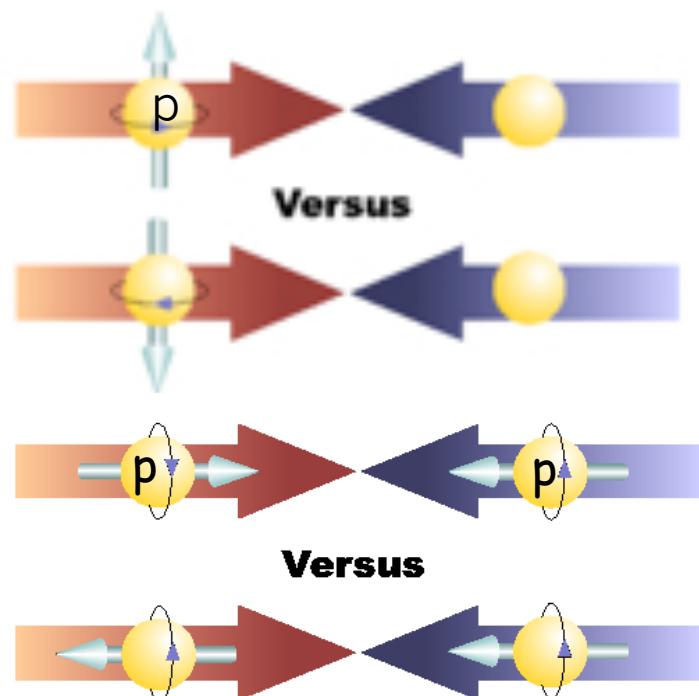


Gluon polarization measurements and the possible role of diffractive process in the transverse single spin asymmetry measurements in RHIC-PHENIX

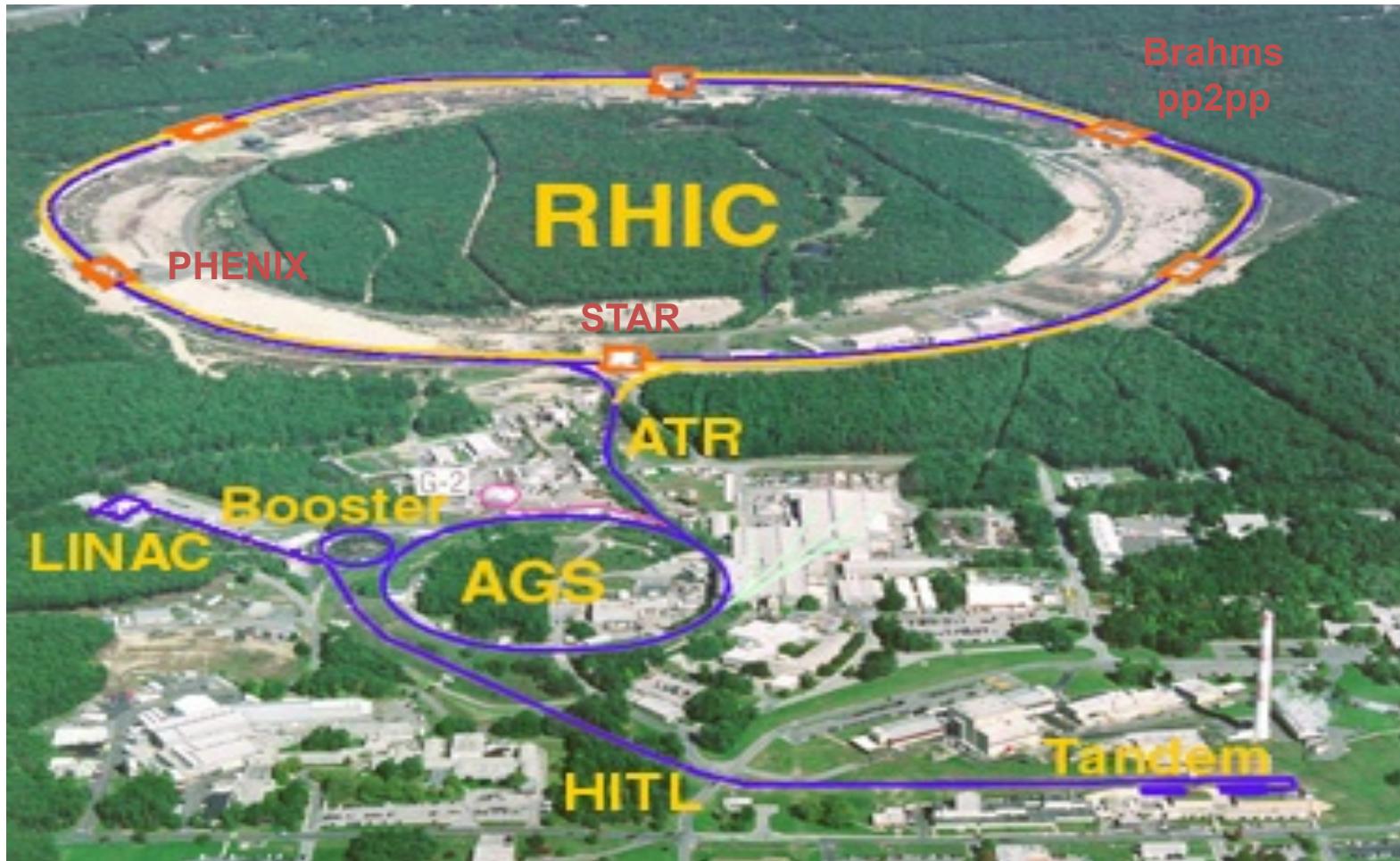
Itaru NAKAGAWA
on behalf of RHIC Spin
Collaboration
RIKEN/RBRC

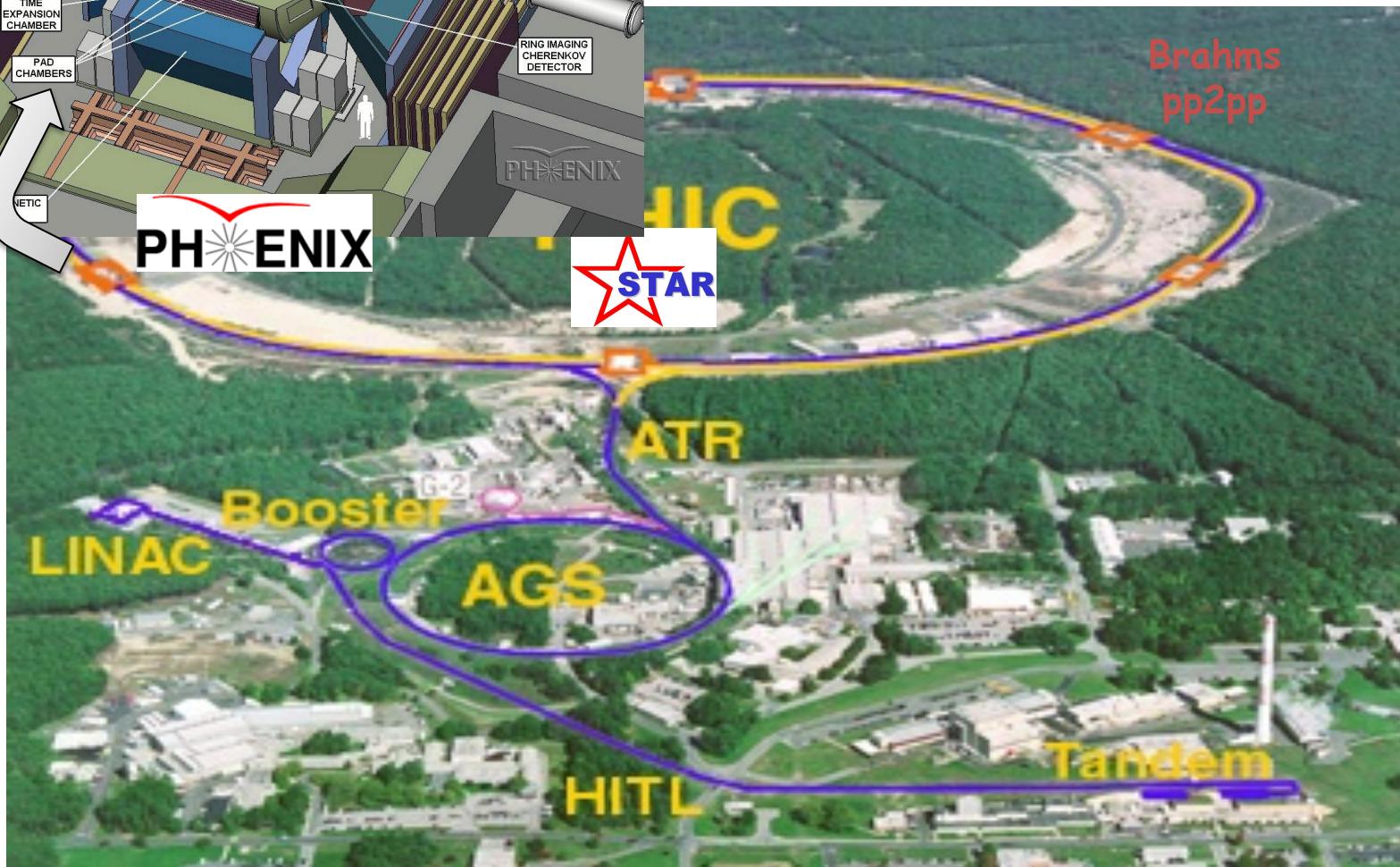
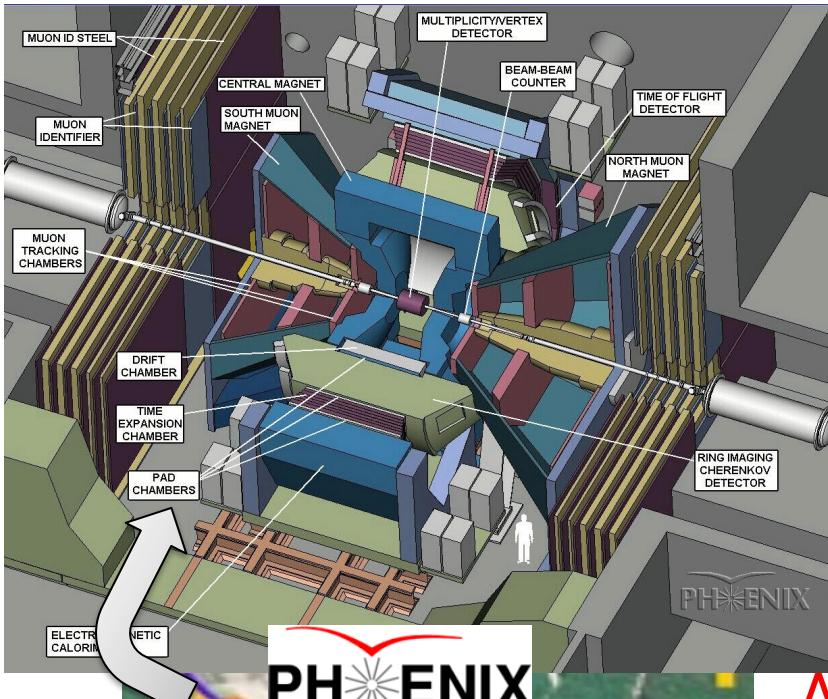
Outline

- RHIC
 - PHENIX @ Polarized Proton @ RHIC
- Transverse
 - very forward neutron
- Longitudinal
 - Gluon Spin

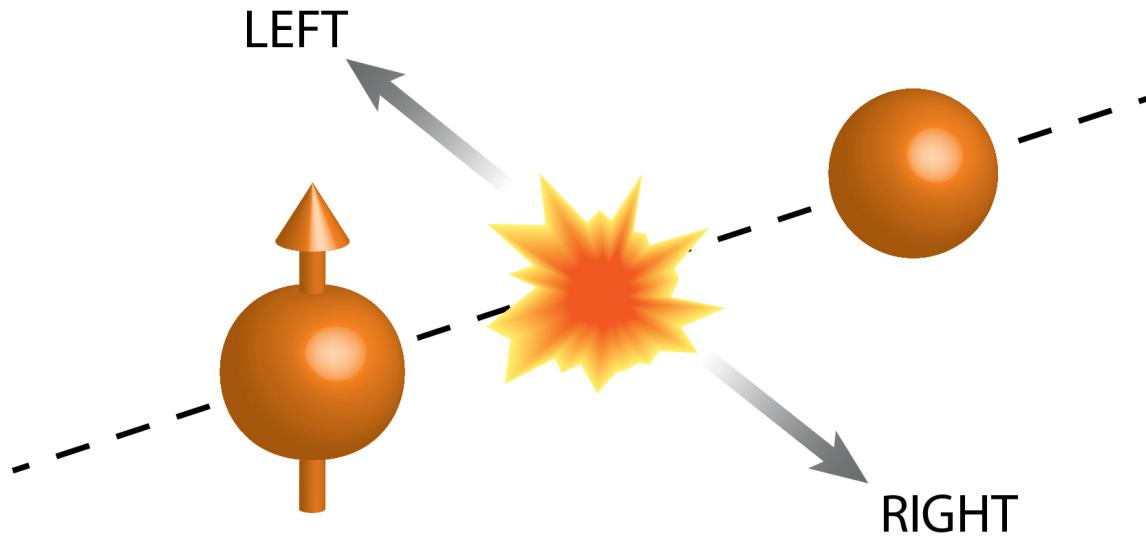


The Relativistic Heavy Ion Collider accelerator complex at Brookhaven National Laboratory





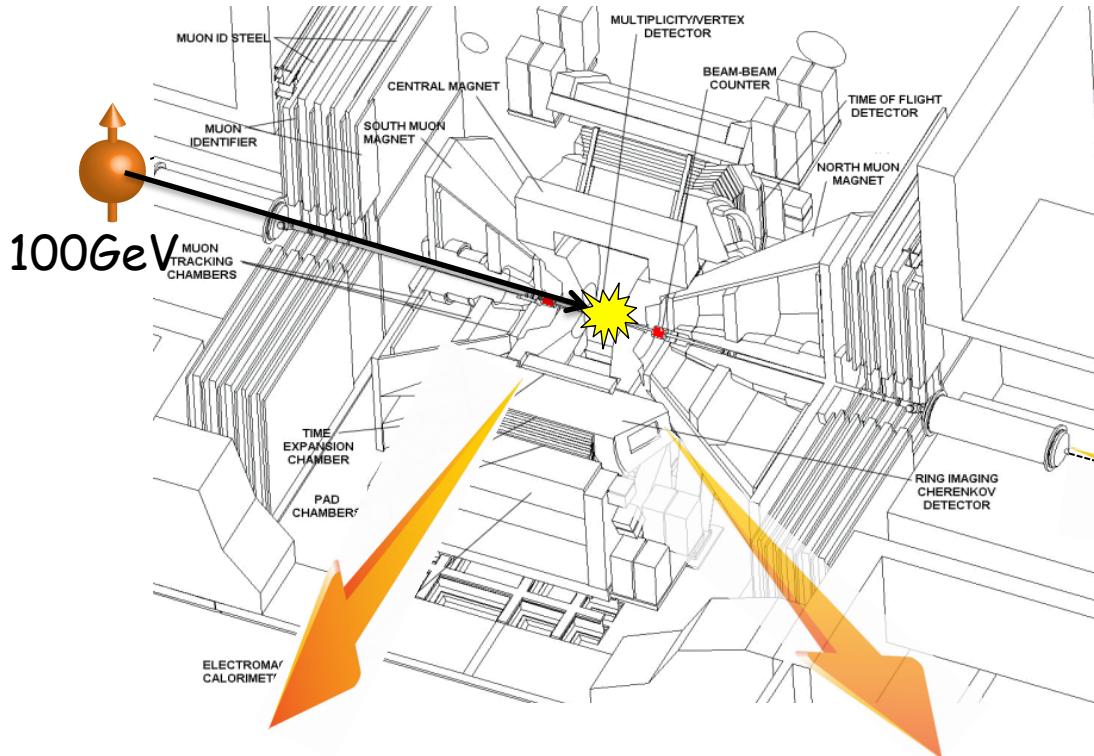
Collision Energy : $62 \sim 510$ GeV
Polarization : 50 - 60 %
Luminosity : $10 \text{ pb}^{-1}/\text{week}$



Forward π^0 and very forward neutron

TRANSVERSE SINGLE SPIN ASYMMETRY

Transverse Asymmetry Observables



Hard Process
(partonic)

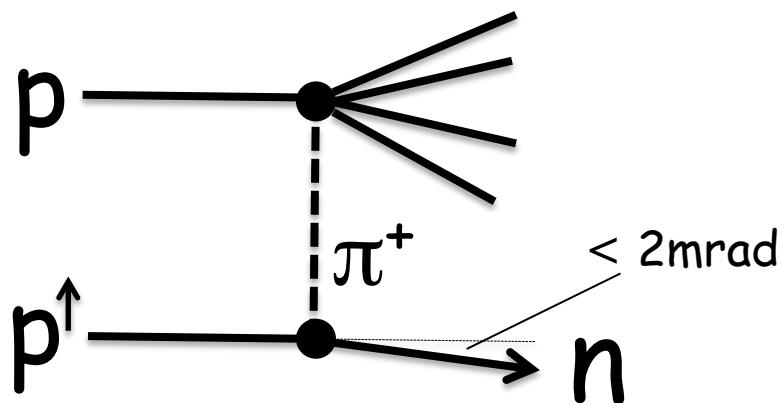
Soft Process
(mesonic)

very forward n
 $(6 < |\eta| < 8)$

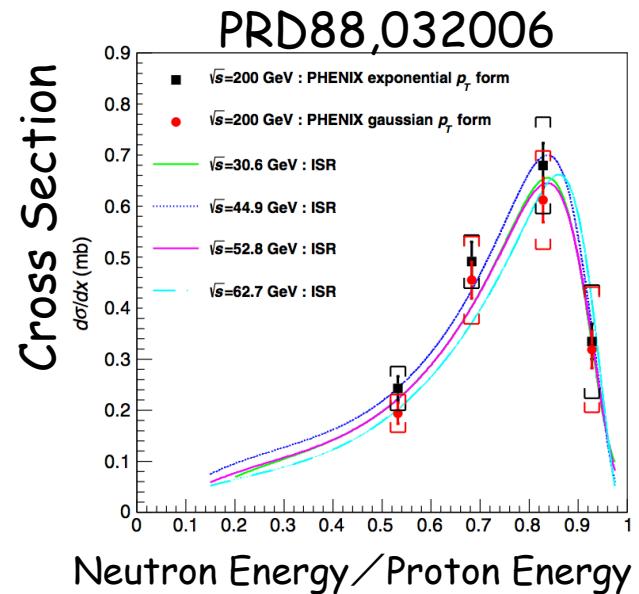
ZDC

Production Mechanism of Forward Neutron

Cross Section



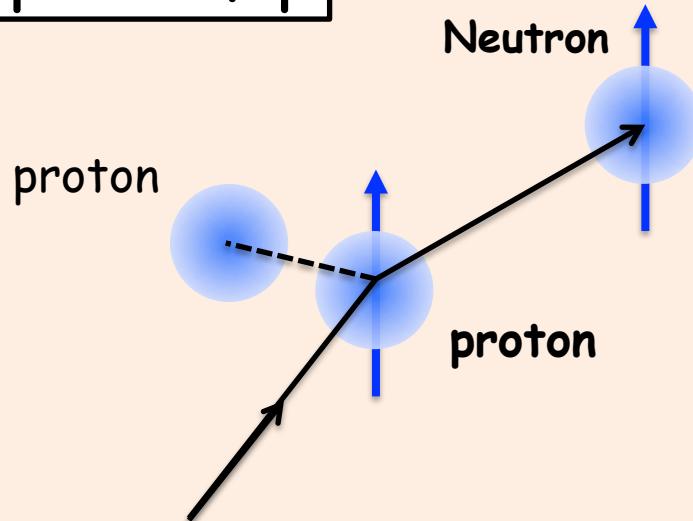
Momentum Transfer $\sim 100\text{MeV}$



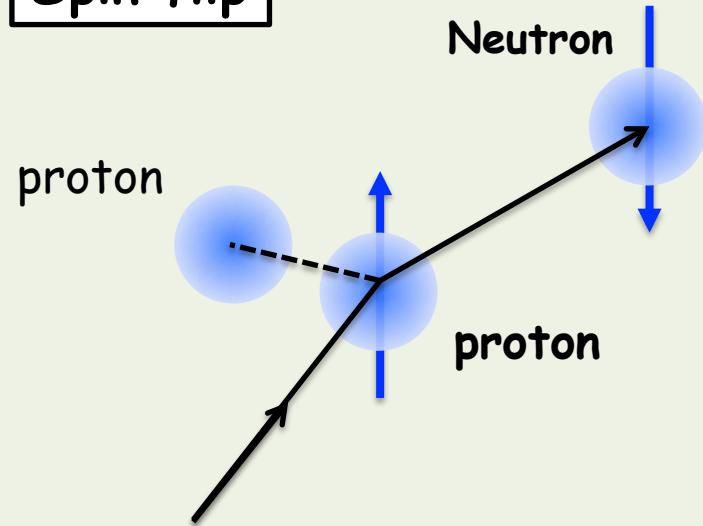
Well Explained by
One-Pion Exchange

$p^\uparrow + p$ Forward Neutron A_N

Spin non-flip



Spin flip



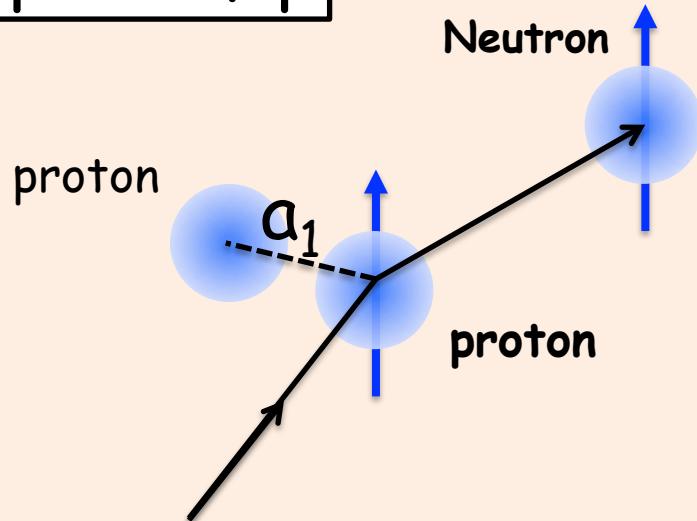
$$A_N \approx \frac{\left(\phi_{non-flip}^* \phi_{flip} \sin \delta \right)}{\left| \phi_{non-flip} \right|^2 + \left| \phi_{flip} \right|^2}$$

δ : phase shift

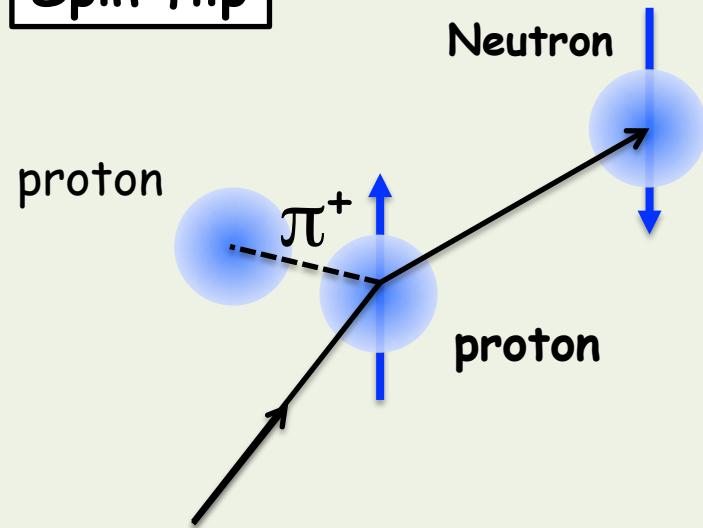
Unpolarized Cross Section

$p^\uparrow + p$ Forward Neutron A_N

Spin non-flip



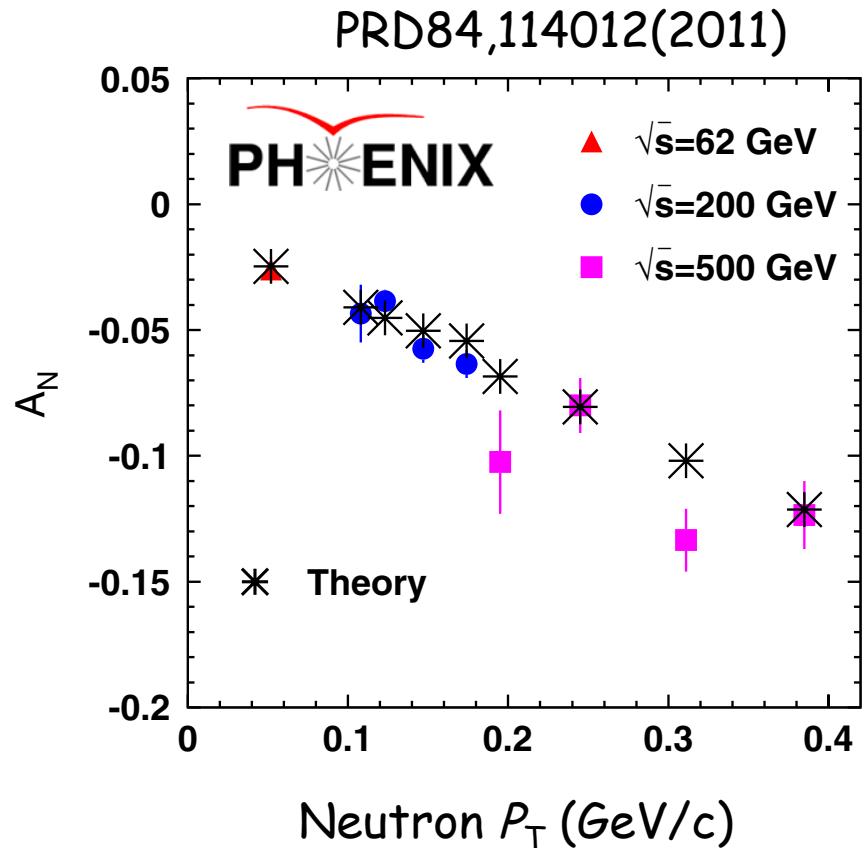
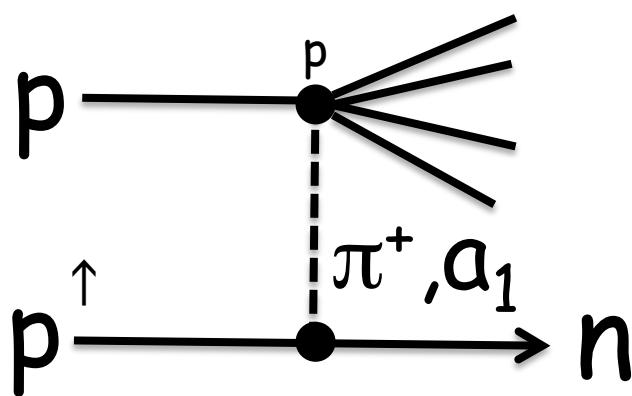
Spin flip



$$A_N \approx \frac{\left(\phi_{non\text{-}flip}^* \phi_{flip} \sin \delta \right)}{\left| \phi_{non\text{-}flip} \right|^2 + \left| \phi_{flip} \right|^2}$$

δ : phase shift

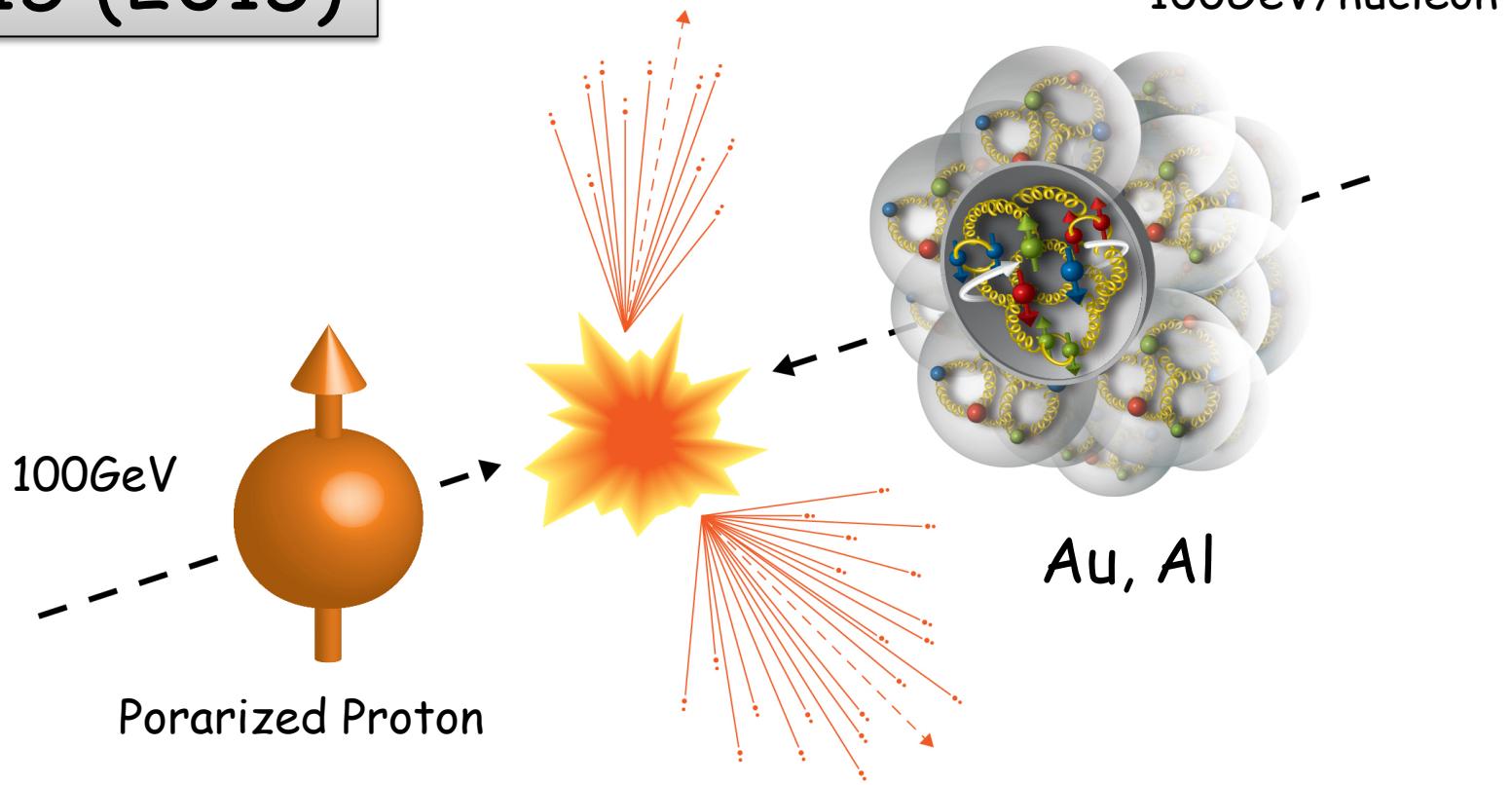
$p^{\uparrow} + p$ Forward Neutron A_N



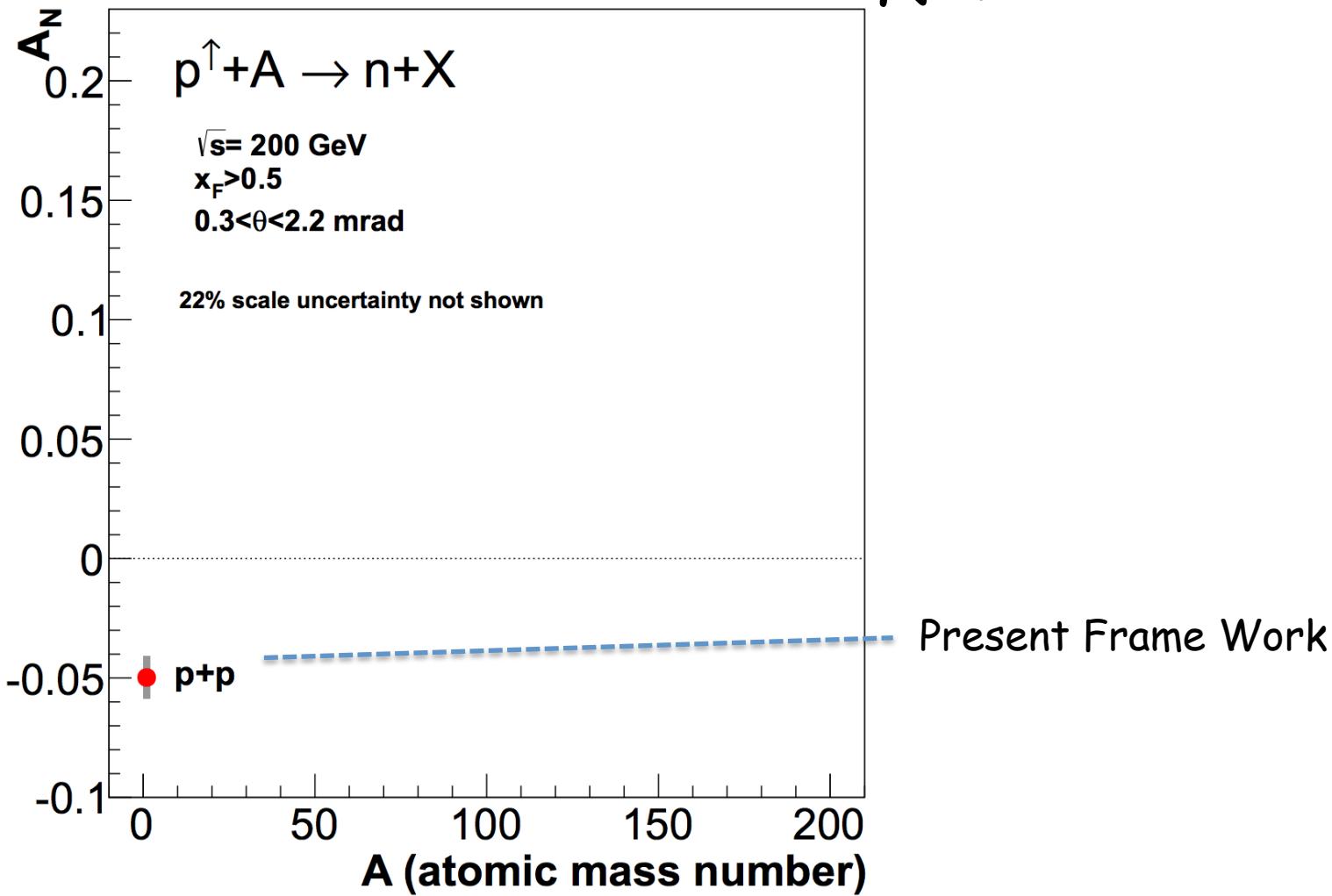
Data are well reproduced by the interference between π and a_1 Reggeon

The First Time Ever High Energy $p^\uparrow + A$ Collisions

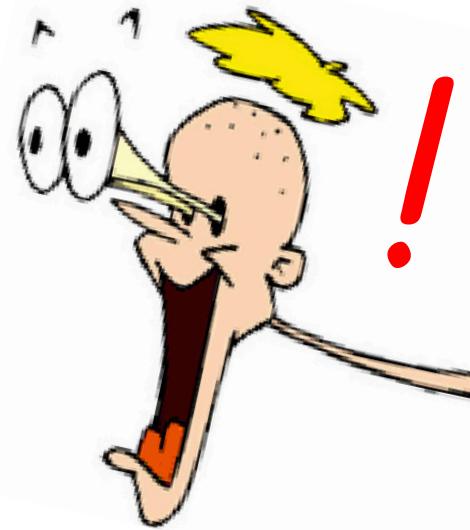
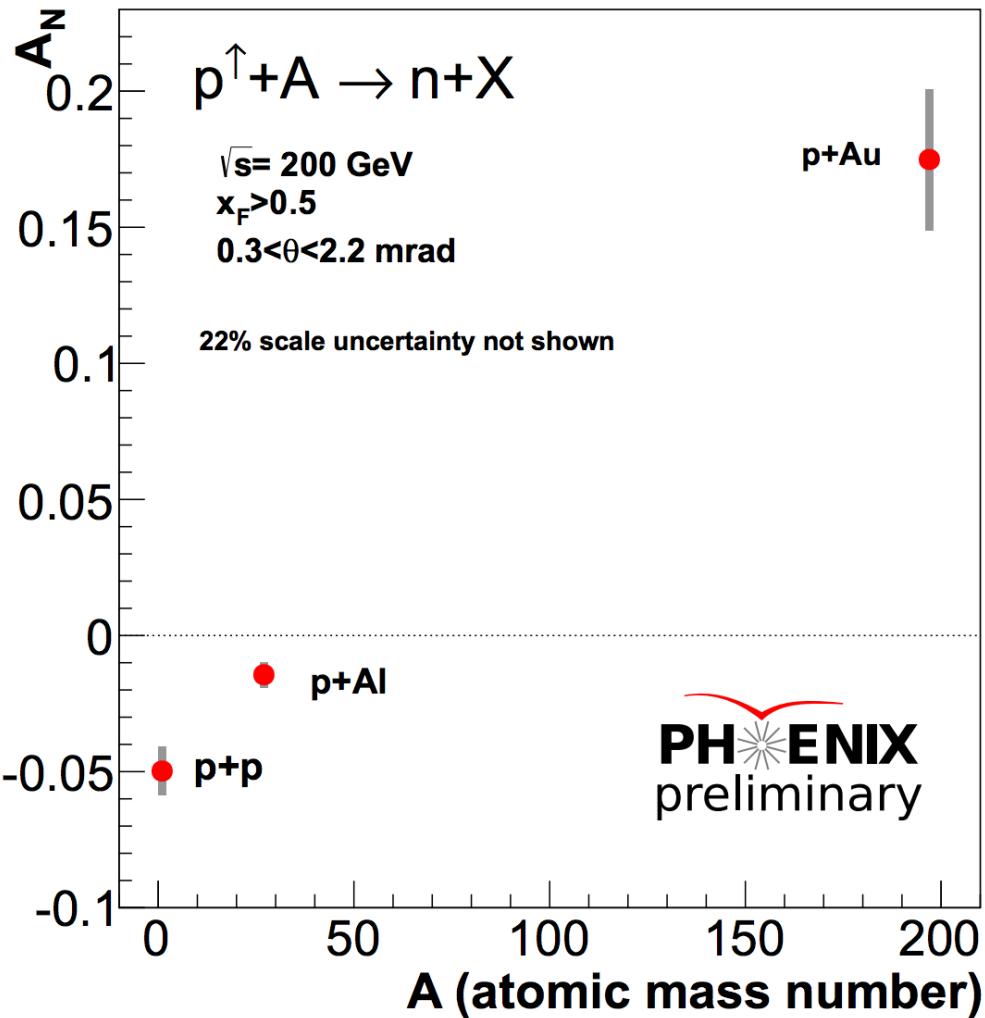
Run15 (2015)



A -Dependent A_N (inclusive)



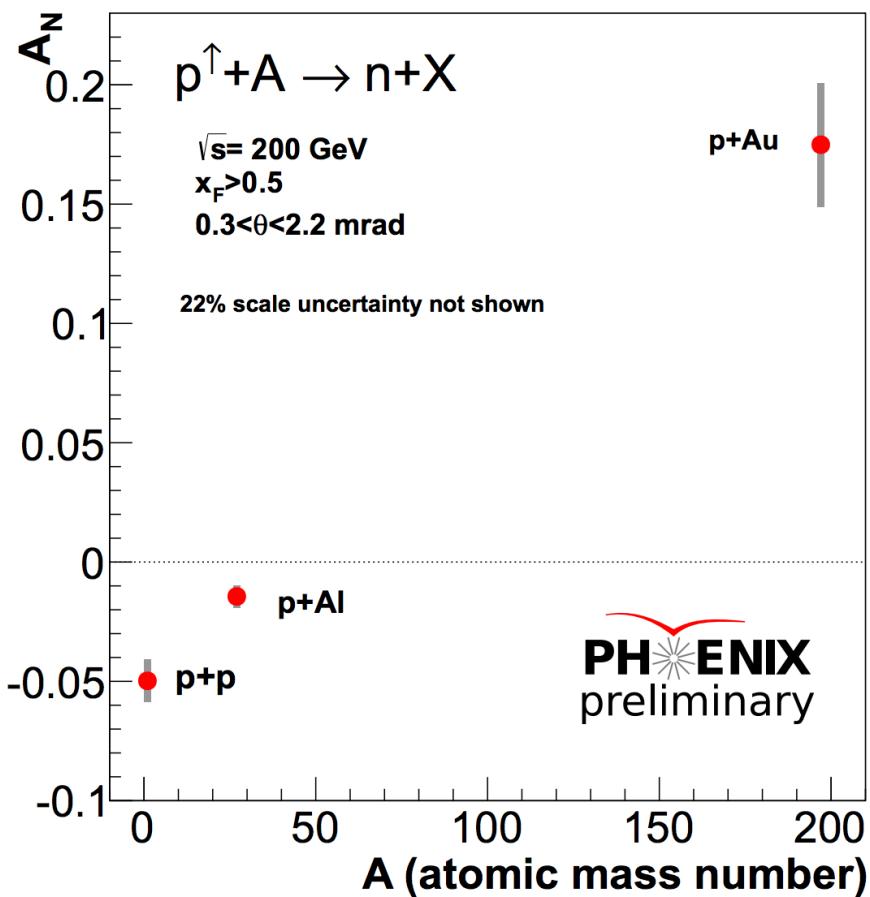
A -Dependent A_N (inclusive)



Absolutely Unexpected!

Analysis by
Minjung Kim (SNU/RIKEN)

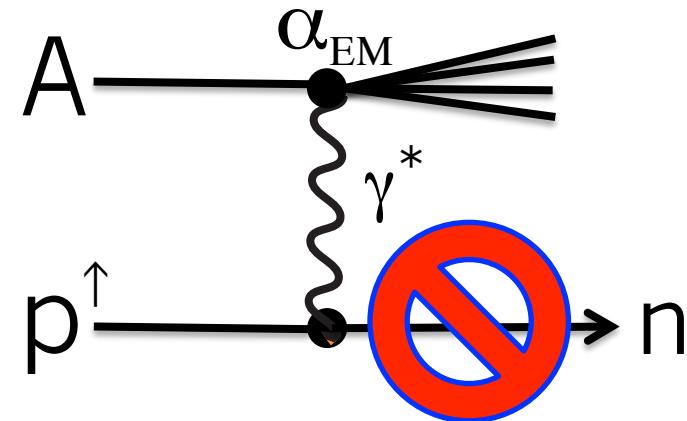
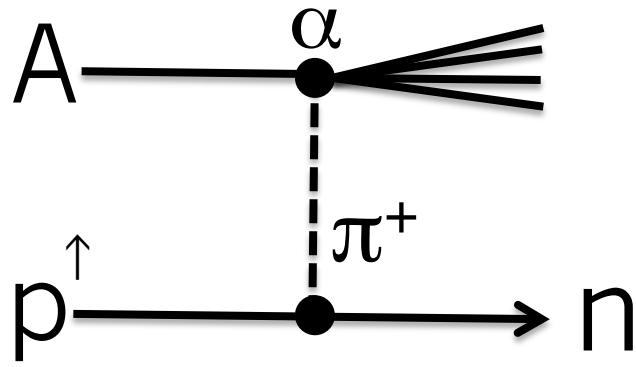
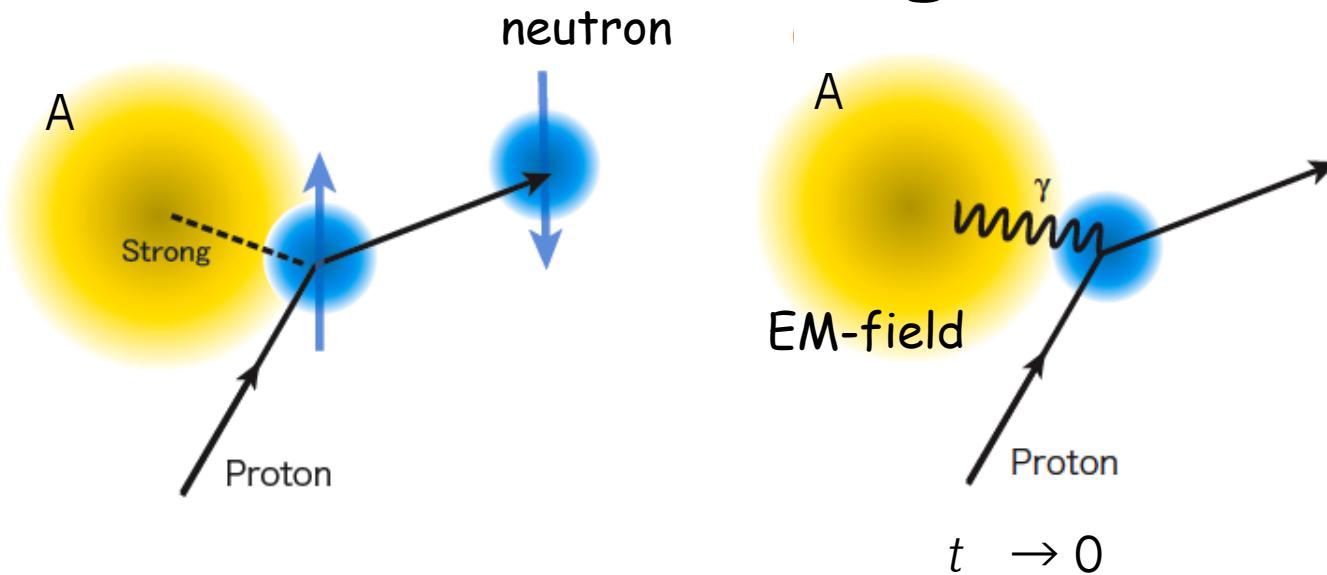
What is going on ?



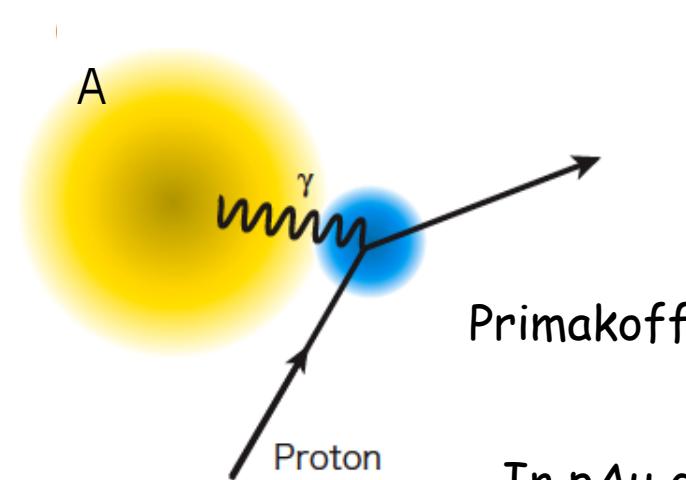
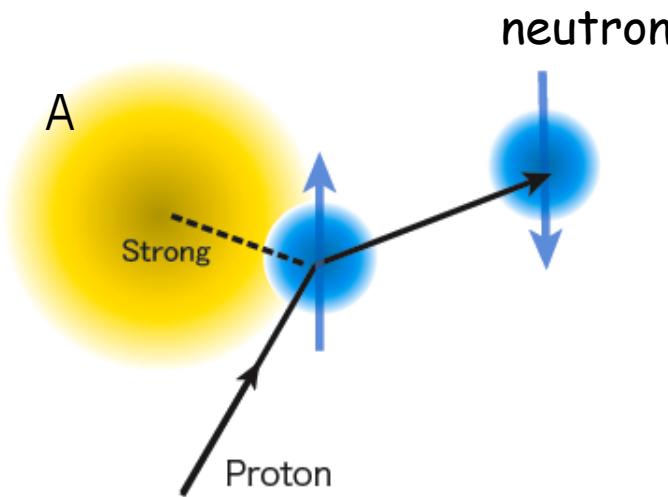
	# of proton	# of neutron
p	1	0
Al	13	14
Au	79	118

- Isospin Symmetry
- Surface Structure of Nucleus
- QED Process
- Gluon Saturation
- Else...

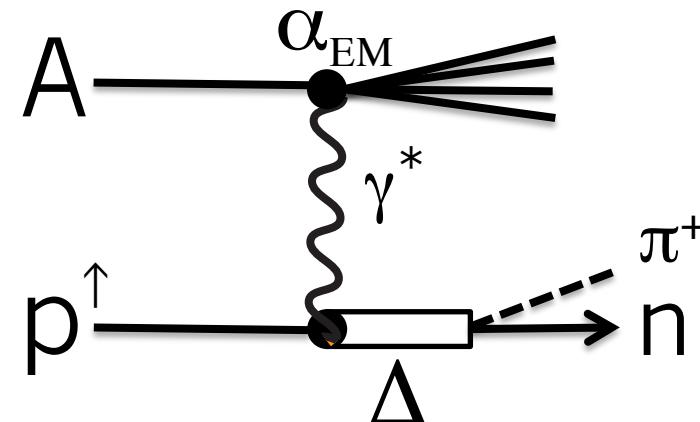
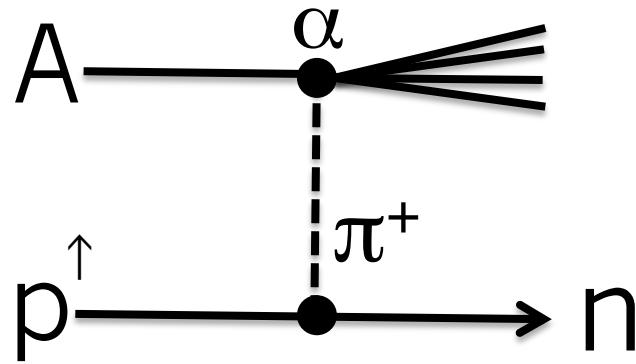
Ultra Peripheral Effect (UPC) Electro-Magnetic



Ultra Peripheral Effect (UPC) Electro-Magnetic



$$t \rightarrow 0 \quad \text{In pAu case} \quad Z \cdot \alpha_{EM} \sim \frac{79}{137}$$



Full Description of A_N

$$A_N \propto 2 \operatorname{Im} \left\{ \phi_{non-flip}^* \phi_{flip} \right\}$$

$$\phi_{flip} = \phi_{flip}^{had} + \phi_{flip}^{EM}$$

$$\phi_{non-flip} = \phi_{non-flip}^{had} + \phi_{non-flip}^{EM}$$

$d_{1 \sim 4}$: relative phase of amplitudes

$$A_N \propto 2 \operatorname{Im} \left(\phi_{non-flip}^{had} + \phi_{non-flip}^{EM} \right) \left(\phi_{flip}^{had*} + \phi_{flip}^{EM*} \right)$$

$$= 2 \left(\phi_{non-flip}^{had} \phi_{flip}^{had} \sin \delta_1 + \phi_{non-flip}^{EM} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had} \phi_{flip}^{EM} \sin \delta_3 + \phi_{non-flip}^{EM} \phi_{flip}^{EM} \sin \delta_4 \right)$$

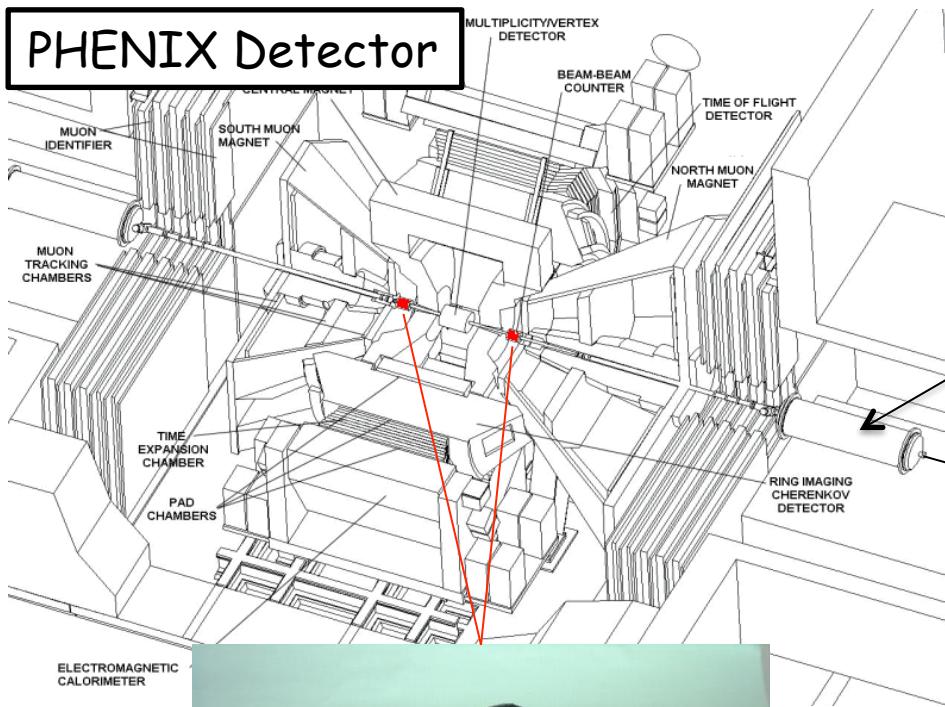


$\rightarrow 0$ (For pp)

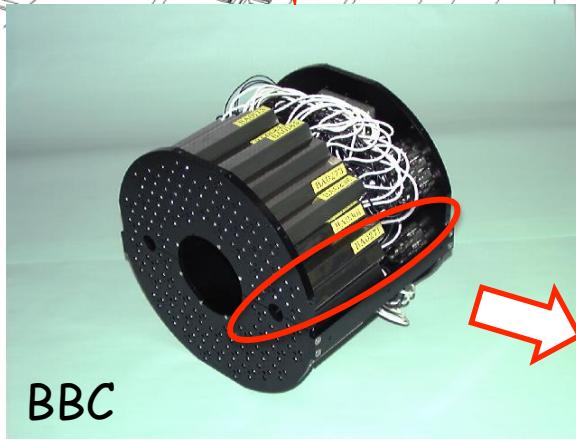


Beam-Beam Counter

PHENIX Detector

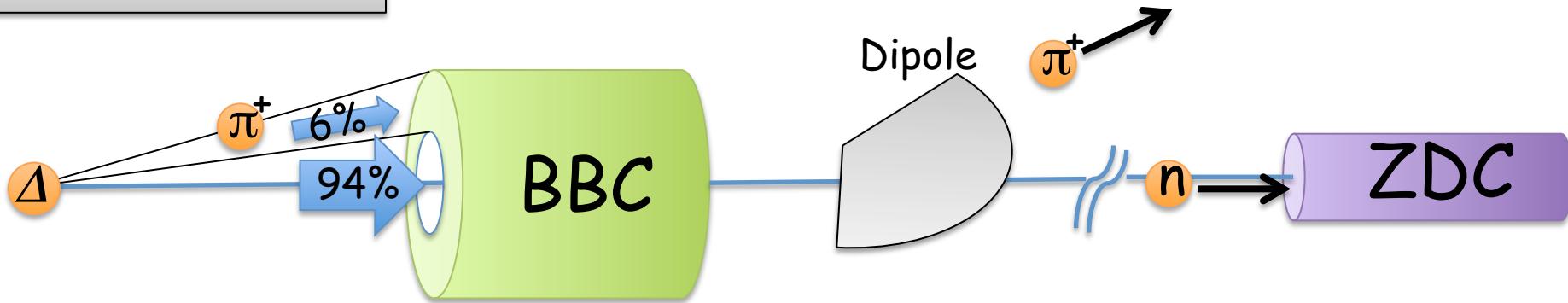


Dipole Bending Magnet



Can we identify UPC events?

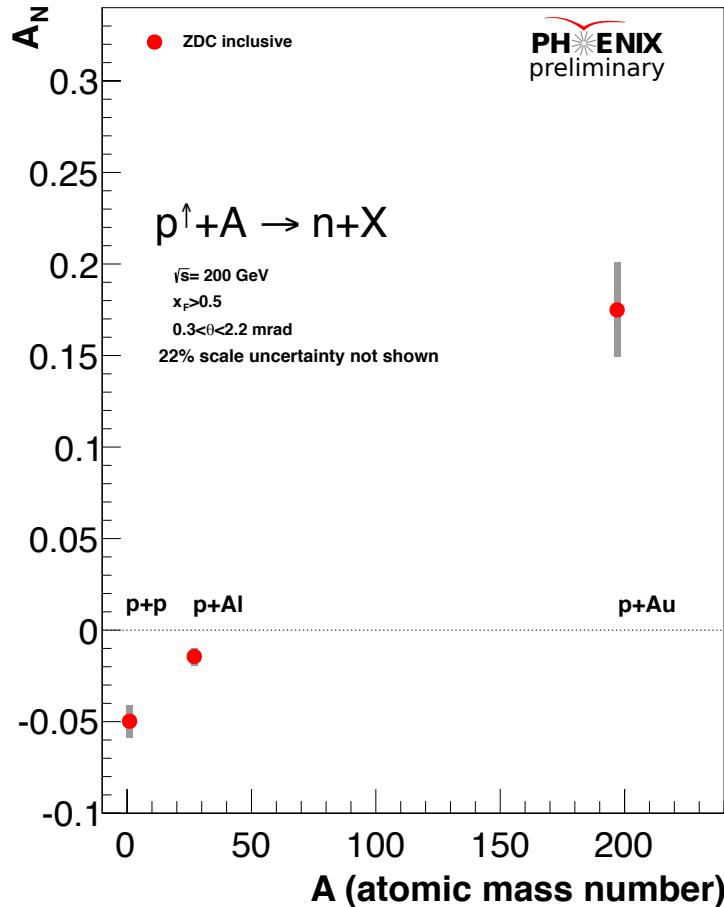
Semi-Inclusive



Primakoff MC : SOPHIA
G. Mitsuka, Eur. Phys. J.C. (2015) 75:614

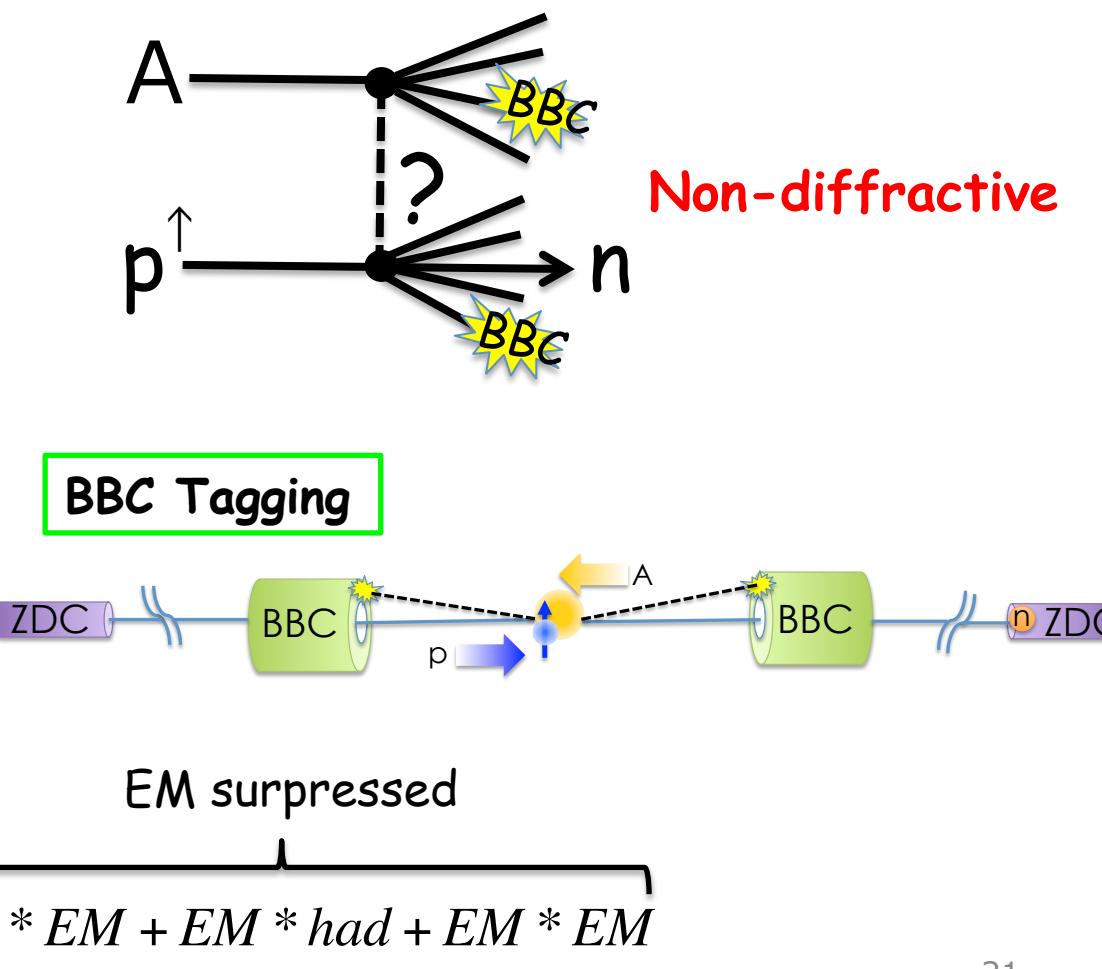
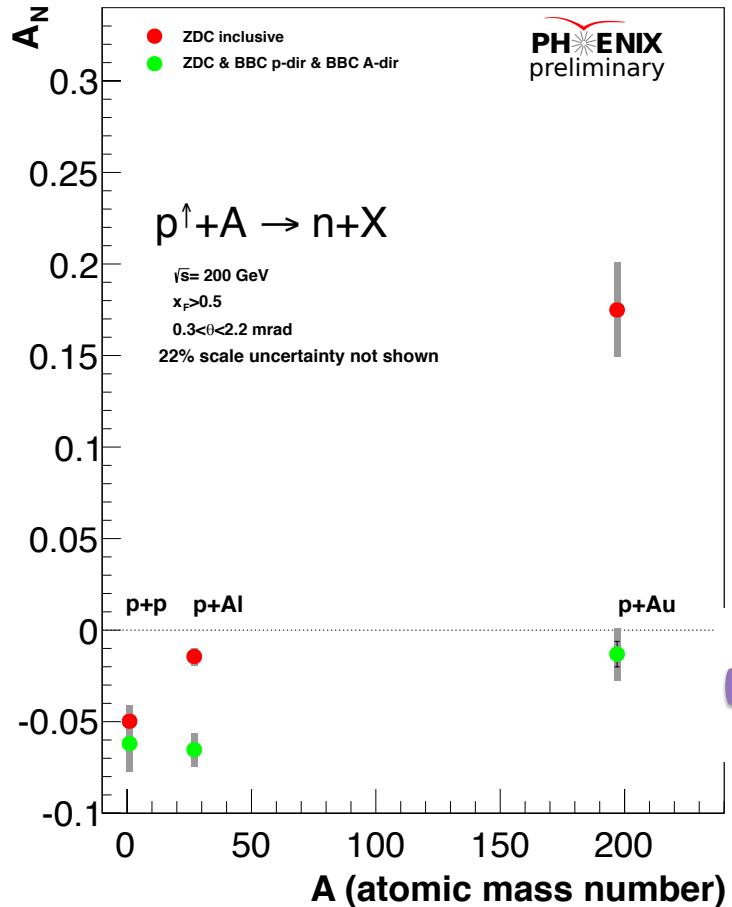
BBC	tag	veto
Primakoff	Small	Large

BBC Tagging and Vetoing

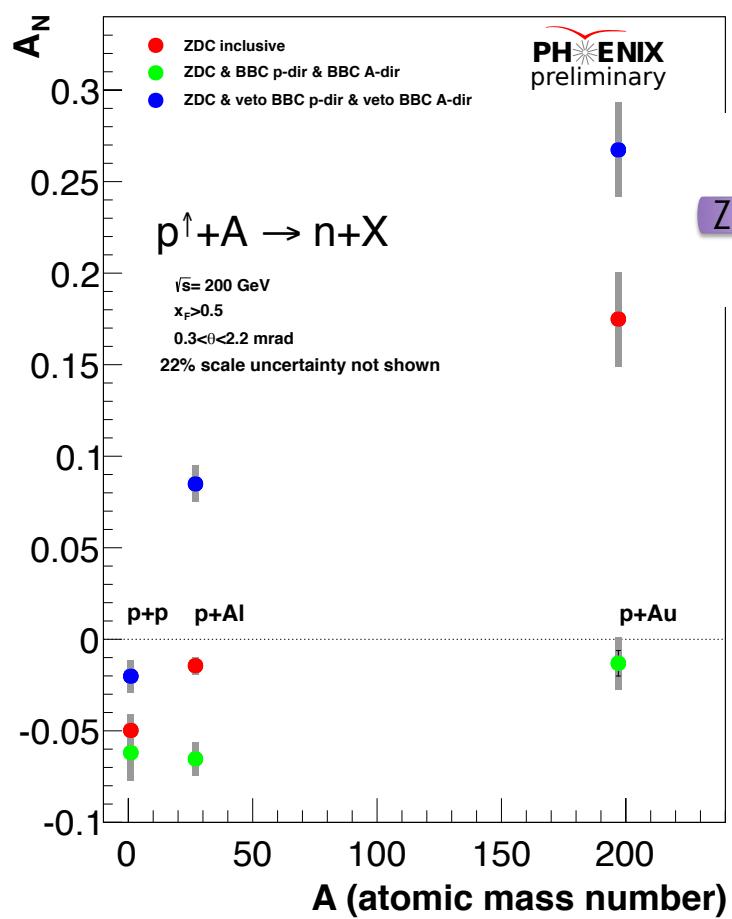


$$A_N \sim [had * had + had * EM + EM * had + EM * EM]$$

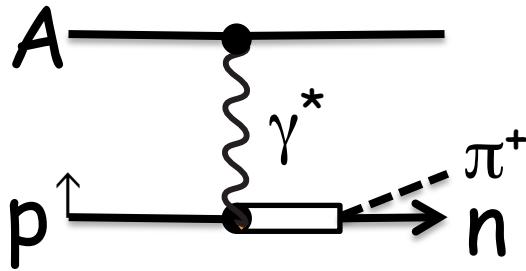
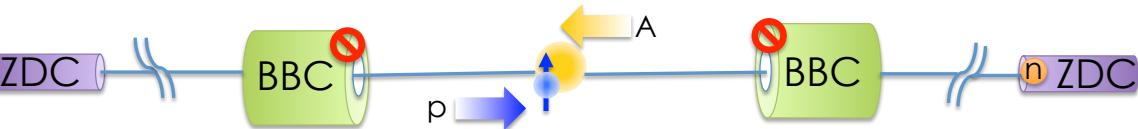
BBC Tagging and Vetoing



BBC Tagging and Vetoing



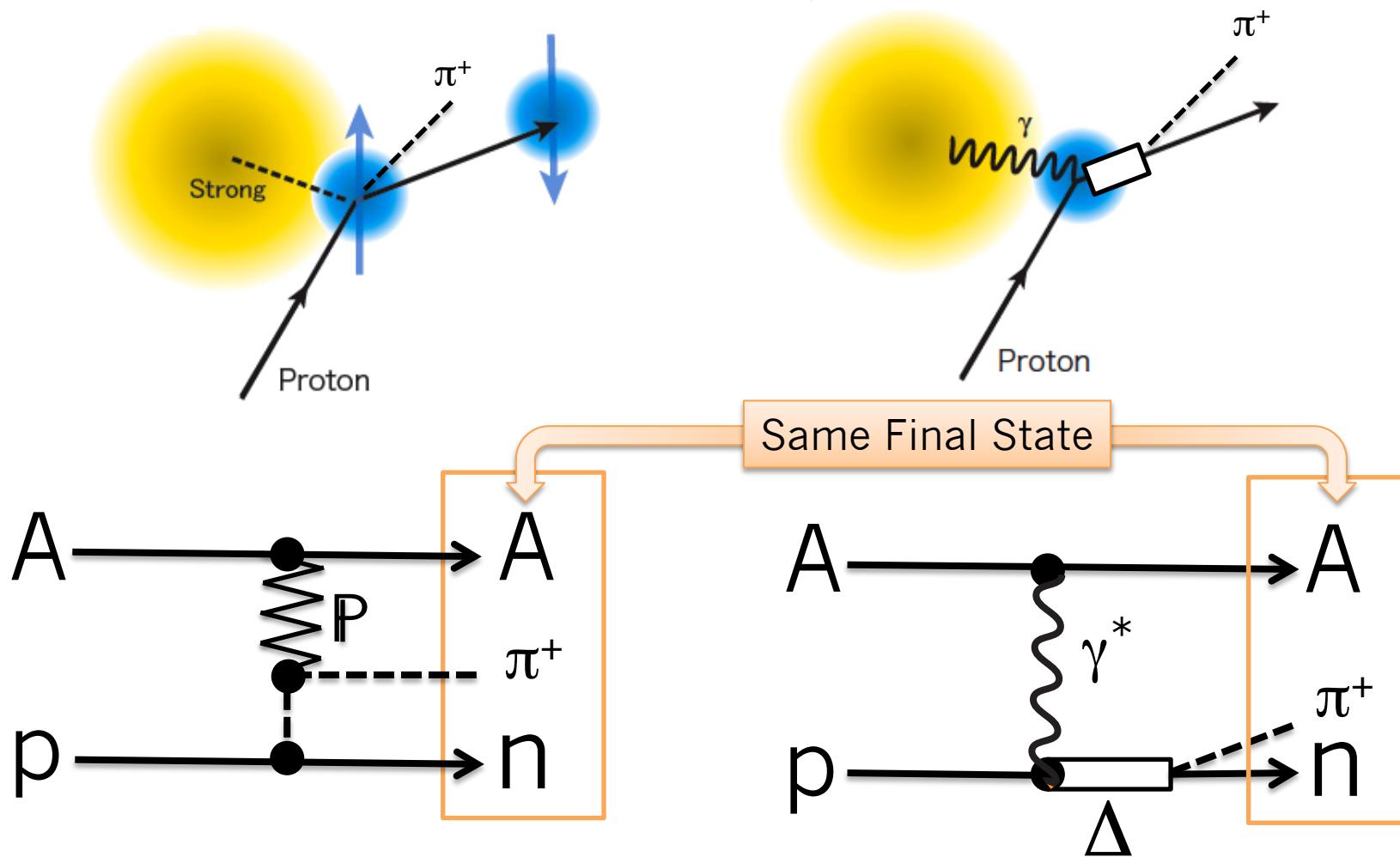
BBC Vetoing



EM enhanced

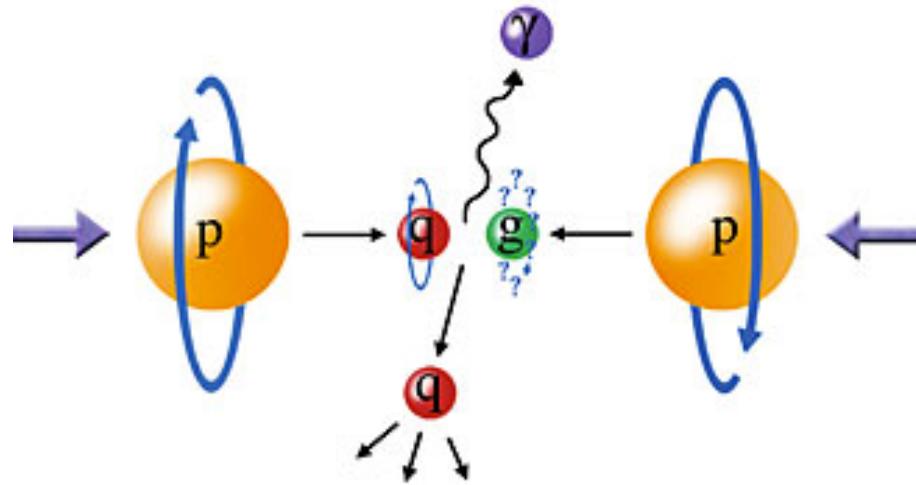
$$A_N \sim \text{had} * \text{had} + \text{had} * \text{EM} + \text{EM} * \text{had} + \text{EM} * \text{EM}$$

Coulomb-Nuclear Interference in Forward Neutron Production



$$A_N \sim \text{had} * \text{had} + \boxed{\text{had} * \text{EM} + \text{EM} * \text{had} + \text{EM} * \text{EM}}$$

Diffractive Process Required?

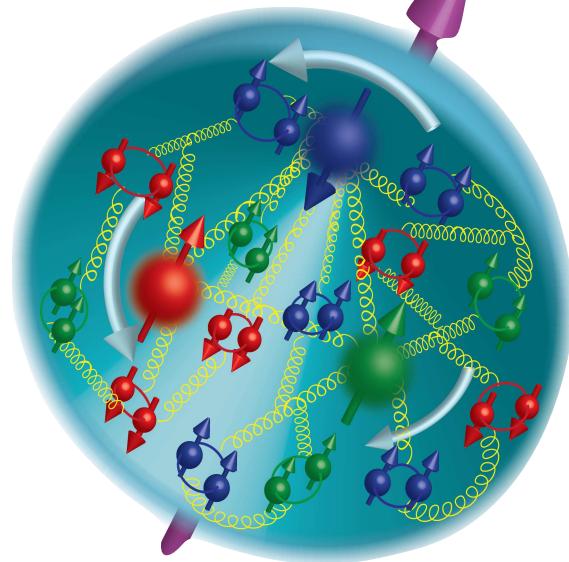


$\Delta G, \Delta \bar{q}$ Polarization

LONGITUDINAL PROGRAM

Proton Spin Decomposition

$$\text{Proton spin} = \frac{1}{2}$$



Longitudinal Spin Sum Rule

$$S_z = \frac{1}{2} \Delta\Sigma + \Delta G + L_z$$

↓ ↓ ↓

~30% ? ??

Quark Spin

$$\Delta\Sigma = \int (\Delta u + \Delta d + \Delta s + \boxed{\Delta \bar{u} + \Delta \bar{d} + \Delta \bar{s}}) dx$$

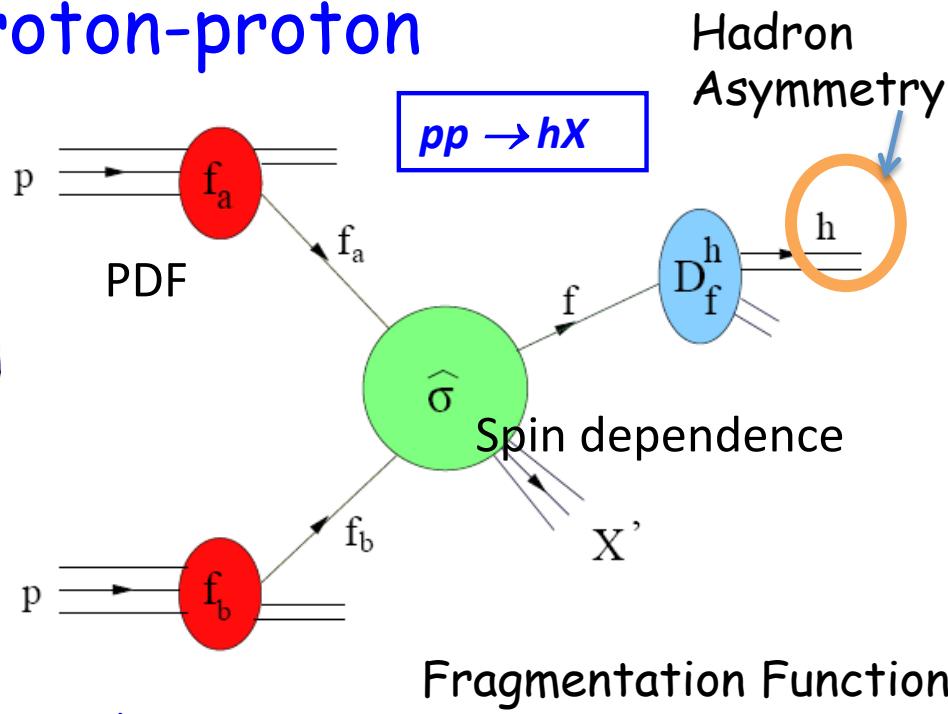
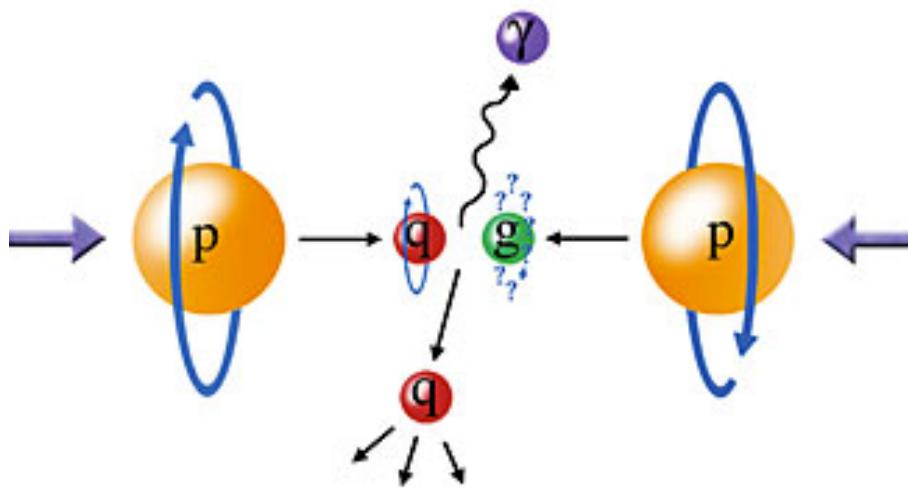
Sea quark

Gluon Spin

$$\Delta G(x) = \boxed{\int \Delta g(x) dx}$$

Longitudinal Spin Structure

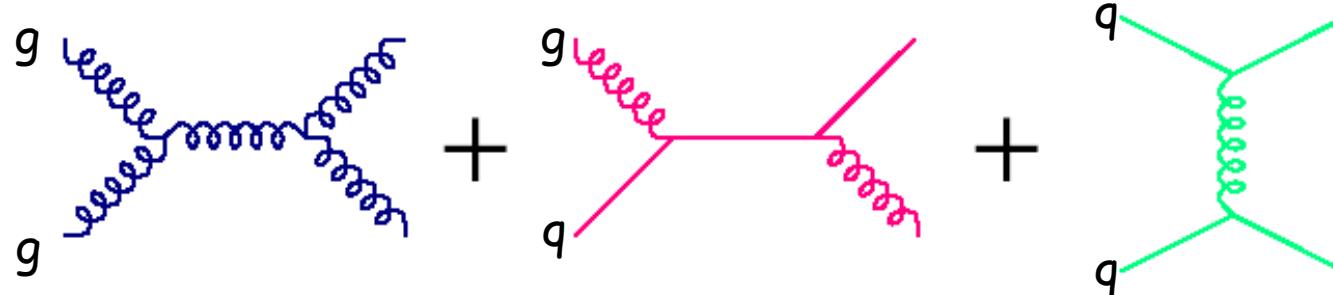
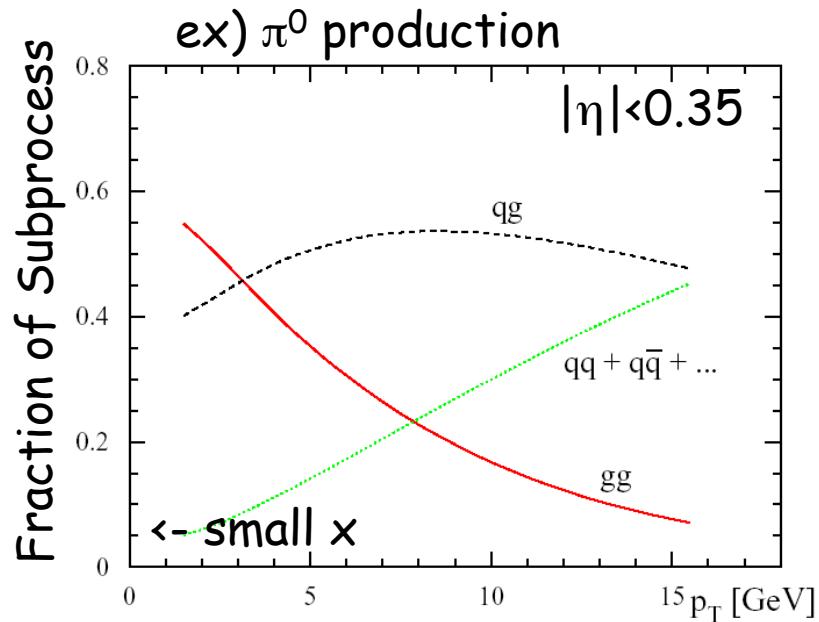
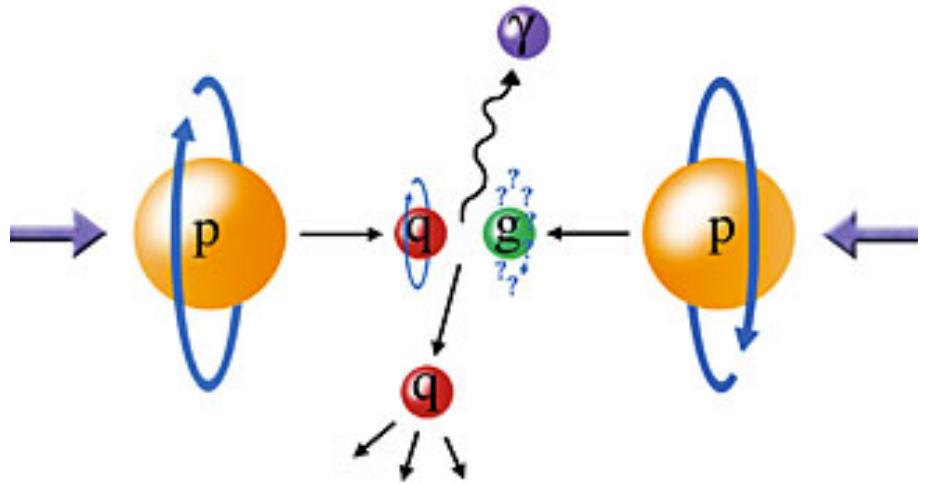
Longitudinally Polarized proton-proton



Double Helicity Asymmetry A_{LL}

$$A_{LL} = \frac{d\sigma^{++} - d\sigma^{+-}}{d\sigma^{++} + d\sigma^{+-}} = \frac{\sum_{a,b} \Delta f_a \otimes \Delta f_b \otimes d\hat{\sigma}^{f_a f_b \rightarrow fX} \cdot \hat{a}_{LL}^{f_a f_b \rightarrow fX} \otimes D_f^h}{\sum_{a,b} f_a \otimes f_b \otimes d\hat{\sigma}^{f_a f_b \rightarrow fX} \otimes D_f^h}$$

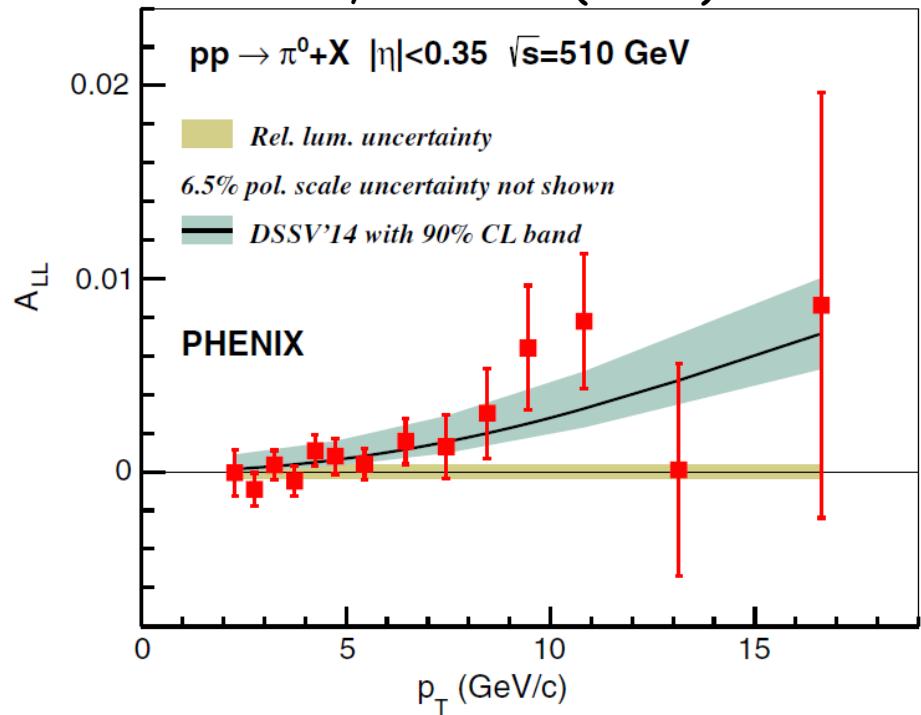
Subprocess of pp



$$A_{LL} \propto [\omega_{gg}] \Delta g \Delta g + [\omega_{gq} \Delta q] \Delta g + [\omega_{qq} \Delta q \Delta q]$$

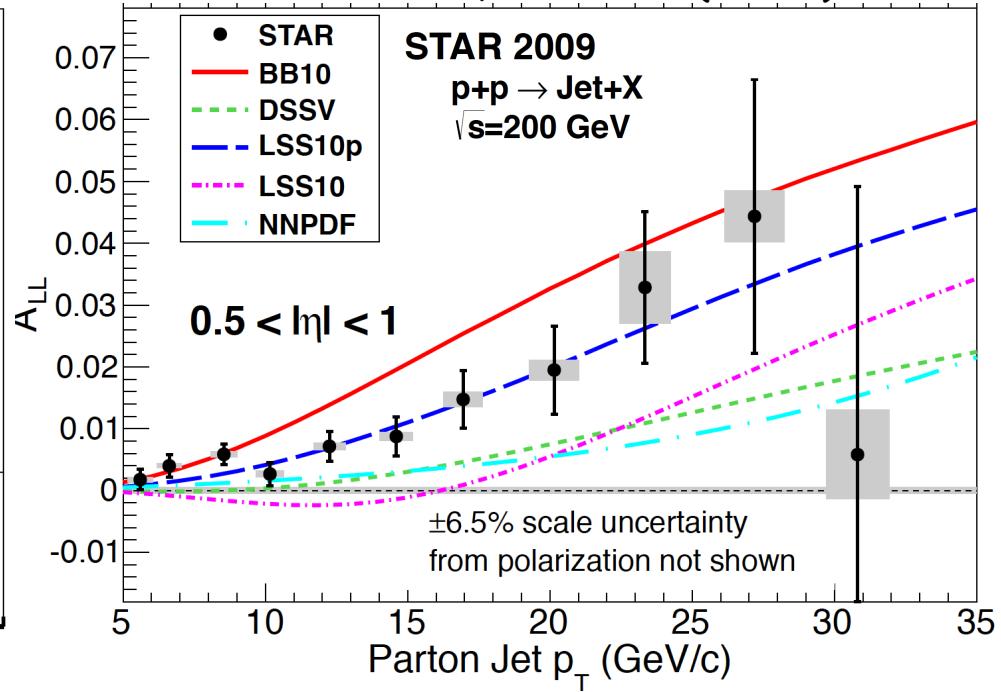
Latest A_{LL} Results

PRD 93, 011501R (2016)



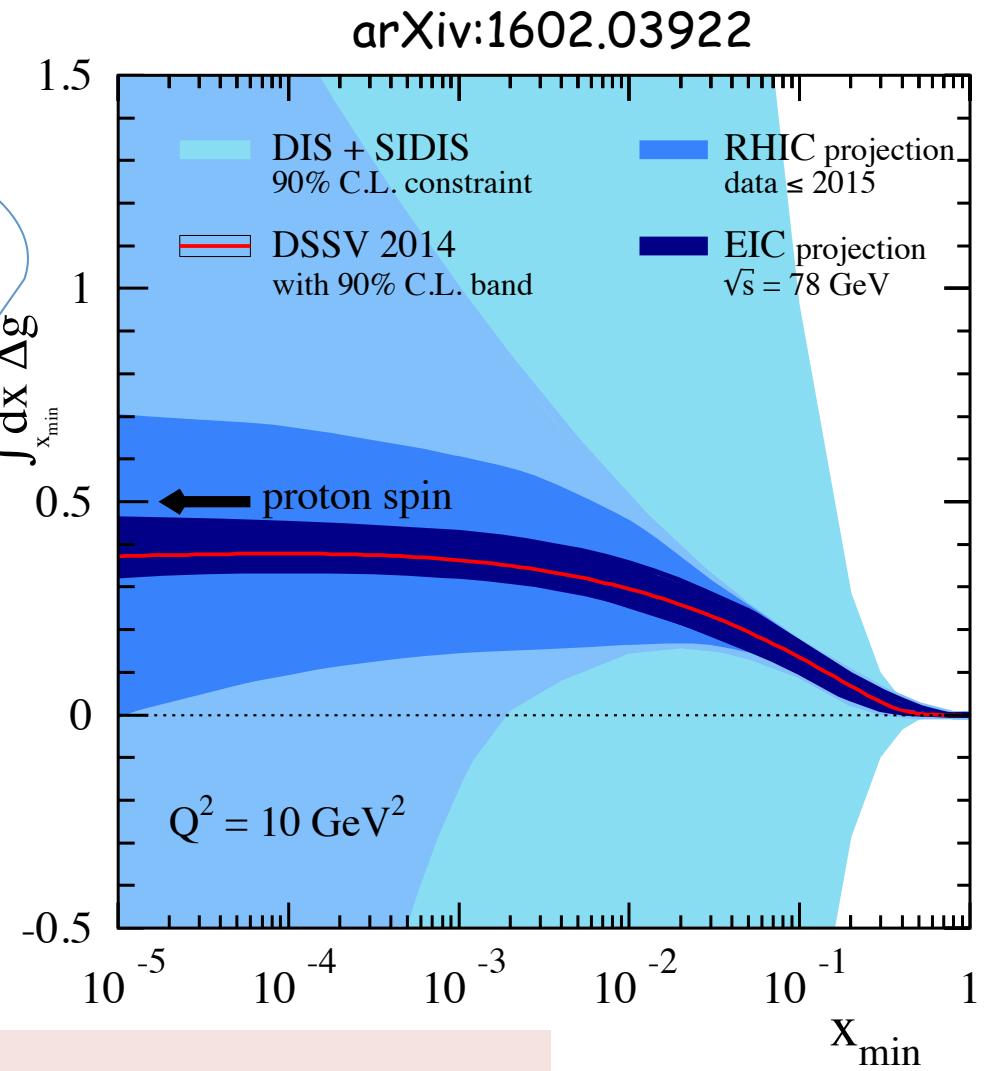
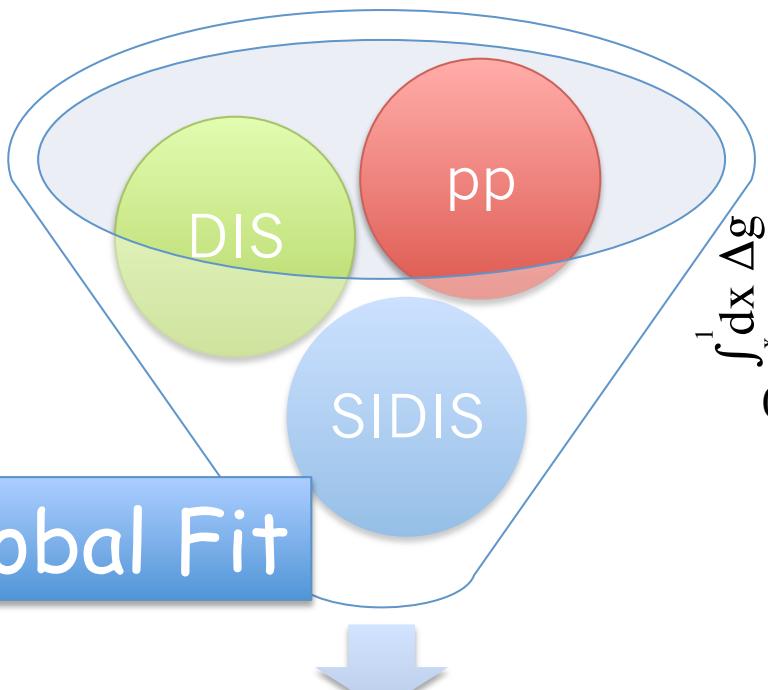
$pp \rightarrow \pi^0 + X$

PRL115, 092002 (2015)



$pp \rightarrow \text{jet} + X$

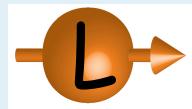
$\Delta g(x)$ extraction : Global Fit



$$\int_{0.05}^{0.2} \Delta g(x) dx = 0.2 \pm^{0.06}_{0.07}$$

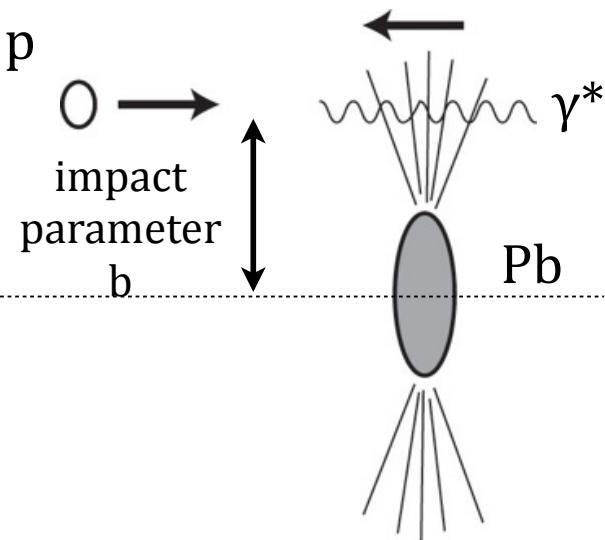
Summary

- RHIC spin program has been providing new insights to spin structure of proton
- Absolutely unexpected strong A-dependence in very forward neutron TSSA of p+A.
- EM process likely to be taken into account.
- Latest data shows non-zero gluon spin fraction



BACKUP SLIDES

Event Generation of UPCs



Event Generation of UPCs

Flux of quasi photons

Cross section of p- γ

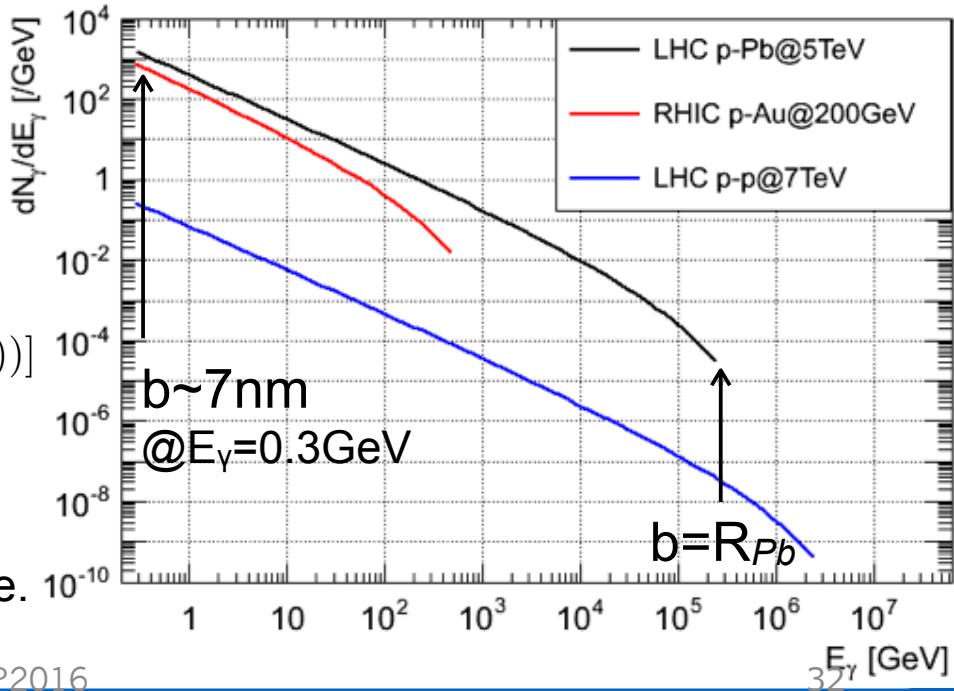
Event Generation of p- γ

Flux of quasi photons

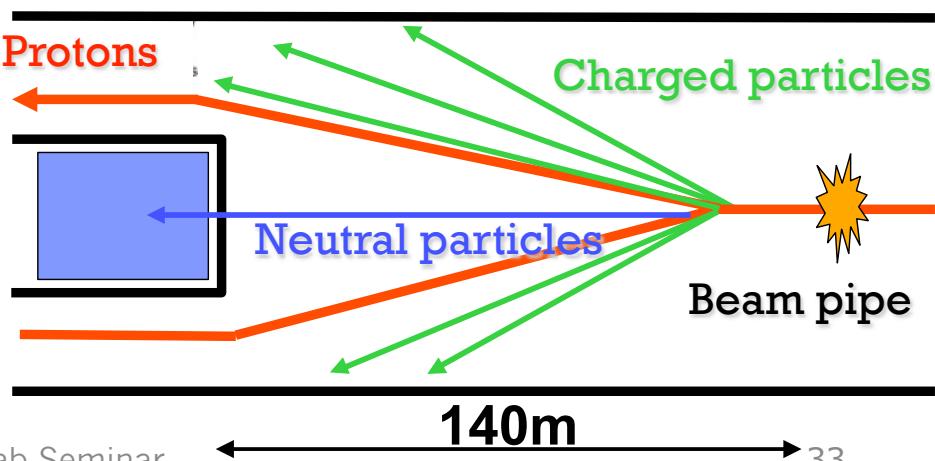
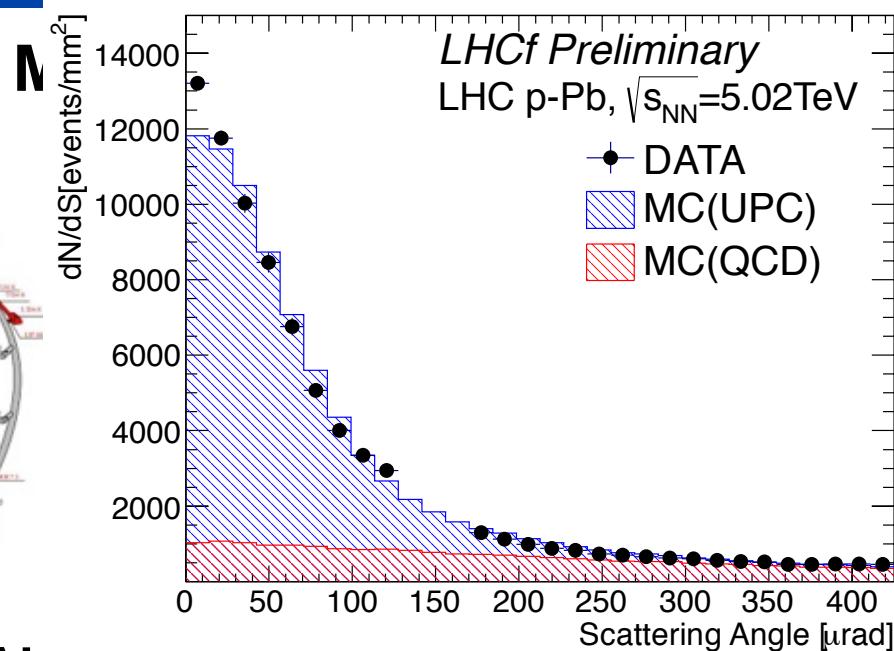
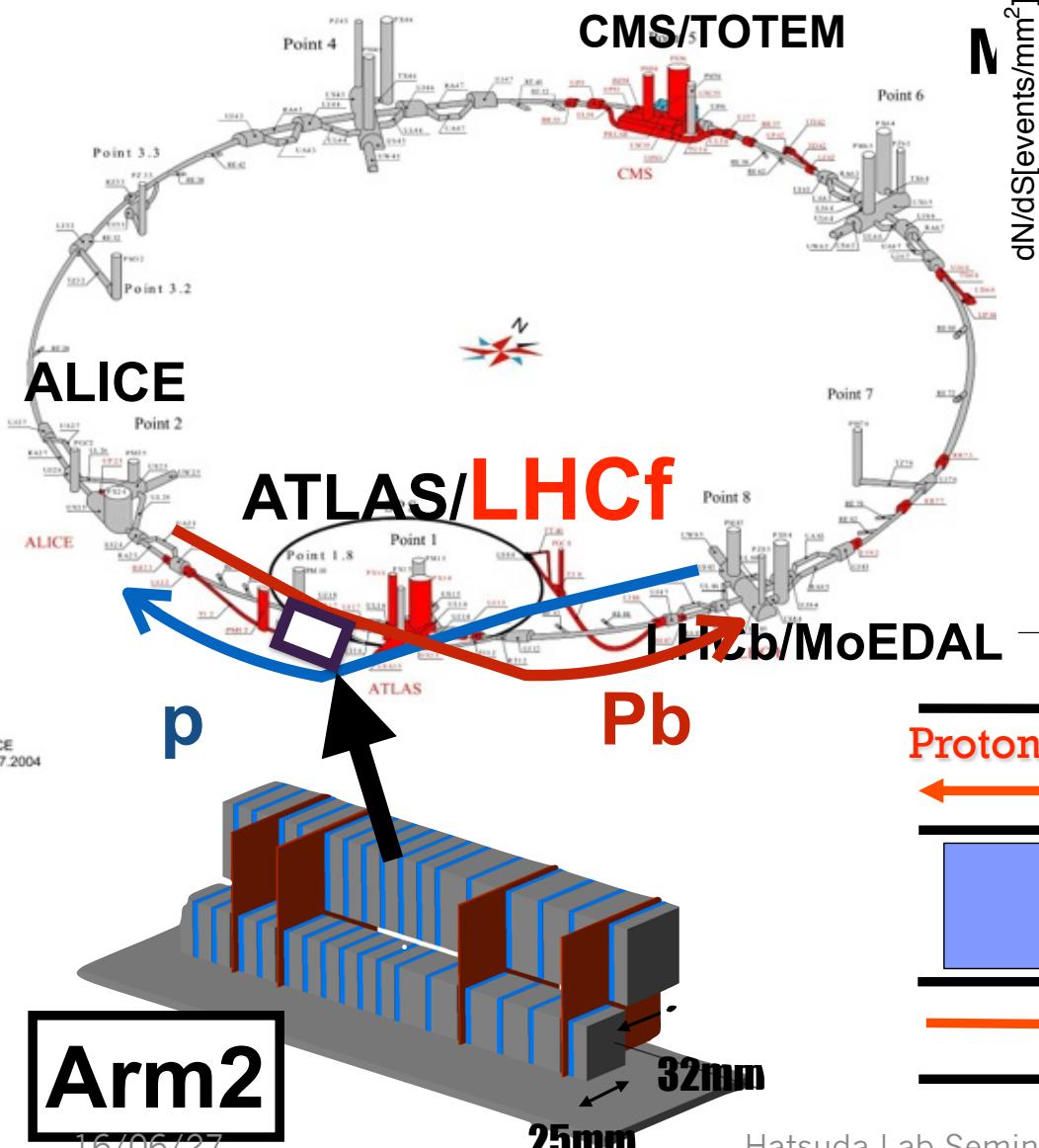
Weizsäcker-Williams method

$$\begin{aligned} \frac{dN_{\gamma^*}}{dE_{\gamma^*}} &= \frac{2Z^2\alpha}{\pi E_{\gamma^*}} [xK_0(x)K_1(x) + \frac{x^2}{2}(K_0^2(x) - K_1^2(x))] \\ &\simeq \frac{2Z^2\alpha}{\pi E_{\gamma^*}} (\log \frac{1.123}{x} - 1) \quad (if E_{\gamma^*} \ll E_{max}) \end{aligned}$$

E_{γ^*} : energy of photons at the proton rest frame.



The LHCf experiment



UPC Monte Carlo

Eur. Phys. J. C (2015) 75:614
DOI 10.1140/epjc/s10052-015-3848-0

THE EUROPEAN
PHYSICAL JOURNAL C



CrossMark

Special Article - Tools for Experiment and Theory

Forward hadron production in ultra-peripheral proton–heavy-ion collisions at the LHC and RHIC

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Università degli Studi di Firenze and INFN Sezione di Firenze, Via Sansone 1, 50019 Sesto Fi-

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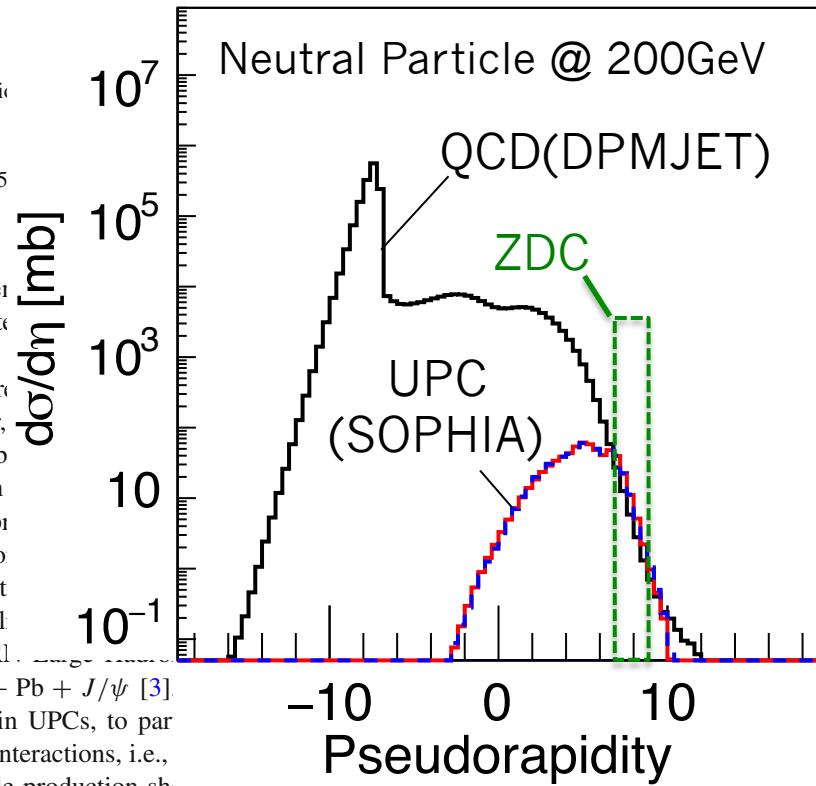
Abstract We present a hadron production study in the forward rapidity region in ultra-peripheral proton–lead ($p + \text{Pb}$) collisions at the LHC and proton–gold ($p + \text{Au}$) collisions at RHIC. The present paper is based on the Monte Carlo simulations of the interactions of a virtual photon emitted by a fast moving nucleus with a proton beam. The simulation consists of two stages: the STARLIGHT event generator simulates the virtual photon flux, which is then coupled to the SOPHIA, DPMJET, and PYTHIA event generators for the simulation of particle production. According to these Monte Carlo simulations, we find large cross sections for ultra-peripheral collisions particle production, especially in the very forward region. We show the rapidity distributions for charged and neutral particles, and the momentum distributions for neutral pions and neutrons at high rapidities. These processes lead to substantial background contributions to the investigations of collective nuclear effects and spin physics. Finally we pro-

virtual photons $e\gamma$ may anyway interact with nuclei, usually referred to as Ref. [1,2] for a review.

Up to now, the gluon distribution function in photoproduction collisions can predict the density in proton–nucleus collisions at the momentum fraction x measurements at the CERN NA3 experiment at the CERN NA3 experiment.

$p + \text{Pb} \rightarrow p + \text{Pb} + J/\psi$ [3]

has been paid, in UPCs, to pair production of photon–proton interactions, i.e., less such particle production should be considered as well in the investigation of collective nuclear effects. Because of

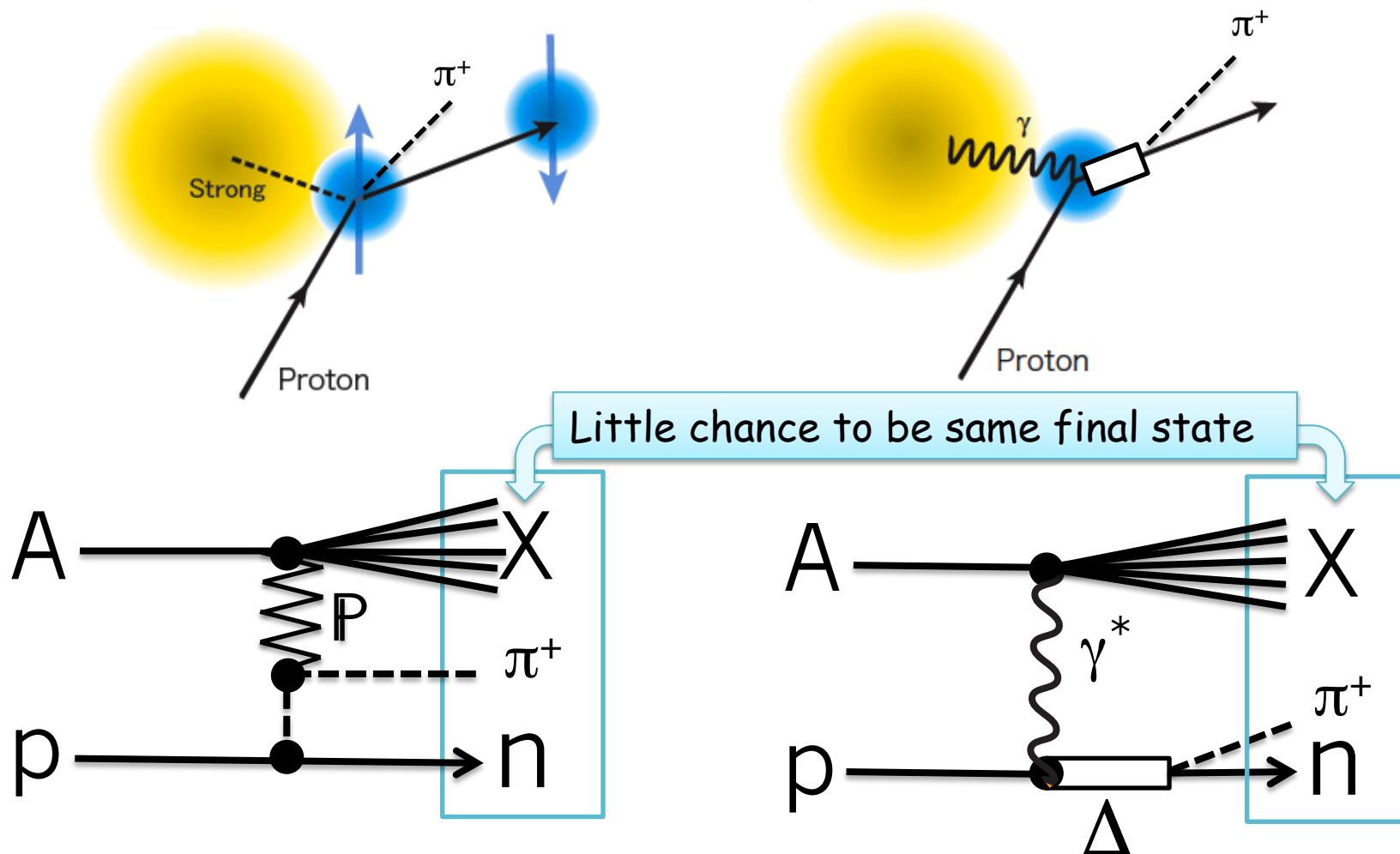


Predicts comparable yields between QCD and UPC processes

16/06/27

Hatsuda Lab Seminar

Non-Diffractive Events

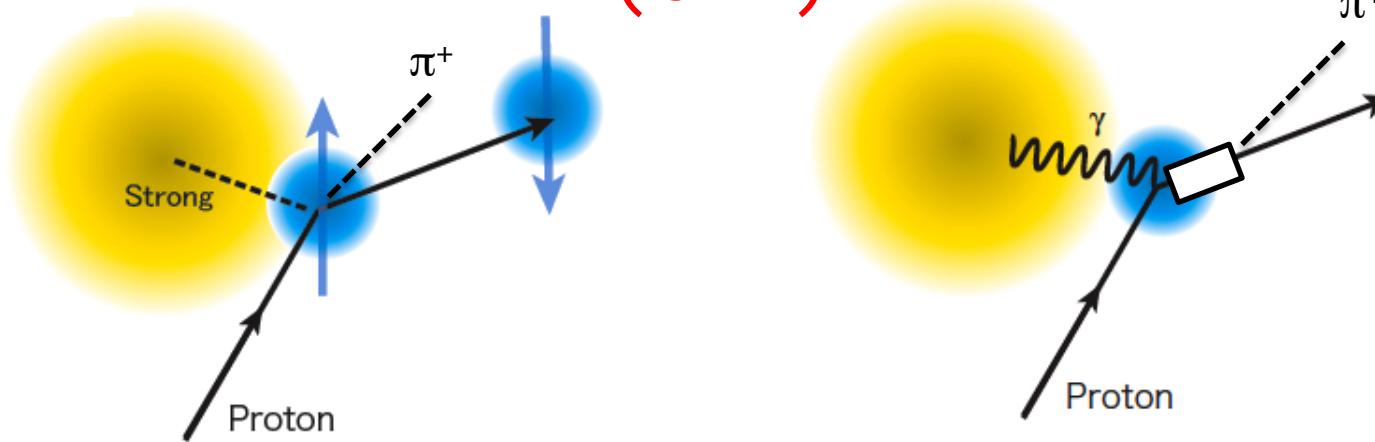


$$A_N \sim \text{had} * \text{had} + \boxed{\text{had} * \text{EM} + \text{EM} * \text{had}} + \text{EM} * \text{EM}$$

Suppressed?

Coulomb-Nuclear Interference

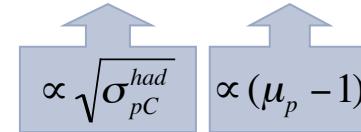
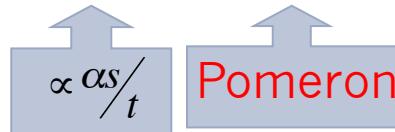
(CNI)



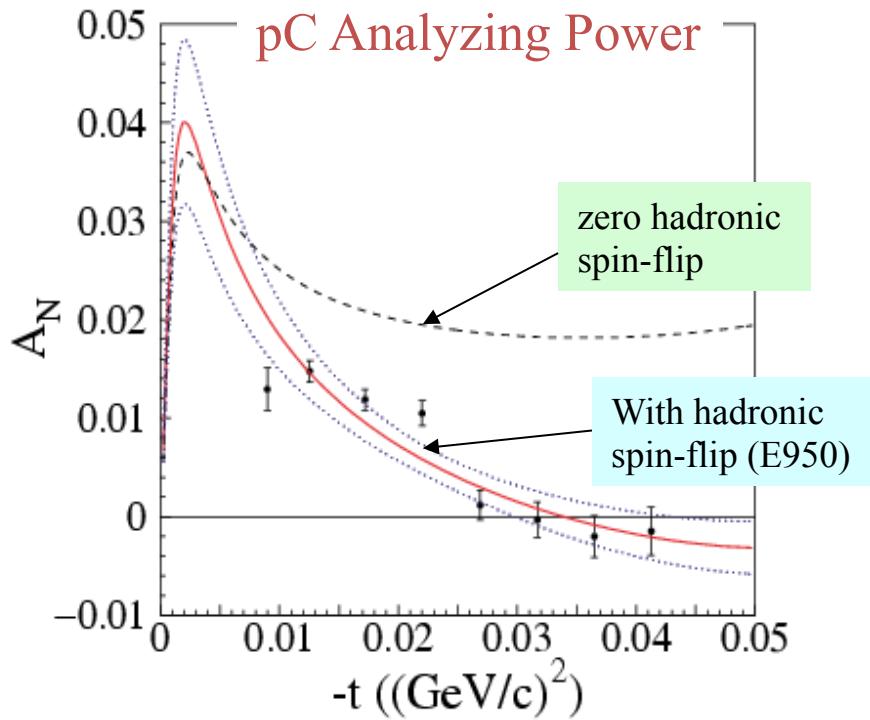
$$A_N \sim had * had + had * EM + EM * had + EM * EM$$

Elastic A_N at Coulomb Nuclear Interference

$$A_N \sim 2 \operatorname{Im} \left\{ \phi_{non-flip}^{EM^*} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had^*} \phi_{flip}^{EM} \sin \delta_3 \right\}$$



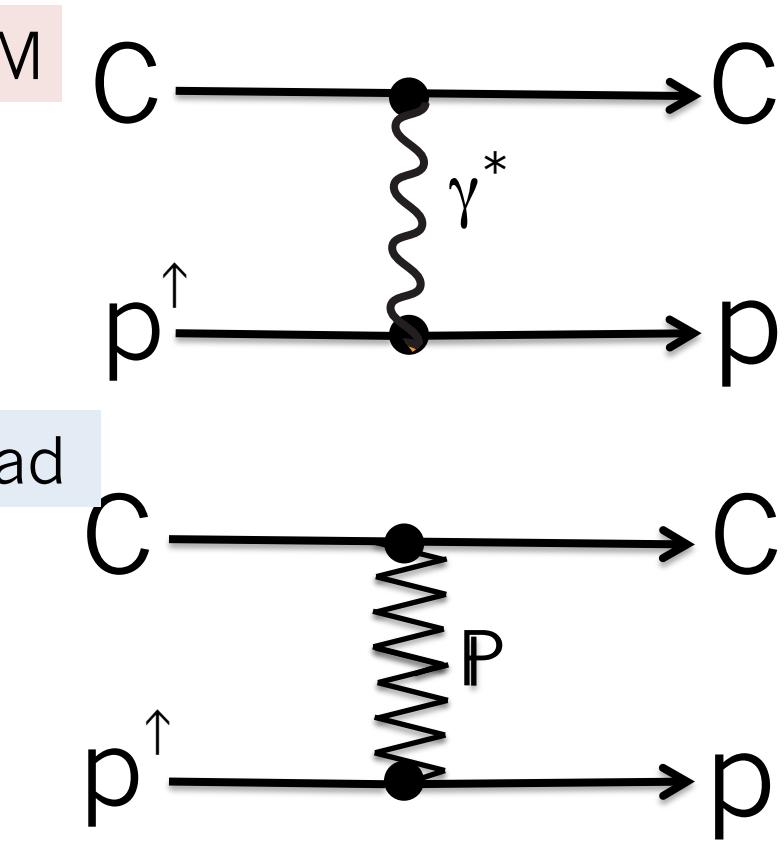
(High energy & small t limit)



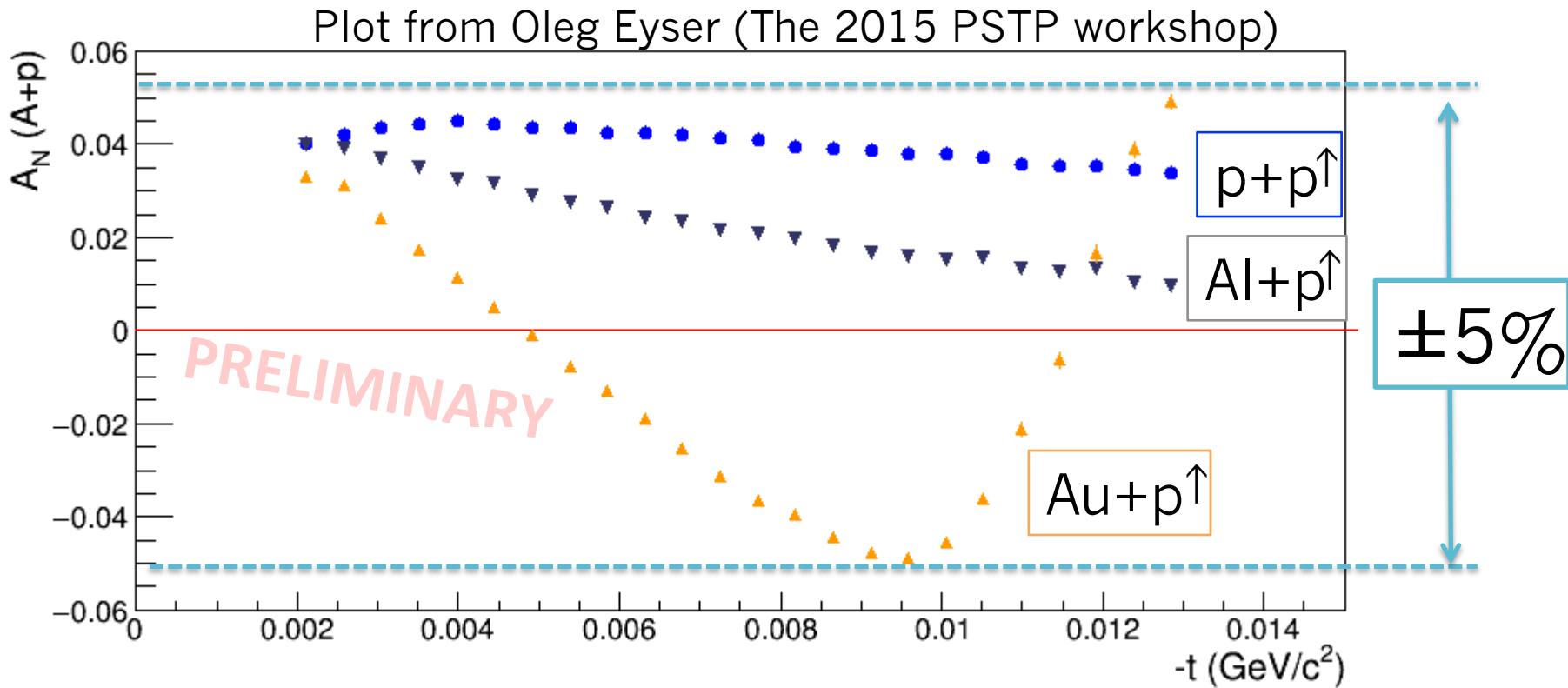
Phys. Rev. Lett., 89, 052302 (2002)

16/06/27 $E_{beam} = 21.7 \text{ GeV}$

Hatsuda Lab Seminar



Run15 Au,Al beam + p \uparrow target

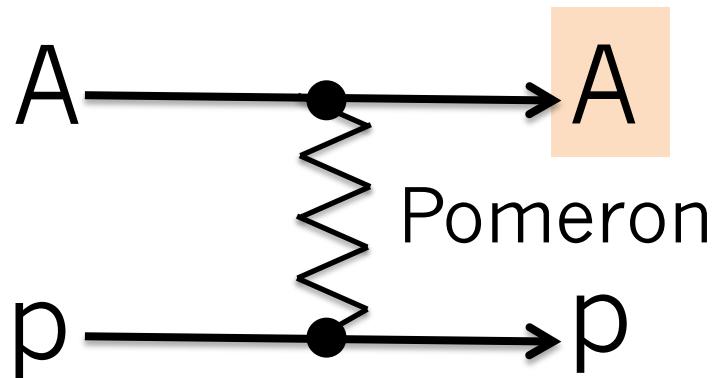


- Strong A-Dependence
- Flips sign of A_N in $Au+p^\uparrow$
- $0.002 < -t < 0.014 \text{ (Gev/c)}^2$

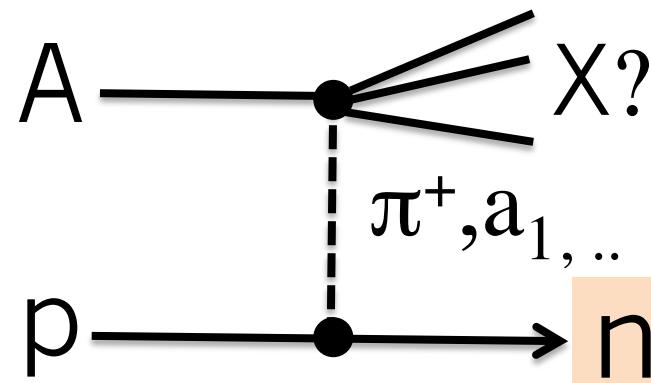
Forward Neutron
 $0.02 < -t < 0.5 \text{ (Gev/c)}^2$

Underlying Mechanism Comparison

Polarimeter



Forward n



Elastic	Inelastic
$\sqrt{s} = 14 \text{ GeV}$	$\sqrt{s} = 200 \text{ GeV}$
$0.002 < -t < 0.014$	$0.02 < -t < 0.5$
$\Delta I = 0$	$\Delta I = 1$

Interim Summary (CNI)

- CNI is not necessarily new.
- Elastic in much smaller $-t$ than present neutron results observed similar behavior, but the asymmetry is in the order of 5% or less.
- There are possible diagram to hadron and EM interfere for neutron production as well.

**Measurement of the Analyzing Power in the Primakoff Process with a
High-Energy Polarized Proton Beam**

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The origin of Asymmetry

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cording to Eq. (1), the asymmetry seen in photoproduction due to the interference between Δ and N^* is expected in coherent Coulomb π^0 production by polarized protons, using the same region of the π^0 - p invariant mass. Therefore this process may be used to measure the polarization of the proton at high energies.⁶ Until now, there has been no measurement of the asymmetry in the nuclear coherent process.

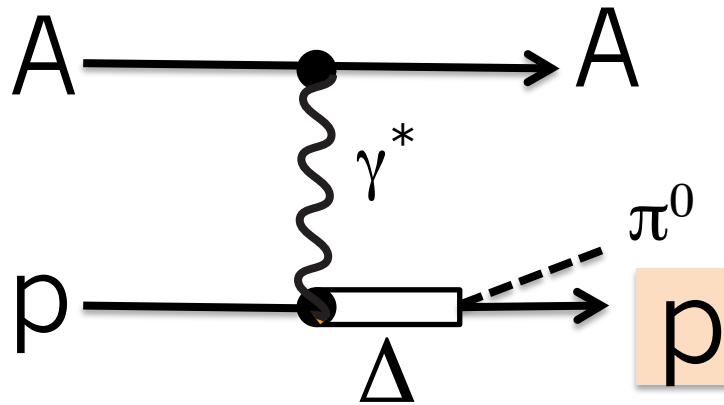
The cross section for the Coulomb coherent process (1) has a sharp peak at $t' \sim 10^{-5}$ (GeV/c)² and decreases rapidly as t'/t^2 . The "width" of the Coulomb peak is determined by the detector resolution. Diffractive dissociation due to the strong interaction is also present, but it has a much slower t' dependence.

We have measured the analyzing power (azimuthal asymmetry) of nuclear Coulomb coherent production from a Pb target by using the newly constructed 185-GeV/c Fermilab polarized proton beam.⁷ The beam polarization is 45% and this is further described in Ref. 7. To reduce certain systematic errors, the spin direction of the incident proton was flipped every 10 min using a spin-rotator system.⁷

monitored to a 1% accuracy with a xenon-flash-tube system. A 30-GeV positron beam was used to calibrate the calorimeter. The measured energy resolution is 3% (rms) at 30 GeV and the position resolution is 2 mm (rms). The measured π^0 energies in this experiment ranged from 25 to 75 GeV.

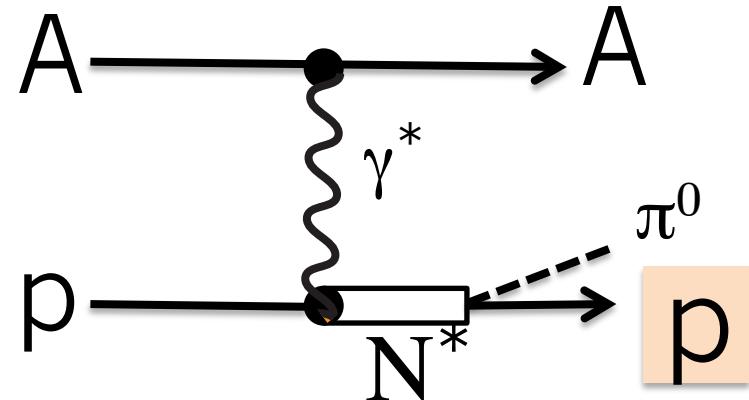
A set of thin plastic scintillation counters (TP1) is placed downstream of the magnet and provides the trigger for the scattered protons. The set consists of four counters arranged to distinguish protons scattered to the left, right, up, and down. The calorimeter also has left, right, up, and down sections, and signals from each section are summed for the trigger. In the coherent process where t' is almost zero, the π^0 and scattered protons are coplanar. Thus the trigger logic is such that the energy deposit is larger than 25 GeV in the left half of the calorimeter, less than 5 GeV in the right half, and a proton hits the right segment of TP1. There are four such combinations to cover the whole range of azimuthal angles. To reject the events which have any extra particle besides a proton and π^0 , veto counters are included in the trigger logic.

Origin of the Asymmetry



$\Delta^*(1232)$ P33

↳ $N\pi$ ($>99\%$)



$N^*(1440)$ P11

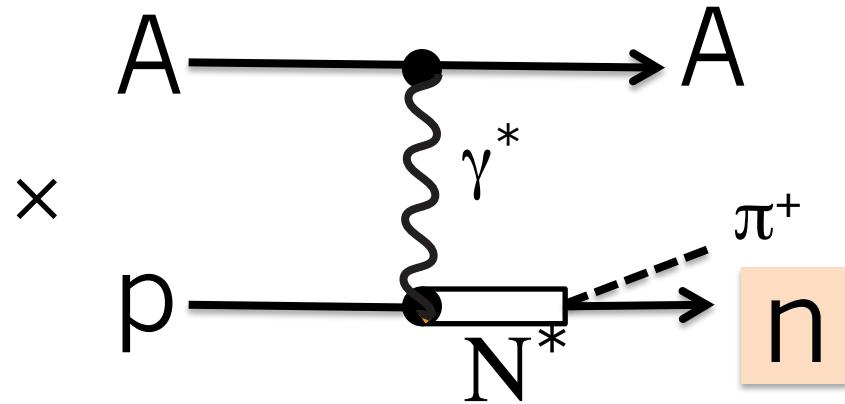
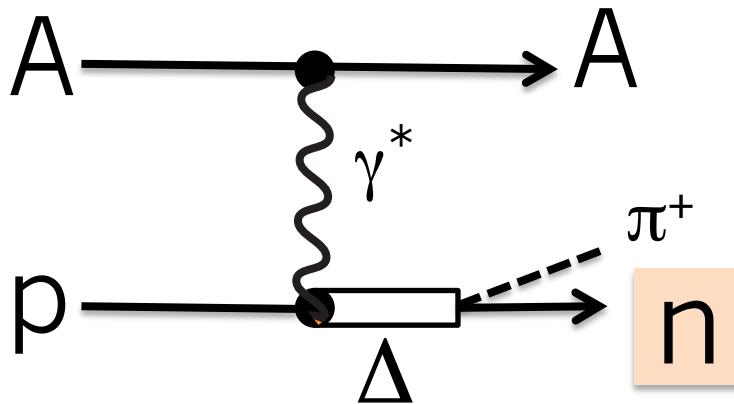
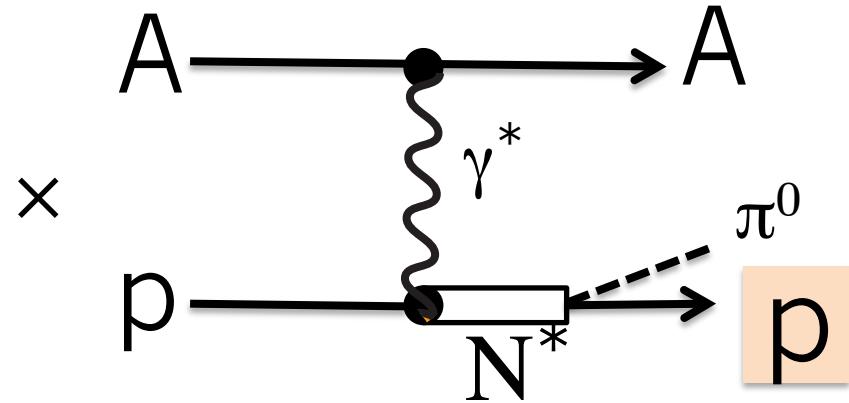
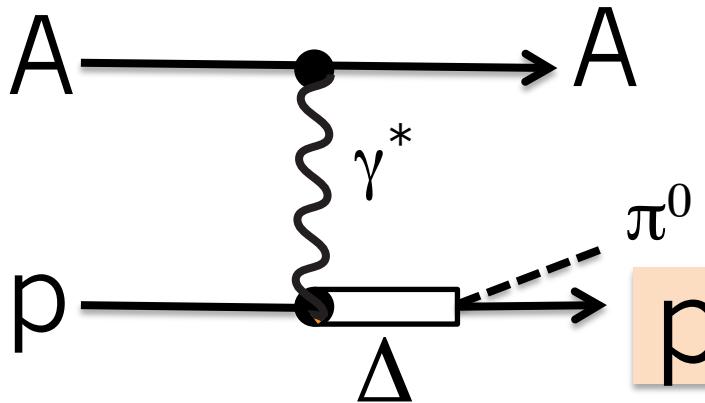
↳ $N\pi$ ($60\sim70\%$)
 $N\pi\pi$ ($30\sim40\%$)

$$A_N \propto \phi_{non-flip}^{had} \phi_{flip}^{had} \sin \delta_1 + \phi_{non-flip}^{EM} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had} \phi_{flip}^{EM} \sin \delta_3 + \boxed{\phi_{non-flip}^{EM} \phi_{flip}^{EM} \sin \delta_4}$$

Comparison between two

	Fermi	PHENIX	STAR17
Beam Energy [GeV]	185	100	255
\sqrt{s} [GeV]	19.5	200	22
Target	Pb	Au	Al/Sn/Au
Observables	$p + \pi^0$	$n (+ \text{ charged})$	$n (+ \text{ charged})$ $\pi^0 ?$
t'	< 0.001	$0.02 < -t < 0.5$	
M	$1.36 < M(\pi^0 p) < 1.52$?	?
A_N	$-0.57 \pm (0.12)_{\text{sta}} + 0.21 - 0.18$	$+0.27 \pm 0.003$ (BBC veto)	

Origin of the Asymmetry

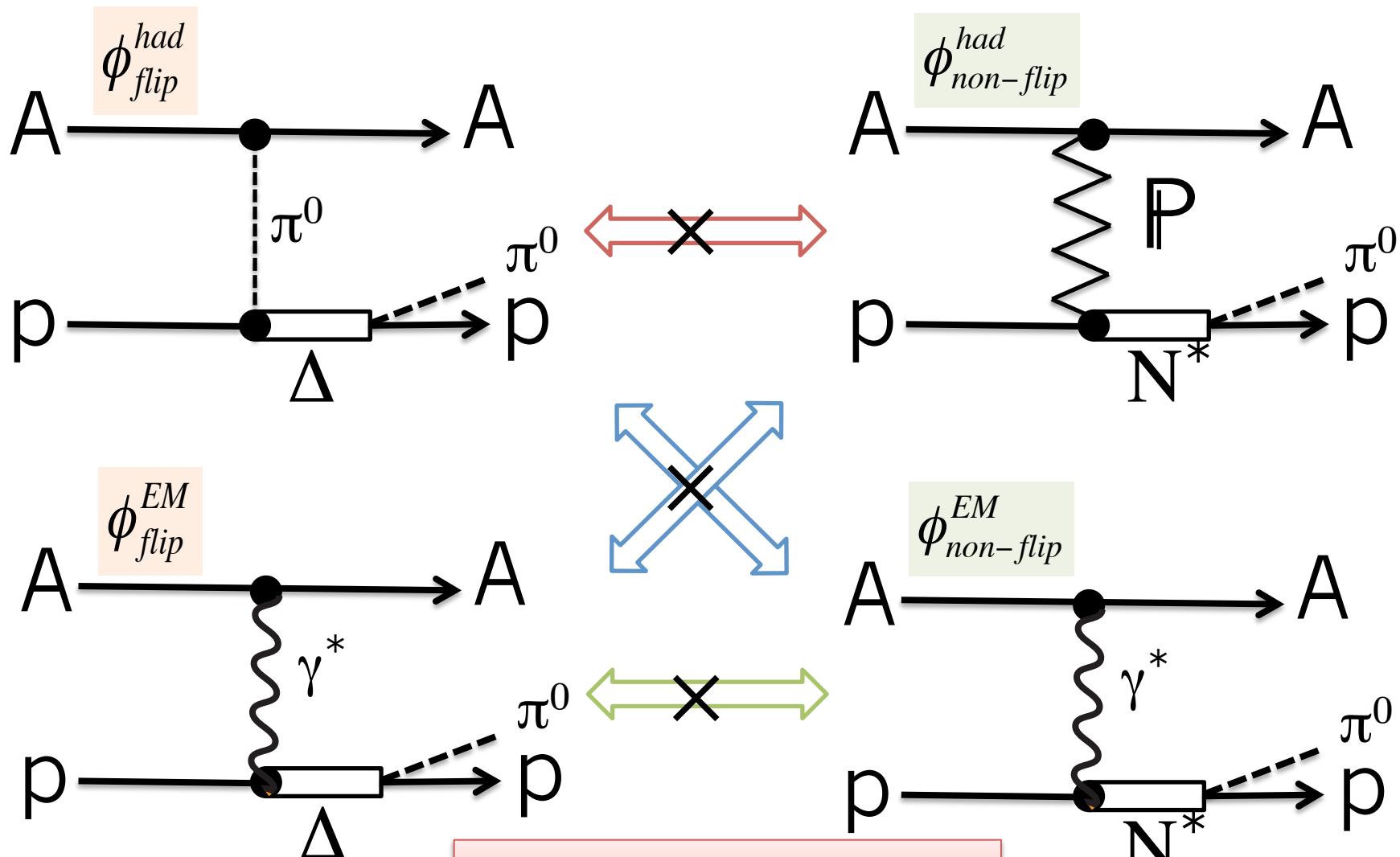


$$A_N \propto \phi_{non-flip}^{had} \phi_{flip}^{had} \sin \delta_1 + \phi_{non-flip}^{EM} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had} \phi_{flip}^{EM} \sin \delta_3 + \boxed{\phi_{non-flip}^{EM} \phi_{flip}^{EM} \sin \delta_4}$$

45

$$A_N \propto \phi_{non-flip}^{had} \phi_{flip}^{had} \sin \delta_1 + \phi_{non-flip}^{EM} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had} \phi_{flip}^{EM} \sin \delta_3 + \phi_{non-flip}^{EM} \phi_{flip}^{EM} \sin \delta_4$$

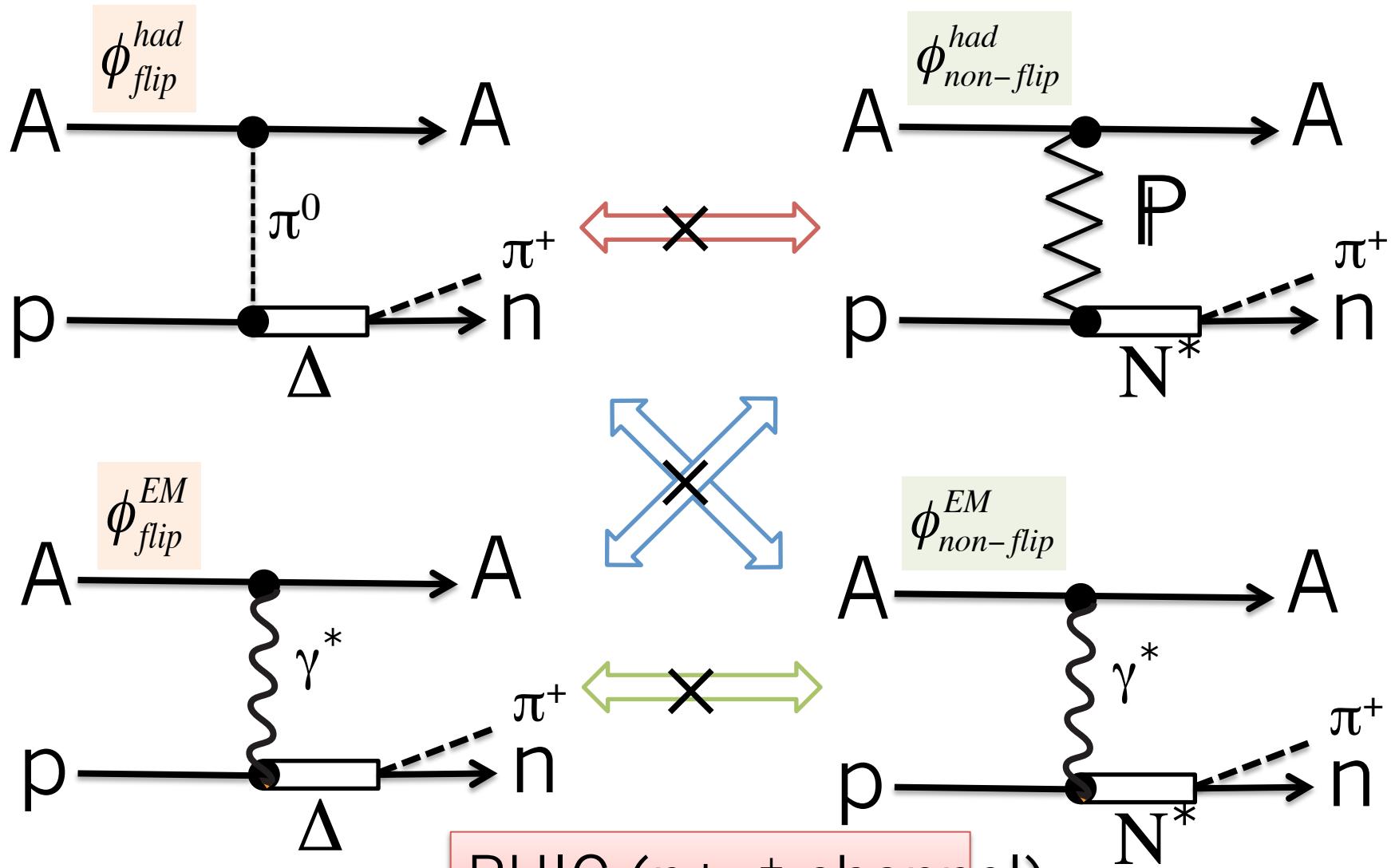
hadronic CNI Primakoff



Fermi ($p + \pi^0$ channel)

$$A_N \propto \phi_{non-flip}^{had} \phi_{flip}^{had} \sin \delta_1 + \phi_{non-flip}^{EM} \phi_{flip}^{had} \sin \delta_2 + \phi_{non-flip}^{had} \phi_{flip}^{EM} \sin \delta_3 + \phi_{non-flip}^{EM} \phi_{flip}^{EM} \sin \delta_4$$

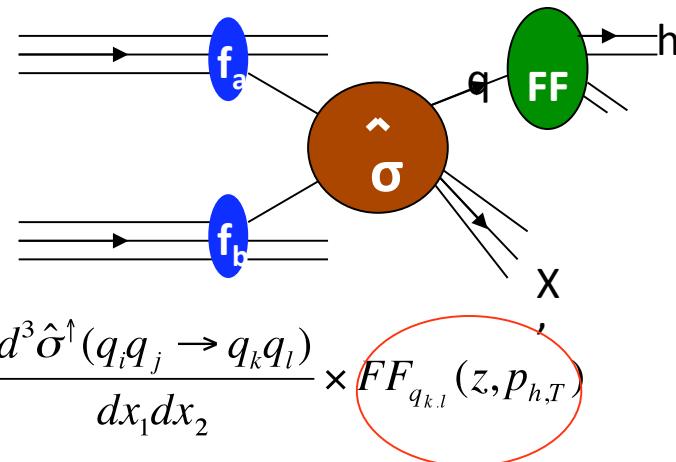
hadronic CNI Primakoff



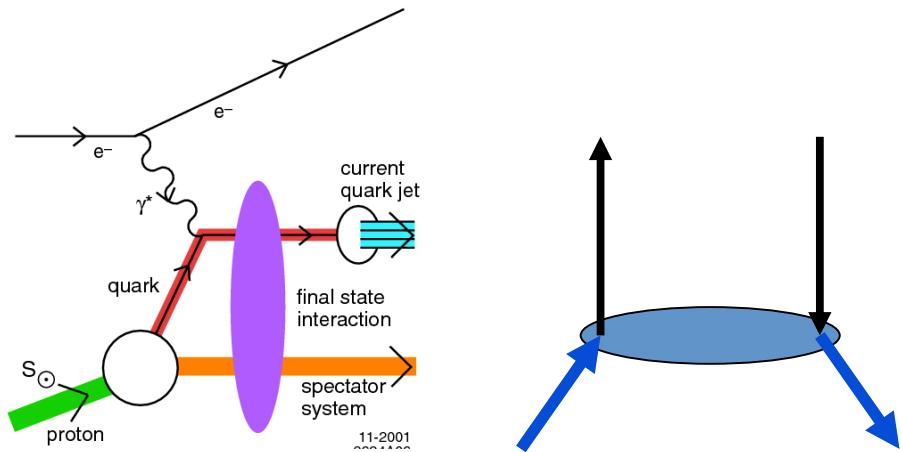
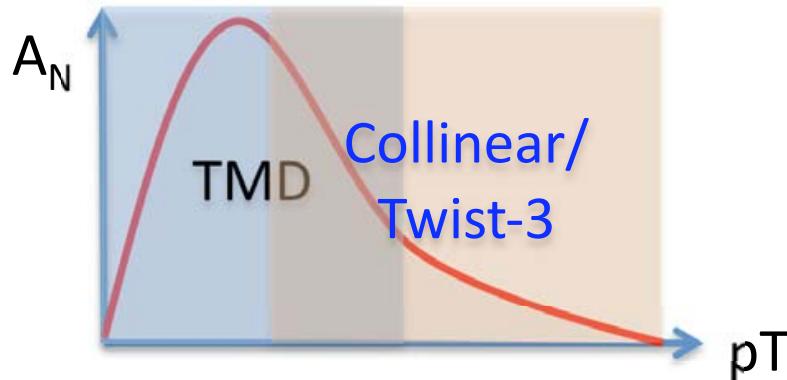
Theory: K_T vs Collinear Factorization

- Tran. Mom. Dep. Funs
 - Sivers Fun
 - Collins Fun

$$\frac{d^3\sigma^\uparrow(pp^\uparrow \rightarrow h + X)}{dx_1 dx_2 dz} \propto q_i^\uparrow(x_1, k_{q,T}) \cdot q_j(x_2) \times \frac{d^3\hat{\sigma}^\uparrow(q_i q_j \rightarrow q_k q_l)}{dx_1 dx_2} \times FF_{q_{k,l}}(z, p_{h,T})$$

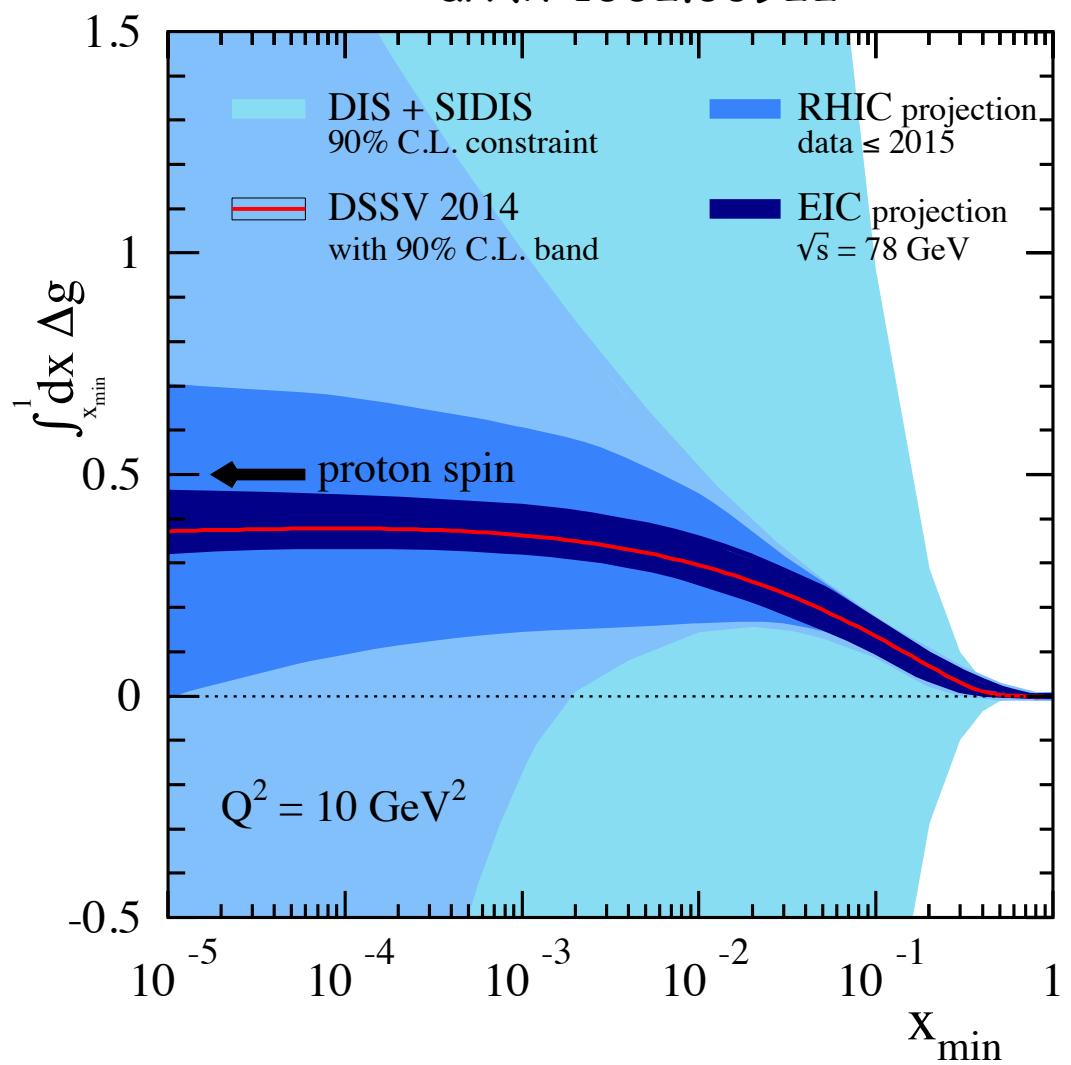


- Twist-3 collinear
 - Quark-gluon correl.
 - Gluon-gluon correl.



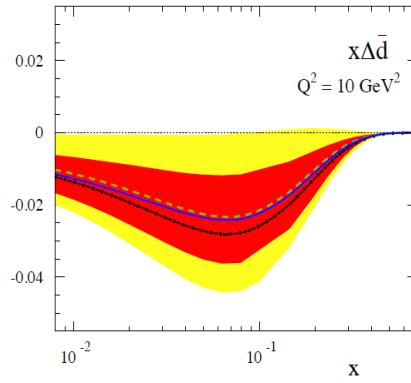
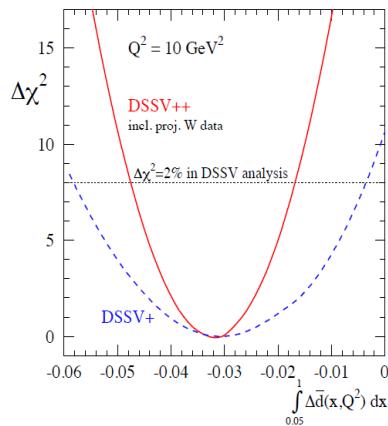
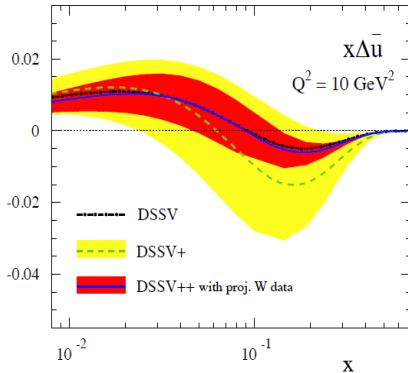
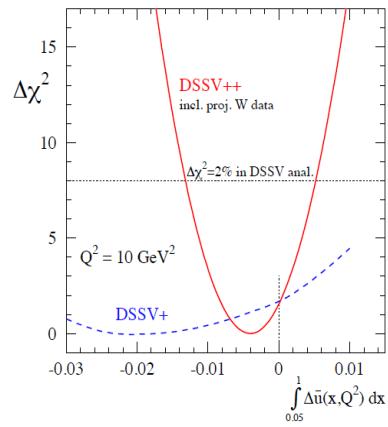
Future

arXiv:1602.03922

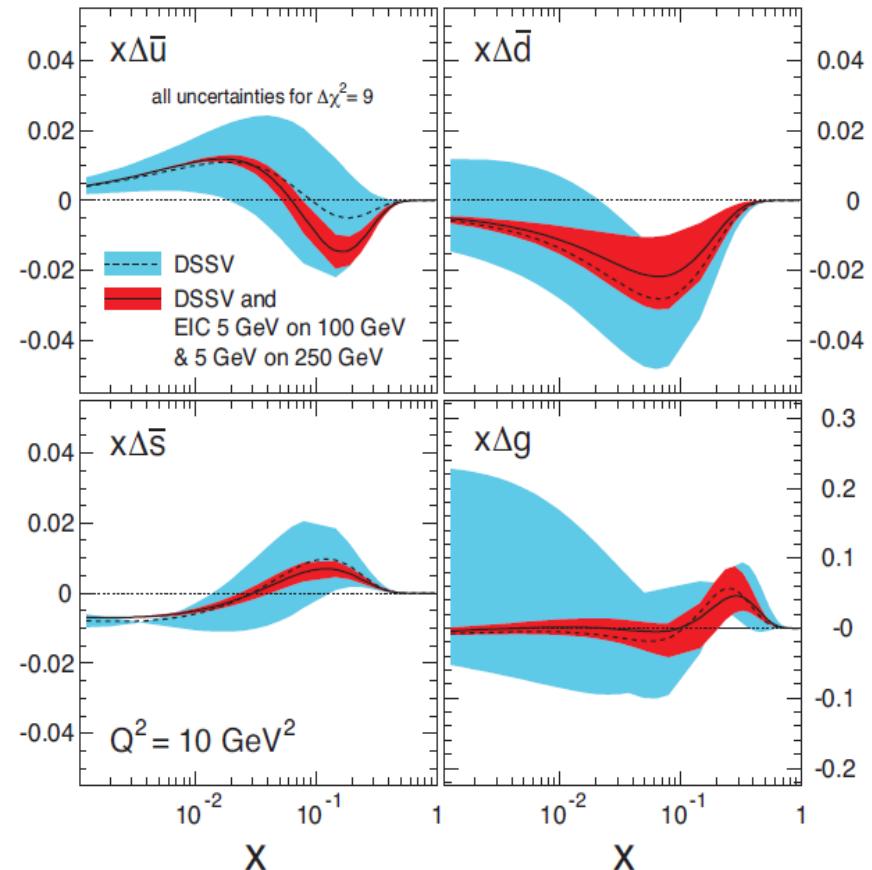


Present and Future

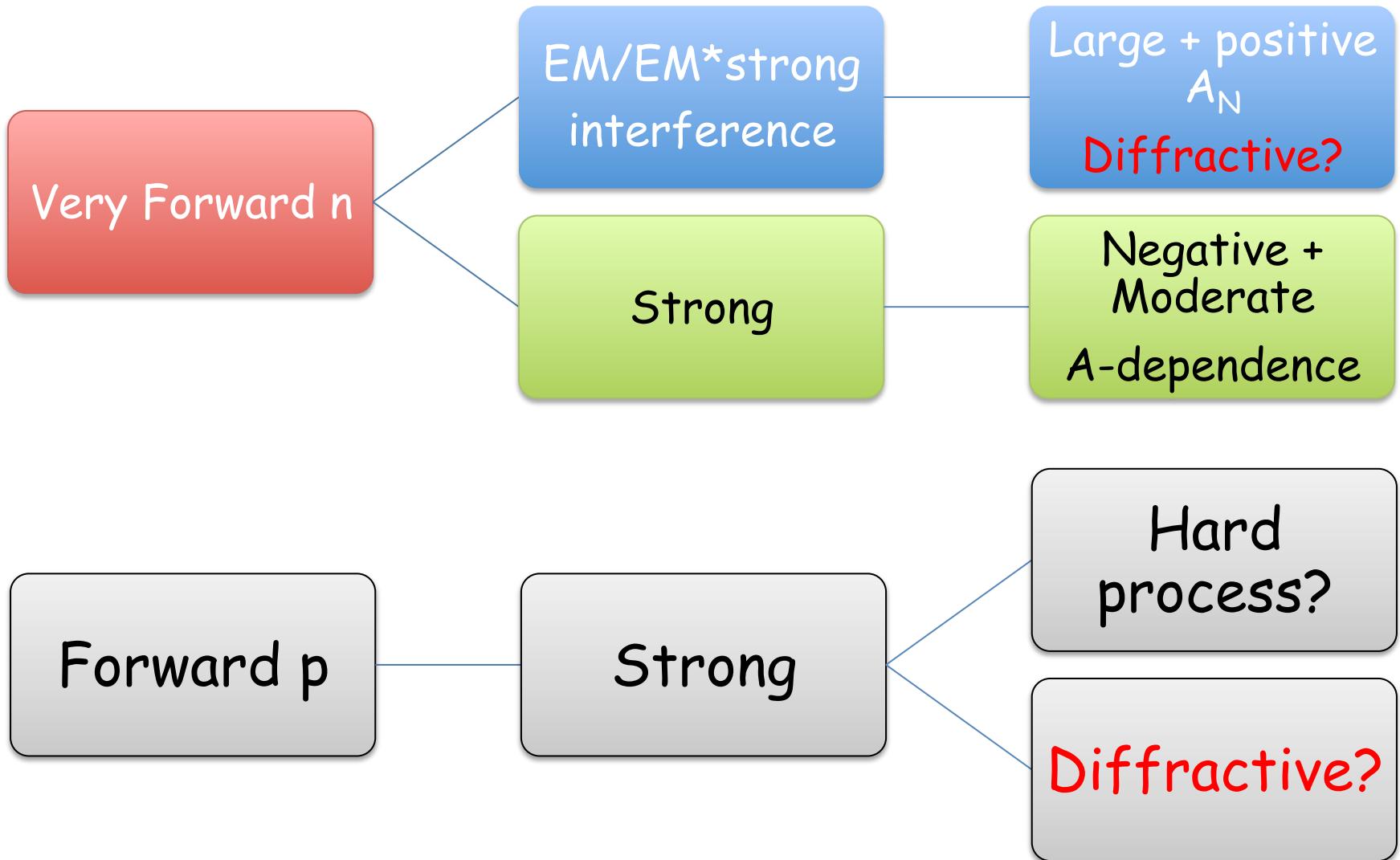
arXiv: 1501.01220



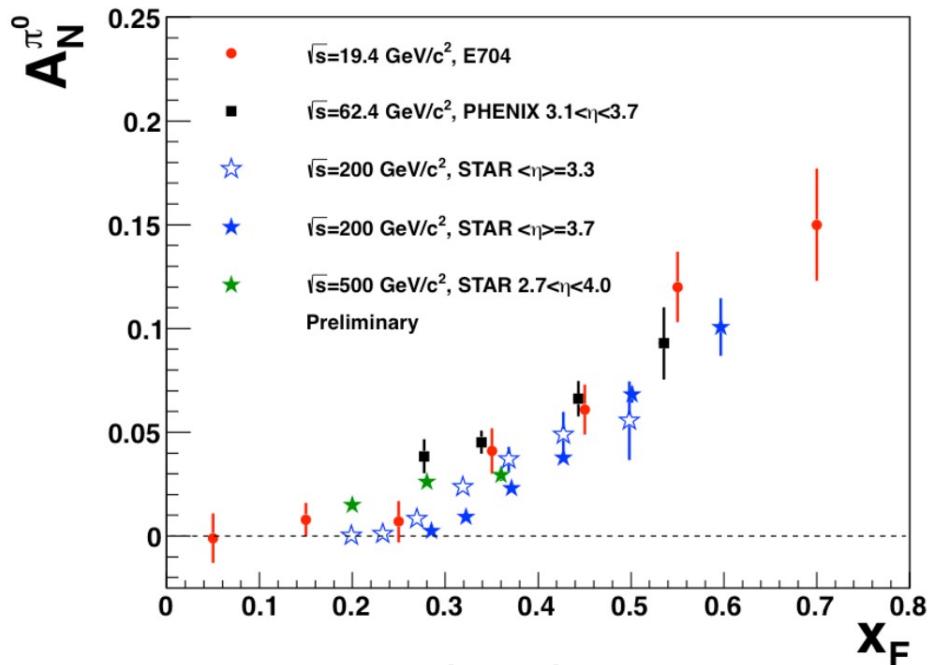
arXiv: 1212.1701



Transverse Summary

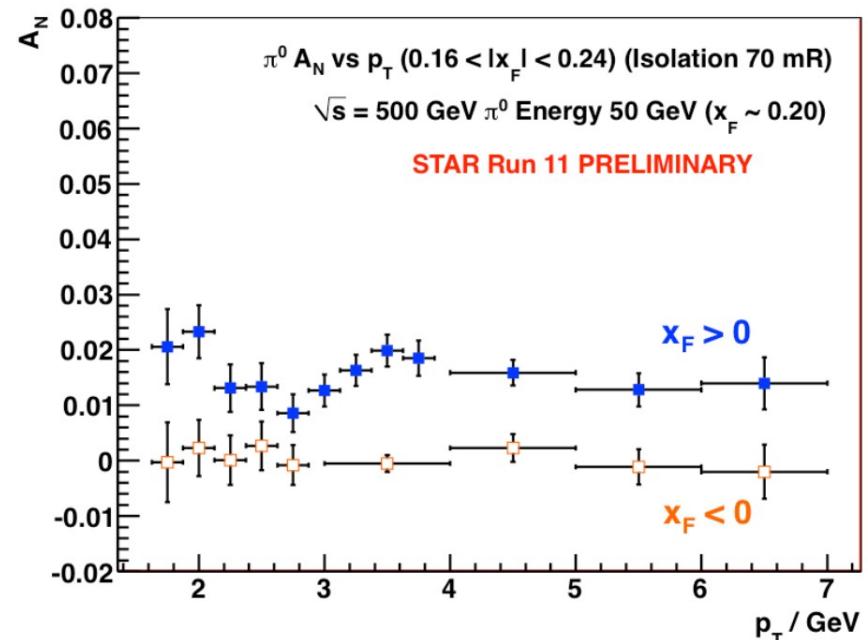


x_F and p_T Dependence



PLB,261,201(1991)
PRL,101,222001(2008)
PRD,90,012006(2014)

$s^{1/2}=200 \text{ GeV}$ and 500 GeV show same rise of A_N vs. x_F as lower $s^{1/2}$ measurements

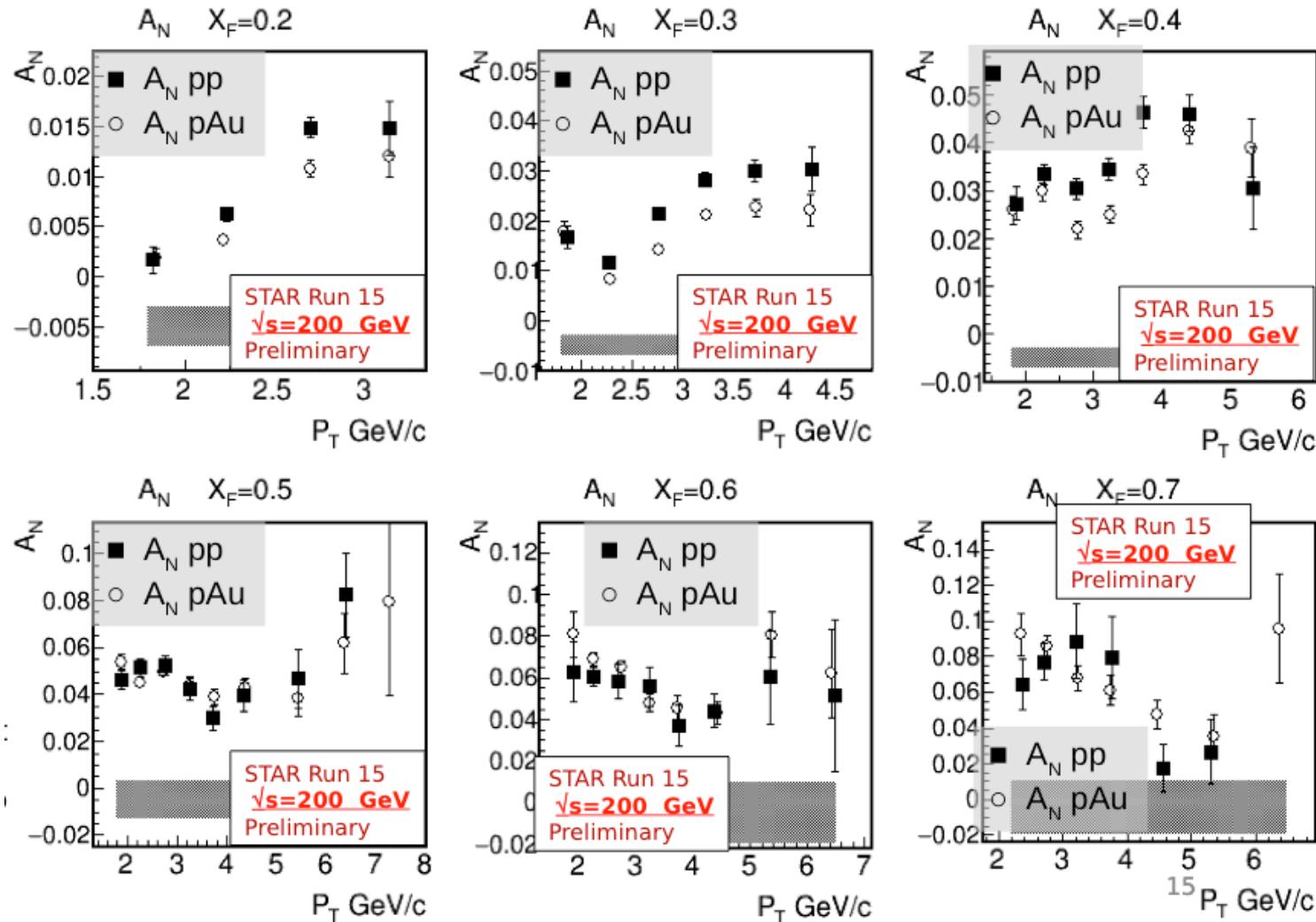


Steve Heppelmann – CIPANP 2012

- Collins, Sivers, Twist-3 suggest $A_N \sim 1/p_T$
- Flat p_T -dependence observed and raises the question as to what causes it



Forward $\pi^0 A_N$ in p+Au



Did not observe A -dependence unlike very forward neutron A_N in PHENIX