma_{\text{tot}} = (101.7 \pm 2.9) \text{ mb}$ $\sigma_{el} = (27.1 \pm 1.4)$ mb $\sigma_{\text{inel}} = (74.7 \pm 1.7) \text{ mb}$

- **TOTEM** has excluded a single-exponential shape of $d\sigma_{el}/dt$ with more than 7 σ .
- TOTEM also measured $\sigma_{\text{tot}} = (102.9 \pm 2.3)$ mb and $\rho = 0.12 \pm 0.03$ with the $\beta^* = 1$ km dataset.

Backup slides

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Backup - ATLAS test of nuclear slope parametrization

- ATLAS has tried different parametrizations for the nuclear slope.
- The upper limit of the fit range was increased to $|t|=$ 0.3 GeV² in order to increase the sensitivity of additional parameters.
- The quality of the fit is increased due to the higher number of free parameters.

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Backup - Results for the nuclear B-slope

- \bullet ATLAS measurement: *B* = 19.73 + 0.24 GeV⁻²
- TOTEM measurement: $B = 19.9 + 0.3$ GeV⁻² \bullet
- Pre-LHC expectations was a linear evolution of the *B*-slope with ln(*s*)
- LHC measurements of the *B*-slope favours a second $ln^2(s)$ term.

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Backup - *t*-reconstruction methods

The scattering angle is calculated from the proton transverse positions far from IP:

$$
\begin{pmatrix} u \\ \theta_u \end{pmatrix} = \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} u^* \\ \theta_u^* \end{pmatrix} , \quad u = (x, y) .
$$

• Subtraction method:

$$
\theta_u^* = \frac{u_A - u_C}{M_{12,A} + M_{12,C}}, \quad u = x, y
$$

• Local angle method method:

$$
\theta_x^* = \frac{\theta_{x,A} - \theta_{x,C}}{M_{22,A} + M_{22,C}} \,, \quad \theta_y^* \text{ as for subtraction}
$$

• Local subtraction method:

$$
\theta_{x,S}^* = \frac{M_{11,S}^{241} \cdot x_{237,S} - M_{11,S}^{237} \cdot x_{241,S}}{M_{11,S}^{241} \cdot M_{12,S}^{237} - M_{11,S}^{237} \cdot M_{12,S}^{241}} \ , \quad S = A, C \ , \quad \theta_{y}^* \text{ as for subtraction}
$$

a Lattice method:

$$
\theta_x^* = M_{12}^{-1} \cdot x + M_{22}^{-1} \cdot \theta_x , \quad \theta_y^* \text{ as for subtraction}
$$

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- The beam optics has direct influence on the *t*-reconstruction through the transport matrix. \bullet
- \bullet Different *t*-reconstructions gives different results ⇒ the initial **design** optics needs modifications.
- Both TOTEM and ATLAS use elastic data to constrain an optics fit including magnet strengths whereby an **effective** optics in obtained.

$$
\frac{y_L}{y_R} = \frac{M_{12}^L}{M_{12}^R} \; .
$$

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